

Boundedness of the gradient of a solution to the Neumann-Laplace problem in a convex domain

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Abstract. It is shown that solutions of the Neumann problem for the Poisson equation in an arbitrary convex n -dimensional domain are uniformly Lipschitz. Applications of this result to some aspects of regularity of solutions to the Neumann problem on convex polyhedra are given.

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Key words: Neumann problem, Laplace operator, Sobolev spaces, convex domain, boundedness of the gradient, eigenvalues of the Neumann-Beltrami operator

Résumé. On démontre que les solutions du problème de Neumann pour l'équation de Poisson dans un domaine convexe arbitraire de dimension n sont uniformément Lipschitz. Les applications de ce résultat à quelques aspects de régularité de solutions du problème de Neumann sur les polyèdres convexes sont données.

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Mots clé: problème de Neumann, opérateur de Laplace, espaces de Sobolev, domaine convexe; bornitude du gradient, valeurs propres de l'opérateur de Neumann-Beltrami.

1 Introduction

Let Ω be a bounded convex domain in \mathbb{R}^n and let $W^{l,p}(\Omega)$ stand for the Sobolev space of functions in $L^p(\Omega)$ with distributional derivatives of order l in $L^p(\Omega)$. By $L^p_{\perp}(\Omega)$ and $W^{l,p}_{\perp}(\Omega)$ we denote the subspaces of functions v in $L^p(\Omega)$ and $W^{l,p}(\Omega)$ subject to $\int_{\Omega} v dx = 0$.

Let $f \in L^2_{\perp}(\Omega)$ and let u be the unique function in $W^{1,2}(\Omega)$, also orthogonal to 1 in $L^2(\Omega)$, and satisfying the Neumann problem

$$-\Delta u = f \quad \text{in } \Omega, \quad (1)$$

$$\frac{\partial u}{\partial \nu} = 0 \quad \text{on } \partial\Omega, \quad (2)$$

where ν is the unit outward normal vector to $\partial\Omega$ and the problem (1), (2) is understood in the variational sense. It is well known that the inverse mapping

$$L^2_{\perp}(\Omega) \ni f \rightarrow u \in W^{1,2}_{\perp}(\Omega) \quad (3)$$

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is continuous. (Since any attempt at reviewing the rich history of this fact and other ones, closely related to it, within frames of a short article is hopeless, I restrict myself to a number of relevant references [3], [4], [11], [13], [16]–[18], [20], [23], [28], [29].) As shown in [2] (see also [12] for a different proof, and [1], [8]–[10] for the Dirichlet problem), the operator

$$L_{\perp}^p(\Omega) \ni f \rightarrow u \in W_{\perp}^{2,p}(\Omega) \quad (4)$$

is also continuous if $1 < p < 2$. One cannot guarantee the continuity of (4) for any $p \in (2, \infty)$ without additional information about the domain. The situation is the same as in the case of the Dirichlet problem (see [4], [8]–[10]), which, moreover, possesses the following useful property: if Ω is convex, the gradient of the solution is uniformly bounded provided the right-hand side of the equation is good enough. This property can be easily checked by using a simple barrier. Other approaches to similar results were exploited in [21] and [14] for different equations and systems but only for the Dirichlet boundary conditions.

In this respect, other boundary value problems are in a nonsatisfactory state. For instance, it was unknown up to now whether solutions of the problem (1), (2) with a smooth f are uniformly Lipschitz under the only condition of convexity of Ω .

The main result of the present paper is *the boundedness of $|\nabla u|$ for the solution u of the Neumann problem (1), (2) in any convex domain $\Omega \subset \mathbb{R}^n$, $n \geq 2$.*

A direct consequence of this fact is the sharp lower estimate $\Lambda \geq n - 1$ for the first nonzero eigenvalue Λ of the Neumann problem for the Beltrami operator on a convex subdomain of a unit hemisphere. It was obtained by a different argument for manifolds of positive Ricci curvature by J. F. Escobar in [6], where the case of equality was settled as well. This estimate answered a question raised by M. Dauge [5], and it leads, in combination with known techniques of the theory of elliptic boundary value problems in domains with piecewise smooth boundaries (see [5], [25]–[27]), to estimates for solutions of the problem (1), (2) in various function spaces. Two examples are given at the end of this article.

2 Main result

In what follows, we need a constant C_{Ω} in the relative isoperimetric inequality

$$s(\Omega \cap \partial g) \geq C_{\Omega} |g|^{1-1/n}, \quad (5)$$

where g is an arbitrary open set in Ω such that $|g| \leq |\Omega|/2$ and $\Omega \cap \partial g$ is a smooth (not necessarily compact) submanifold of Ω . By s we denote the $(n - 1)$ -dimensional area and by $|g|$ the n -dimensional Lebesgue measure. The Poincaré-Gagliardo-Nirenberg inequality

$$\inf_{t \in \mathbb{R}} \|v - t\|_{L^{n/(n-1)}(\Omega)} \leq \text{const.} \|\nabla v\|_{L^1(\Omega)}, \quad \forall v \in W^{1,1}(\Omega), \quad (6)$$

where $\text{const.} \leq C_{\Omega}^{-1}$ is a consequence of (5) (see Theorem 3.2.3 [24]).

Theorem. *Let $f \in L_{\perp}^q(\Omega)$ with a certain $q > n$. Then the solution $u \in W_{\perp}^{1,2}(\Omega)$ of the problem (1), (2) satisfies the estimate*

$$\|\nabla u\|_{L^{\infty}(\Omega)} \leq c(n, q) C_{\Omega}^{-1} |\Omega|^{(q-n)/qn} \|f\|_{L^q(\Omega)}, \quad (7)$$

where c is a constant depending only on n and q .

Proof. It suffices to prove (7) assuming that $f \in C_0^{\infty}(\Omega)$. Let us approximate Ω by a sequence $\{\Omega_m\}_{m \geq 1}$ of convex domains with smooth boundaries, $\Omega_m \supset \bar{\Omega}$. This can be done, for instance, by approximating Ω by a family of equidistant surfaces and by smoothing them with small perturbation of normal vectors. Then (6) implies

$$\inf_{t \in \mathbb{R}} \|w - t\|_{L^{\frac{n}{n-1}}(\Omega_m)} \leq (1 + \varepsilon) C_{\Omega}^{-1} \|\nabla w\|_{L^1(\Omega_m)}, \quad (8)$$

for all $w \in W^{1,1}(\Omega_m)$, where ε is an arbitrary positive number and $m = m(\varepsilon)$.

By u_m we denote a solution of the problem (1), (2) in Ω_m with f extended by zero outside Ω . One can easily see that $\nabla u_m \rightarrow \nabla u$ in $L^2(\Omega)$. Hence, it is enough to obtain (7) assuming that $\partial\Omega$ is smooth.

Let $t > \tau > 0$ and let Ψ be a piecewise linear continuous function on \mathbb{R} specified by $\Psi(\xi) = 0$ for $\xi < \tau$ and $\Psi(\xi) = 1$ for $\xi > t$. Note that

$$(\Delta u)^2 - |\nabla_2 u|^2 = (u_{x_j} \Delta u)_{x_j} - (u_{x_j} u_{x_i x_j})_{x_i}, \quad (9)$$

where

$$|\nabla_2 u| = \left(\sum_{1 \leq i, j \leq n} u_{x_i x_j}^2 \right)^{1/2}.$$

Hence

$$\begin{aligned} \int_{\Omega} \Psi(|\nabla u|) (f^2 - |\nabla_2 u|^2) dx &= \int_{\partial\Omega} \Psi(|\nabla u|) (\nu_j u_{x_j} \Delta u - \nu_i u_{x_j} u_{x_i x_j}) ds_x \\ &+ \int_{\Omega} \Psi'(|\nabla u|) ((|\nabla u|)_{x_j} u_{x_j} f + (|\nabla u|)_{x_i} u_{x_j} u_{x_i x_j}) dx, \end{aligned} \quad (10)$$

where (ν_1, \dots, ν_n) are components of the outward unit normal. By the Bernshtein-type identity (see, for instance, [11] or [17]), the first integral on the right-hand side of (10) equals

$$-2 \int_{\partial\Omega} Q(\nabla_{tan} u, \nabla_{tan} u) ds_x,$$

where Q is the second fundamental quadratic form on $\partial\Omega$ and ∇_{tan} is the tangential gradient. The form Q is nonpositive by convexity of Ω , which leads, together with (10), to the inequality

$$\int_{\Omega} \Psi'(|\nabla u|) ((|\nabla u|)_{x_j} u_{x_j} f + (|\nabla u|)_{x_i} u_{x_j} u_{x_i x_j}) dx \leq \int_{\Omega} \Psi(|\nabla u|) f^2 dx. \quad (11)$$

By the co-area formula [7], the left-hand side of (11) is identical to

$$(t - \tau)^{-1} \int_{\tau}^t \int_{|\nabla u|=\sigma} ((|\nabla u|)_{x_j} u_{x_j} f + (|\nabla u|)_{x_i} u_{x_j} u_{x_i x_j}) \frac{ds_x}{|\nabla|\nabla u||} d\sigma,$$

which is equal to

$$(t - \tau)^{-1} \int_{\tau}^t \int_{|\nabla u|=\sigma} \left(\frac{\partial u}{\partial \nu} f - |\nabla u| \frac{\partial}{\partial \nu} |\nabla u| \right) ds_x d\sigma \quad (12)$$

because $\nabla|\nabla u| = -\nu|\nabla|\nabla u||$ on the level surface $|\nabla u| = \sigma$, where ν is the unit normal, outward with respect to the set $\{x : |\nabla u| > \sigma\}$. The expression (12) can be written as

$$(\tau - t)^{-1} \int_{\tau}^t \int_{|\nabla u|=\sigma} \left(\frac{\partial u}{\partial \nu} f - |\nabla u| |\nabla|\nabla u|| \right) ds_x d\sigma.$$

Passing here to the limit as $\tau \uparrow t$ and using (11), we arrive at the estimate

$$\int_{|\nabla u|=t} (|\nabla u| |\nabla|\nabla u|| - f \frac{\partial u}{\partial \nu}) ds_x \leq \int_{|\nabla u|>t} f^2 dx,$$

which implies

$$t \int_{|\nabla u|=t} |\nabla|\nabla u|| ds_x \leq t \int_{|\nabla u|=t} |f| ds_x + \int_{|\nabla u|>t} f^2 dx. \quad (13)$$

We define the median of $|\nabla u|$ as

$$med |\nabla u| = \sup\{t \in \mathbb{R} : |\{|\nabla u| > t\}| \geq |\Omega|/2\},$$

and we note that

$$|\{|\nabla u| > med |\nabla u|\}| \leq |\Omega|/2$$

and

$$|\{|\nabla u| \geq \text{med } |\nabla u|\}| \geq |\Omega|/2.$$

Clearly,

$$\text{med } |\nabla u| \leq \left(\frac{2}{|\Omega|}\right)^{1/2} \|\nabla u\|_{L^2(\Omega)} \quad (14)$$

and, by Hölder's inequality and (8),

$$\begin{aligned} \|\nabla u\|_{L^2(\Omega)} &\leq \inf_{\gamma \in \mathbb{R}} \|u - \gamma\|_{L^{\frac{n}{n-1}}(\Omega)}^{1/2} \|f\|_{L^n(\Omega)}^{1/2} \\ &\leq (1 + \varepsilon) C_\Omega^{-1/2} \|\nabla u\|_{L^1(\Omega)}^{1/2} \|f\|_{L^n(\Omega)}^{1/2}. \end{aligned}$$

Hence,

$$\|\nabla u\|_{L^2(\Omega)} \leq (1 + \varepsilon)^2 C_\Omega^{-1} |\Omega|^{1/2} \|f\|_{L^n(\Omega)},$$

which, in combination with (14) and Hölder's inequality, implies

$$\text{med } |\nabla u| \leq 2^{1/2} (1 + \varepsilon)^2 C_\Omega^{-1} |\Omega|^{(q-n)/qn} \|f\|_{L^q(\Omega)}. \quad (15)$$

We introduce the function

$$\psi : [\text{med } |\nabla u|, \max |\nabla u|] \rightarrow [0, \infty)$$

by the equality

$$\psi(t) = \int_{\text{med } |\nabla u|}^t \left(\int_{|\nabla u|=\sigma} |\nabla |\nabla u|| ds_x \right)^{-1} d\sigma. \quad (16)$$

Let $E_\psi = \{x : |\nabla u(x)| = t(\psi)\}$ and $M_\psi = \{x : |\nabla u(x)| > t(\psi)\}$. Putting $t = t(\psi)$ in (13) and integrating in ψ over \mathbb{R}_+ , we arrive at

$$\max |\nabla u|^2 \leq (\text{med } |\nabla u|)^2 + 2 \max |\nabla u| \int_0^\infty \int_{E_\psi} |f| ds_x d\psi + 2 \int_0^\infty \int_{M_\psi} f^2 dx d\psi.$$

Recalling (15), we see that in order to obtain (7), it suffices to prove the inequalities

$$\int_0^\infty \int_{E_\psi} |f| ds_x d\psi \leq c C_\Omega^{-1} |\Omega|^{(q-n)/qn} \|f\|_{L^q(\Omega)} \quad (17)$$

and

$$\left(\int_0^\infty \int_{M_\psi} f^2 dx d\psi \right)^{1/2} \leq c C_\Omega^{-1} |\Omega|^{(q-n)/qn} \|f\|_{L^q(\Omega)}, \quad (18)$$

where $c = c(n, q)$.

The following argument leading to (17) is an obvious modification of the proof of Lemma 4 [22]. We start with the estimate for the area of the set E_ψ

$$s(E_\psi)^2 \leq -\frac{d}{d\psi} |M_\psi| \quad (19)$$

obtained in Lemma 2 [22]. By the triple Hölder inequality and (19)

$$\begin{aligned} \int_{E_\psi} |f| ds_x &\leq \left(\int_{E_\psi} |f|^q \frac{ds_x}{|\nabla |\nabla u||} \right)^{1/q} \left(\int_{E_\psi} |\nabla |\nabla u|| ds_x \right)^{1/q} s(E_\psi)^{1-2/q} \\ &\leq s(E_\psi)^{-1} \left(-\frac{d}{d\psi} |M_\psi| / t'(\psi) \right)^{1-1/q} \left(\int_{E_\psi} |f|^q \frac{ds_x}{|\nabla |\nabla u||} \right)^{1/q} \int_{E_\psi} |\nabla |\nabla u|| ds_x. \end{aligned}$$

This, in combination with the inequality

$$s(E_\psi) \geq (1 + \varepsilon)^{-1} C_\Omega |M_\psi|^{1-1/n}$$

(see (5)), implies the estimate for the integral on the left-hand side of (17)

$$\begin{aligned} \int_0^\infty \int_{E_\psi} |f| ds_x d\psi &\leq (1+\varepsilon) C_\Omega^{-1} \left(\int_0^{|\Omega|/2} \mu^{\frac{(1-n)q}{n(q-1)}} d\mu \right)^{1-1/q} \left(\int_0^\infty \int_{E_\psi} |f|^q \frac{ds_x}{|\nabla|\nabla u|} t'(\psi) d\psi \right)^{1/q} \\ &\leq (1+\varepsilon) C_\Omega^{-1} \left(\frac{n(q-1)}{q-n} \left(\frac{|\Omega|}{2} \right)^{\frac{q-n}{n(q-1)}} \right)^{1-1/q} \|f\|_{L^q(\Omega)} \end{aligned}$$

and the proof of (17) is complete.

We turn to inequality (18). By (19), its left-hand side does not exceed

$$\left(- \int_0^\infty \int_{M_\psi} |f|^2 dx \frac{d|M_\psi|}{s(E_\psi)^2} \right)^{1/2}$$

which is dominated by

$$\begin{aligned} &(1+\varepsilon) C_\Omega^{-1} \left(- \int_0^\infty \int_{M_\psi} |f|^2 dx \frac{d|M_\psi|}{|M_\psi|^{\frac{n-1}{n}}} \right)^{1/2} \\ &\leq (1+\varepsilon) C_\Omega^{-1} \left(\int_0^{|\Omega|/2} \int_0^\sigma f_*(\tau)^2 d\tau \frac{d\sigma}{\sigma^{\frac{n-1}{n}}} \right)^{1/2}, \end{aligned}$$

where f_* is the nonincreasing rearrangement of $|f|$. Now, (18) follows by integration by parts and Hölder's inequality. The proof of Theorem is complete.

3 Regularity of solutions to the Neumann problem in a convex polyhedron

The following assertion essentially stemming from the above theorem is a particular case of Escobar's result in [6] mentioned in Introduction.

Corollary. *Let ω be a convex subdomain of the upper unit hemisphere S_+^{n-1} . The first positive eigenvalue Λ of the Beltrami operator on ω with zero Neumann data on $\partial\omega$ is not less than $n-1$.*

Proof. Let $\lambda(\lambda+n-2) = \Lambda$ and $\lambda > 0$. In the convex domain

$$\Omega = \left\{ x \in \mathbb{R}^n : 0 < |x| < 1, \frac{x}{|x|} \in \omega \right\},$$

we define the function

$$u(x) = |x|^\lambda \Phi\left(\frac{x}{|x|}\right) \eta(|x|), \quad (20)$$

where Φ is an eigenfunction corresponding to Λ and η is a smooth cut-off function on $[0, \infty)$, equal to one on $[0, 1/2]$ and vanishing outside $[0, 1]$.

Let N be an integer satisfying $4N > n-1 \geq 4(N-1)$ and let $j = 0, 1, \dots, N$,

$$q_j = \begin{cases} \frac{2(n-1)}{n-1-4j}, & \text{if } 0 \leq j < (n-1)/4, \\ \text{arbitrary} & \text{if } j = (n-1)/4, \end{cases}$$

and $q_N = \infty$. Iterating the estimate

$$\|\Phi\|_{L^{q_{j+1}}(\omega)} \leq c \Lambda \|\Phi\|_{L^{q_j}(\omega)}$$

obtained in Theorems 5 and 6 [22], we see that $\Phi \in L^\infty(\omega)$.

The function u , defined by (20), satisfies the Neumann problem (1), (2) with

$$f(x) = -\Phi\left(\frac{x}{|x|}\right) [\Delta, \eta(|x|)] |x|^\lambda.$$

Since $\Phi \in L^\infty(\omega)$, it follows that $f \in L^\infty(\Omega)$ and by Theorem, $|\nabla u| \in L^\infty(\Omega)$, which is possible only if $\lambda \geq 1$, i.e. $\Lambda \geq n - 1$. The proof is complete.

Two applications of the above estimate for Λ will be formulated.

Let Ω be a convex bounded 3-dimensional polyhedron. By the techniques, well-known nowadays (see [5], [25]–[27]), one can show the unique solvability of the variational Neumann problem in $W_\perp^{1,p}(\Omega)$ for every $p \in (1, \infty)$. By definition of this problem, its solution is subject to the integral identity

$$\int_\Omega \nabla u \cdot \nabla \eta \, dx = f(\eta),$$

where f is a given distribution in the space $(W^{1,p'}(\Omega))^*$, $f(1) = 0$ and η is an arbitrary function in $W^{1,p'}(\Omega)$, $p + p' = pp'$.

Let us turn to the second application of Corollary. We continue to deal with the polyhedron Ω in \mathbb{R}^3 . Let $\{\mathcal{O}\}$ be the collection of all vertices and let $\{U_{\mathcal{O}}\}$ be an open finite covering of $\overline{\Omega}$ such that \mathcal{O} is the only vertex in $U_{\mathcal{O}}$. Let also $\{E\}$ be the collection of all edges and let α_E denote the opening of the dihedral angle with edge E , $0 < \alpha_E < \pi$. The notation $r_E(x)$ stands for the distance between $x \in U_{\mathcal{O}}$ and the edge E such that $\mathcal{O} \in \overline{E}$.

With every vertex \mathcal{O} and edge E we associate real numbers $\beta_{\mathcal{O}}$ and δ_E , and we introduce the weighted L^p -norm

$$\|v\|_{L^p(\Omega; \{\beta_{\mathcal{O}}\}, \{\delta_E\})} := \left(\sum_{\{\mathcal{O}\}} \int_{U_{\mathcal{O}}} |x - \mathcal{O}|^{p\beta_{\mathcal{O}}} \prod_{\{E: \mathcal{O} \in \overline{E}\}} \left(\frac{r_E(x)}{|x - \mathcal{O}|} \right)^{p\delta_E} |v(x)|^p \, dx \right)^{1/p},$$

where $1 < p < \infty$. Under the conditions

$$\begin{aligned} 3/p' &> \beta_{\mathcal{O}} > -2 + 3/p', \\ 2/p' &> \delta_E > -\min\{2, \pi/\alpha_E\} + 2/p', \end{aligned}$$

the inclusion of the function f in $L^p(\Omega; \{\beta_{\mathcal{O}}\}, \{\delta_E\})$ implies the unique solvability of the problem (1), (2) in the class of functions with all derivatives of the second order belonging to $L^p(\Omega; \{\beta_{\mathcal{O}}\}, \{\delta_E\})$. This fact follows from Corollary and a result essentially established in Sect. 7.5 [27].

An important particular case when all $\beta_{\mathcal{O}}$ and δ_E vanish, i.e. when we deal with a standard Sobolev space $W^{2,p}(\Omega)$, is also included here. To be more precise, if

$$1 < p < \min\left\{3, \frac{2\alpha_E}{(2\alpha_E - \pi)_+}\right\} \quad (21)$$

for all edges E , then the inverse operator of the problem (1), (2):

$$L_\perp^p(\Omega) \ni f \rightarrow u \in W_\perp^{2,p}(\Omega)$$

is continuous whatever the convex polyhedron $\Omega \subset \mathbb{R}^3$ may be. The bounds for p in (21) are sharp for the class of all convex polyhedra.

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