The Integrability of Negative Powers of the Solution of the Saint Venant Problem *

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Abstract

We initiate the study the finiteness condition $\int_\Omega u(x)^{-\beta}\, dx \leq C(\Omega,\beta) < +\infty$ where $\Omega \subseteq \mathbb{R}^n$ is an open set and u is the solution of the Saint Venant problem $\Delta u = -1$ in Ω, $u = 0$ on $\partial\Omega$. The central issue which we address is that of determining the range of values of the parameter $\beta > 0$ for which the aforementioned condition holds under various hypotheses on the smoothness of Ω and demands on the nature of the constant $C(\Omega,\beta)$. Classes of domains for which our analysis applies include bounded piecewise C^1 domains in $\mathbb{R}^n$, $n \geq 2$, with conical singularities (in particular polygonal domains in the plane), polyhedra in $\mathbb{R}^3$, and bounded domains which are locally of class C^2 and which have (finitely many) outwardly pointing cusps. One of the original motivations for addressing the aforementioned issue was the study of sublevel set estimates for functions v satisfying $v(0) = 0$, $\nabla v(0) = 0$ and $\Delta v \geq c > 0$.

1 Introduction

1.1 Background

Suppose that u is a positive superharmonic function defined in an open, bounded subset Ω of $\mathbb{R}^n$, i.e.

$$\Delta u \leq 0 \quad \text{and} \quad u > 0 \quad \text{in} \quad \Omega. \tag{1.1}$$

Two issues which have received a considerable amount of attention in the literature are:

(i) proving lower pointwise bounds for u in terms of powers of the distance function to the boundary, and

(ii) establishing the membership of u to the Lebesgue scale $L^p(\Omega)$, $0 < p \leq \infty$.

See, for example, [2], [3], [5], [6], [23], [24], [25], [27], [30] and the references therein. We wish to highlight two aspects of the philosophy that has emerged from these studies. First, granted a certain degree of reasonableness of the underlying domain, for superharmonic functions, positivity always entails a quantitative version of itself, in the form of the estimate

$$u(x) \geq C(\Omega, u)\, \delta_\Omega(x)^\alpha, \qquad \text{for all} \ \ x \in \Omega, \tag{1.2}$$

where $C(\Omega, u) > 0$ is a constant depending on u and Ω, for some exponent $\alpha = \alpha(\Omega) \geq 1$ independent of u. Here and elsewhere, for an arbitrary set $\Omega \subseteq \mathbb{R}^n$, we have denoted by δ_Ω the (Euclidean) distance to its boundary, i.e.,

$$\delta_\Omega(x) := \text{dist}\,(x, \partial\Omega), \qquad \forall\, x \in \mathbb{R}^n. \tag{1.3}$$

The second aspect alluded to above is that there is a common integrability threshold for the entire class positive superharmonic functions in the sense that

$$\int_\Omega u(x)^p \, dx < +\infty, \tag{1.4}$$

for some integrability exponent $p = p(\Omega) > 0$ independent of the positive superharmonic function u in Ω.

The specific nature of the exponents $\alpha(\Omega)$ and $p(\Omega)$ is dictated by the degree of regularity exhibited by Ω. For example, (1.2) has been proved for $\alpha = 1$ in a suitable subclass of the class of domains satisfying a uniform interior ball condition which, in turn, contains the class of bounded C^2 domains, by Kuran in [23], and for bounded planar Jordan domains with a Dini-continuous boundary by Kuran and Schiff in [24]. On the other hand, the lower bounds for the Green function established in [25] for Lipschitz domains also lead to estimates of the type (1.2), typically for exponents larger than one.

As far as (1.4) is concerned, in the case when $\Omega \subseteq \mathbb{R}^n$ is a bounded C^∞ domain, Armitage [2], [3] has proved that (1.4) holds for any positive superharmonic function u in Ω, granted that $0 < p < n/(n-1)$. This result has been subsequently extended by Maeda and Suzuki in [25] to the class of bounded Lipschitz domains for a range of p's which depends on the Lipschitz constant of the domain in question, in such a way that $p \nearrow n/(n-1)$ as the domain is progressively closer and closer to being of class C^1 (i.e., as the Lipschitz constant approaches zero). Further refinements of this result, in the class of John domains and Hölder domains (in which scenario p is typically small), have been studied, respectively by Aikawa in [5] and by Stegenga and Ullrich in [30].

1.2 Overview and motivation

In this paper we are concerned with the validity of (1.4) for *negative* values of the integrability exponents, in the case when u is a positive function with $\Delta u < 0$ in Ω. A case in point is the solution of the Saint Venant problem[1] (cf., e.g., [7], [12], [13])

$$\begin{cases} \Delta u = -1 & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega, \end{cases} \tag{1.5}$$

and the question which makes the object of our study is that of determining the range of values of the parameter $\beta > 0$ for which an estimate of the form

$$\int_\Omega u(x)^{-\beta} \, dx \leq C(\Omega, \beta) < +\infty \tag{1.6}$$

holds, under various conditions on Ω and demands on the nature of the constant $C(\Omega, \beta)$. Cases of special interest include the class of nontangentially accessible domains satisfying an inner cone condition (which includes the class of Lipschitz domain) in $\mathbb{R}^n$, polygonal domains in $\mathbb{R}^2$, polyhedral domains in $\mathbb{R}^3$, as well as piecewise smooth domains with conical and cuspidal singularities. Since $\int_\Omega u(x)^{-\beta} \, dx$ is entirely determined by the domain Ω and the parameter β, we shall occasionally refer to this number as the "β-integral of Ω."

Aside from its relevance in potential theory, the problem (1.5) plays a significant role in elasticity theory. For example, the torsional rigidity coefficient of Ω, originally defined as

$$P(\Omega) := \sup_{0 \neq w \in C_0^\infty(\Omega)} \left(\int_\Omega |w| \, dx \right)^2 \left(\int_\Omega |\nabla w|^2 \, dx \right)^{-1} \tag{1.7}$$

turns out to be

$$P(\Omega) = \int_\Omega u \, dx = \int_\Omega |\nabla u|^2 \, dx, \tag{1.8}$$

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[1]Much of our analysis will apply in the more general case that $\Delta u \leq -C_n$ in Ω, where C_n is a positive dimensional constant.

where u is the solution of (1.5) (cf. the discussion in [7], [11], [19]).

Our interest in the estimate (1.6) was originally motivated by problems in harmonic analysis concerning sublevel set estimates for a real-valued, strictly convex function of class C^2 defined in an open, convex set $\Omega \subseteq \mathbb{R}^n$. (It is thus also related to the behaviour of oscillatory integrals; cf. [8], [9], [32].) It this vein, we recall that it has been shown in [8] that there exists a finite dimensional constant $C = C_n > 0$ with the property that, with $|E|$ denoting the Lebesgue measure[2] of a Lebesgue measurable set E,

$$|\Omega| \leq C \|v\|_{L^\infty(\Omega)}^{n/2}, \tag{1.9}$$

provided that, in addition to the already mentioned properties, the Hessian of the function v satisfies

$$\det\Big[\Big(\frac{\partial^2 v}{\partial x_i \partial x_j}\Big)_{1 \leq i,j \leq n}\Big] \geq 1 \quad \text{on} \ \ \Omega. \tag{1.10}$$

As noted in [8], if v is also nonnegative, then by applying (1.9) with $\{x \in \Omega : v(x) < t\}$, $t > 0$, in place of Ω we obtain the sublevel set estimate

$$\big|\{x \in \Omega : v(x) < t\}\big| \leq C\, t^{n/2}, \qquad t > 0. \tag{1.11}$$

On the other hand, granted (1.10), the arithmetic-geometric mean inequality gives

$$\Delta v/n \geq \big(\det((\partial_{ij} v)_{1 \leq i,j \leq n})\big)^{1/n} \geq 1. \tag{1.12}$$

Hence, it is natural to ask, what happens with (1.9) if we only knew $\Delta v \geq n$? Is it reasonable to expect to still have such an estimate which, by the same procedure as above, would then lead to a sub-level set estimate similar to (1.11)? If so, what is the nature of the constant C in (1.11) in this more general situation?

We wish to elaborate on this point and, in particular, make it more transparent how condition (1.6) for the solution of (1.5) comes into play. To set the stage, assume that v is a real-valued, strictly convex function v of class C^2 in a neighbourhood of the origin in $\mathbb{R}^n$ and which is normalised so that

$$v(0) = 0, \qquad \nabla v(0) = 0. \tag{1.13}$$

Next, fix a (small) threshold $t > 0$, define

$$\Omega := \{x : v(x) < t\} \subseteq \mathbb{R}^n, \tag{1.14}$$

and, from now on, restrict v to the open convex set Ω. To continue, denote by $\mathcal{G}$ the region of the space in $\mathbb{R}^{n+1}$ lying directly above the graph of the function v and below the the n-dimensional horizontal plane $x_{n+1} = t$, i.e.,

$$\mathcal{G} := \{(x, x_{n+1}) \in \mathbb{R}^n \times \mathbb{R} : x \in \Omega \text{ and } v(x) < x_{n+1} < t\}. \tag{1.15}$$

In order to estimate $|\mathcal{G}|$, the $(n+1)$-dimensional Lebesgue measure of $\mathcal{G}$, let u solve the auxiliary problem (1.5). We then have

$$|\mathcal{G}| = \int_\Omega (t - v(x))\, dx = \int_\Omega (v(x) - t)(\Delta u)(x)\, dx = \int_\Omega (\Delta v)(x) u(x)\, dx, \tag{1.16}$$

after integrating by parts and using the fact that both u and $v - t$ vanish on $\partial\Omega$. Using this formula, for given $\gamma \in (0, 1/2)$ we may then compute (making use of the obvious inequality $|\mathcal{G}| \leq t\,|\Omega|$)

$$\begin{aligned}
\int_\Omega (\Delta v)^\gamma \, dx &= \int_\Omega ((\Delta v) u)^\gamma u^{-\gamma}\, dx \leq \Big(\int_\Omega (\Delta v) u\, dx\Big)^\gamma \Big(\int_\Omega u^{-\gamma/(1-\gamma)}\, dx\Big)^{1-\gamma} \\
&= |\mathcal{G}|^\gamma \Big(\int_\Omega u^{-\gamma/(1-\gamma)}\, dx\Big)^{1-\gamma} \leq t^\gamma |\Omega|^\gamma \Big(\int_\Omega u^{-\gamma/(1-\gamma)}\, dx\Big)^{1-\gamma},
\end{aligned} \tag{1.17}$$

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[2]Later on, we shall also occasionally use the notation $\mathcal{L}^n(E)$ in place of $|E|$.

hence

$$\int_\Omega (\Delta v)^\gamma \, dx \le \|v\|_{L^\infty(\Omega)}^\gamma |\Omega|^\gamma \Big(\int_\Omega u^{-\beta} \, dx \Big)^{1-\gamma}, \tag{1.18}$$

where have set $\beta := \frac{\gamma}{1-\gamma} \in (0,1)$. Note that in the case in which

$$\int_\Omega u(x)^{-\beta} \, dx \le C_\beta \, |\Omega|^{1-2\beta/n}, \tag{1.19}$$

this analysis gives

$$\int_\Omega (\Delta v)^\gamma \le C_\gamma \|v\|_{L^\infty(\Omega)}^\gamma \, |\Omega|^{1-\frac{2\gamma}{n}}. \tag{1.20}$$

The upshot of this analysis is that by using the weaker condition $\Delta v \ge n$ (in place of the quantitative non-degeneracy of the Hessian matrix for v, as in (1.10)), one deduces from (1.20) that

$$n^\gamma |\Omega| \le C_\gamma \, t^\gamma |\Omega|^{1-2\gamma/n}, \tag{1.21}$$

which leads to

$$|\{x : \ v(x) < t\}| \le D \, t^{n/2}, \qquad t > 0, \tag{1.22}$$

with D depending only on the dimension n. This is of course contingent upon (1.20) holding for some $\gamma \in (0, 1/2)$ (possibly depending on n) with the constant C_γ (which is related to C_β from (1.19) via $C_\gamma = (C_\beta)^{1-\gamma}$) being independent of the parameter t.

However, we cannot expect an inequality such as (1.22) to hold for an arbitrary strictly convex v, defined on a convex domain containing 0, which satisfies (1.13). For example consider, for small ϵ, the function $v_\epsilon(x) = x_1^2 + \epsilon x_2^2$ defined on $\mathbb{R}^2$, for which (1.22) is easily seen to fail. Upon reflection, this is related to the fact that inequality (1.20), considered for arbitrary convex domains Ω and strictly convex v defined on Ω, is dilation invariant, but, unlike its counterpart for the Hessian problem, is not affine invariant. Thus we cannot expect inequality (1.19) to hold uniformly over all convex domains Ω, and indeed at the end of Section 3 we demonstrate this explicitly. On the other hand, Proposition 2.13 below shows that for convex sets Ω containing the unit ball and contained in some dimensional multiple of the unit ball, (1.19) does hold for $\beta < 1/2$. We have not been able to discount the possibility that for such sets Ω, (1.19) might even hold for all $\beta < 1$ with a constant depending only on β and n. For some evidence in that direction, in the two-dimensional case at least, see Theorem 4.7 and Proposition 4.9 below.

1.3   Description of results and layout of the paper

The discussion in §1.1-§1.2 highlights the significance of the problem (1.5) as well as the relevance of the finiteness condition (1.6). Note that the solution u of (1.5) satisfies $1/u \in L^\infty_{loc}(\Omega)$, so the finiteness condition in (1.6) is related to the rate at which u vanishes on the boundary. While, from this point of view, a pointwise lower bound such as (1.2) provides, in principle, a venue for deducing an estimate of the form (1.6), the range of negative integrability exponents obtained by such a method is typically far from optimal, so a number of new ideas are required. A succinct summary of our main results is as follows:

Theorem 1.1. *The β-integral associated with a bounded domain $\Omega \subseteq \mathbb{R}^n$ is finite in any of the following situations:*

(i) $\beta \in (0,1)$ and Ω is a bounded piecewise C^1 domain in $\mathbb{R}^n$, $n \ge 2$, with conical singularities;

(ii) $\beta \in (0,1)$ and Ω is a polyhedron in $\mathbb{R}^3$;

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(iii) $\beta \in (0,1)$ *and* $\Omega \subseteq \mathbb{R}^n$, $n \geq 2$, *is a bounded domain, locally of class C^2 and which has an outwardly pointing cusp at $0 \in \partial\Omega$. Specifically, it is assumed that there exists a small number $\varepsilon > 0$ and a function $\mathcal{F} \in C^2([0,1])$ with $\mathcal{F} > 0$ on $(0,1]$ and $(\mathcal{F}^2)''(0) = 0$, for which $\{x \in \Omega : x_n \leq 1\}$ coincides with the cuspidal set $\{x = (x', x_n) : 0 < x_n \leq 1, |x'| < \varepsilon\mathcal{F}(x_n)\}$. In the case when $n = 2$ and $\beta \in (1/2, 1)$, the following (necessary) finiteness condition is also assumed:*

$$\int_0^1 \mathcal{F}(\tau)^{1-2\beta}\, d\tau < +\infty. \tag{1.23}$$

Of course, in part (iii) of Theorem 1.1, the same type of result holds for any bounded piecewise C^2 domain with (finitely many) exterior cusps. We also wish to emphasise that part (i) of Theorem 1.1 covers, in particular, the case of polygons in the plane. One special case is treated in Proposition 4.9 where it is shown that, for each fixed $\beta \in (0,1)$, the β-integral of a regular polygon with N sides which is circumscribed by the unit disc has following asymptotic

$$\frac{4^\beta \pi}{1-\beta} + \mathcal{O}(N^{\beta-1}) \qquad \text{as} \;\; N \to \infty. \tag{1.24}$$

In the case of a bounded piecewise C^1 domain Ω with conical singularities, our approach is to estimate the contribution from individual conical points by carefully devising appropriate barrier functions which compare favourably with the solution of (1.5). The contribution from the region Ω away from the boundary singularities is then estimated separately, by relying on results valid on smooth domains. See Theorem 4.3 and Theorem 4.7 which are the main results phrased in the two-dimensional setting, as well as Theorem 5.2 which contains an extension to the higher dimensional case.

Let us now review the content of the various sections of this paper. In Section 2 we derive estimates for the solution of the Saint Venant problem in rather general domains, satisfying weak regularity properties, described in terms of basic geometric measure theoretic conditions. This portion of our analysis points to the value $\beta = 1/2$ as the natural critical exponent for the condition (1.6) in this degree of generality for the underlying domain Ω. Improvements of this result in the case when Ω satisfies an inner cone condition (cf. Definition 3.4) are subsequently discussed in Section 3. As a preamble, here we briefly review the construction and properties of classical barrier functions in cones. We then derive a lower pointwise bound for the Green function (akin to work in [25]) which is then used to prove Theorem 3.7, the main result in this section. A consequence of this theorem is that (1.6) holds for any $\beta \in (0,1)$ in the case when Ω is a bounded C^1 domain.

Sections 4-5 are devoted to studying the class of bounded piecewise C^1 domains with conical singularities, and the main results here are Theorem 4.3 and Theorem 5.2. The new phenomenon which we discover is that, much as for bounded C^1 domains, (1.6) continues to hold for every $\beta \in (0,1)$ in the aforementioned class of piecewise C^1 domains with conical singularities in $\mathbb{R}^n$, even though direct pointwise estimates (from below) on the solution u of (1.5) indicate that the regularity of u at a given conical singularity deteriorates as the aperture of the cone tends to zero. Hence, informally speaking, for the type of domains considered in these sections, the size (smallness, in fact) of the solution of the Saint Venant problem (1.5) is, in the process of taking an integral average, better controlled than pointwise estimates from below in terms of powers of the distance function to the boundary might originally seem to indicate. Finally, in Section 6, we study the veracity of (1.6) for other classes of domains with isolated singularities, such as polyhedra and piecewise C^1 domains with outwardly pointing cuspidal singularities.

Throughout, we employ the customary convention of using the same letter for denoting constants whose values may change from line to line. Whenever the dependence of the constants in question on certain parameters is important, we indicate this as such. Also, $F \approx G$ means that there exist $C_1, C_2 > 0$ which are independent of the relevant parameters entering the expressions F, G with the property that $C_1 F \leq G \leq C_2 F$.

2   Estimates for the Saint Venant problem in rough domains

Let Ω be a bounded, open subset in $\mathbb{R}^n$, and denote by $W^{1,p}(\Omega)$ the classical L^p-based Sobolev space of order one in Ω, where $1 \le p \le \infty$. Furthermore, we shall use $W_0^{1,p}(\Omega)$ to denote the closure of $C_0^\infty(\Omega)$ in $W^{1,p}(\Omega)$. A standard application of the Lax-Milgram lemma shows that the Saint Venant problem (1.5) has a unique solution in the energy space $W_0^{1,2}(\Omega)$, i.e.,

$$u \in W_0^{1,2}(\Omega), \qquad \Delta u = -1 \ \text{ in } \ \Omega, \tag{2.1}$$

is always well-posed. In fact, the solution u of (2.1) can be expressed as

$$u(x) = \int_\Omega G(x, y)\, dy, \qquad x \in \Omega, \tag{2.2}$$

where $G(\cdot, \cdot)$ is the Green function for the Dirichlet Laplacian in Ω. The latter is the unique function $G : \Omega \times \Omega \to [0, +\infty]$ satisfying

$$G(\cdot, y) \in W^{1,2}(\Omega \setminus B(y, r)) \cap W_0^{1,1}(\Omega), \qquad \forall\, y \in \Omega, \ \forall\, r > 0, \tag{2.3}$$

and

$$\int_\Omega \langle \nabla_x G(x, y), \nabla \varphi(x) \rangle\, dx = \varphi(y), \quad \forall\, \varphi \in C_0^\infty(\Omega). \tag{2.4}$$

See, e.g., [16] and [20] for the proof of the existence and uniqueness of the Green function; a number of other useful properties of the Green function can be found in these works, such as the fact that the Green function is symmetric (i.e., $G(x, y) = G(y, x)$ for all $x, y \in \Omega$) and satisfies the estimates (valid for $n \ge 3$)

$$G(x, y) \le C_n |x - y|^{2-n} \quad \text{ for all } x, y \in \Omega, \tag{2.5}$$

$$G(x, y) \ge C_n |x - y|^{2-n} \quad \text{ for } x, y \in \Omega \text{ with } |x - y| \le \tfrac{1}{2}\, \delta_\Omega(x), \tag{2.6}$$

where the constants depend only on the dimension. The replacement for $|x - y|^{2-n}$ in the case when $n = 2$ is $\log(\mathrm{diam}\,(\Omega)/|x - y|)$. Hence, as a consequence of (2.2) and (2.5),

$$0 < u(x) \le C_n\, [\mathrm{diam}\,(\Omega)]^2, \quad \text{ for each } \ x \in \Omega. \tag{2.7}$$

Remark 2.1. *Let $\Omega \subseteq \mathbb{R}^n$ be an arbitrary open set. Then for every $\beta > 0$ the solution of (2.1) satisfies the bound from below*

$$C(n, \beta)|\Omega|\, [\mathrm{diam}\,(\Omega)]^{-2\beta} \le \int_\Omega u(x)^{-\beta}\, dx. \tag{2.8}$$

Indeed, $|\Omega| = \int_\Omega u(x)^\beta u(x)^{-\beta}\, dx \le C_n^\beta [\mathrm{diam}\,(\Omega)]^{2\beta} \int_\Omega u(x)^{-\beta}\, dx$, by (2.7). In particular, if Ω has the property that $B(0, 1) \subseteq \Omega \subseteq B(0, C_n)$, then for every $\beta > 0$ there holds

$$C(n, \beta) \le \int_\Omega u(x)^{-\beta}\, dx. \tag{2.9}$$

We wish to point out that in the case when Ω is regular for the Dirichlet problem (i.e., the classical Dirichlet problem is well-posed in the class of continuous functions), one actually has $u \in C^0(\overline{\Omega})$. Necessary and sufficient criteria for regularity are well-known. For example, any bounded open set $\Omega \subseteq \mathbb{R}^n$ is regular for the Dirichlet problem if it satisfies an exterior corkscrew condition. The latter piece of terminology is clarified in the definition below.

Definition 2.2. *We say that $\Omega \subset \mathbb{R}^n$ satisfies an interior corkscrew condition if there are constants $M > 1$ and $R > 0$ such that for each $x \in \partial\Omega$ and $r \in (0, r_*)$ there exists*

$$A_r(x) \in \Omega, \text{ called corkscrew point relative to } x,$$
$$\text{so that } |x - A_r(x)| < r \text{ and } \operatorname{dist}(A_r(y), \partial\Omega) > M^{-1} r. \tag{2.10}$$

Also, $\Omega \subset \mathbb{R}^n$ satisfies the an exterior corkscrew condition if $\Omega^c := \mathbb{R}^n \setminus \Omega$ satisfies an interior corkscrew condition.

As explained in §1, the central issue in this paper is that of determining the "largest" value of the parameter $\beta > 0$ for which an estimate of the form (1.6) holds, under various geometrical conditions on Ω. Elucidating the nature of the constant $C(\Omega, \beta)$ appearing in (1.6) is also of interest. A basic tool systematically employed throughout the paper is the Maximum Principle. In order to state a version of this result valid for functions in the Sobolev space $W^{1,2}(\Omega)$ we first recall the following definition.

Definition 2.3. *Let $\Omega \subseteq \mathbb{R}^n$ be a bounded open set and assume that $E \subseteq \overline{\Omega}$. Given $u \in W^{1,2}(\Omega)$, we say that $u \geq 0$ on E in the sense of $W^{1,2}(\Omega)$ if there exists a sequence $u_j \in C^\infty(\Omega) \cap W^{1,2}(\Omega)$, $j \in \mathbb{N}$, which converges to u in $W^{1,2}(\Omega)$ and such that, for each $j \in \mathbb{N}$, there exists an open neighbourhood U_j of E in $\mathbb{R}^n$ with the property that $u_j > 0$ in $U_j \cap \Omega$.*

As is well-known, if $E \subseteq \Omega$ and $u \geq 0$ on E in the sense of $W^{1,2}(\Omega)$ then $u \geq 0$ a.e. on E. Furthermore, if $u \in W^{1,2}(\Omega)$ satisfies $u \geq 0$ a.e. in Ω then $u \geq 0$ in Ω in the sense of $W^{1,2}(\Omega)$. Let us also point out here that if $u \in W^{1,2}(\Omega) \cap C^0(\overline{\Omega})$ satisfies $u|_{\partial\Omega} \geq 0$ then $u \geq 0$ on $\partial\Omega$ in the sense of $W^{1,2}(\Omega)$ (cf. [20]).

Analogously to Definition 2.3, one can define $u \leq 0$ and $u = 0$ on $E \subseteq \overline{\Omega}$ in the sense of $W^{1,2}(\Omega)$. In particular, this allows one to compare any two functions $u, v \in W^{1,2}(\Omega)$ on $E \subseteq \overline{\Omega}$ in the sense of $W^{1,2}(\Omega)$, and also to define the supremum and infimum of a function $u \in W^{1,2}(\Omega)$ on $E \subseteq \overline{\Omega}$ in the sense of $W^{1,2}(\Omega)$. In this context, the following version of the Maximum Principle then holds (cf. [20, Lemma 1.1.17]):

Proposition 2.4. *Let $\Omega \subseteq \mathbb{R}^n$ be a bounded open set and assume that $u \in W^{1,2}(\Omega)$ is a subharmonic function in Ω. Then*

$$\sup_{\Omega} u \leq \sup_{\partial\Omega} u \quad \text{in the sense of } W^{1,2}(\Omega). \tag{2.11}$$

Returning to the main topic of interest for us here, we continue by making a series of simple yet significant remarks.

Remark 2.5. The case of a ball in $\mathbb{R}^n$, i.e., when $\Omega = B(0, R)$, $R > 0$, in which scenario (2.1) has the explicit solution

$$u(x) = \frac{1}{2n}(R^2 - |x|^2), \qquad x \in B(0, R), \tag{2.12}$$

shows that we must necessarily have $\beta < 1$ and that the critical value $\beta = 1$ is unattainable. Indeed, the function in (2.12) satisfies

$$\frac{R}{2n} \, \delta_{B(0,R)}(x) \leq u(x) \leq \frac{R}{n} \, \delta_{B(0,R)}(x), \qquad \forall \, x \in B(0, R), \tag{2.13}$$

Hence, in this case,

$$\int_{B(0,R)} u(x)^{-1} \, dx = +\infty. \tag{2.14}$$

In fact, it can be seen that this is typical of any sufficiently smooth domain (in fact, Theorem 3.7, stated later, shows that any domain of class C^1 will do), namely any $\beta < 1$ will work in (1.6).

Remark 2.6. Regarding the issue whether the β-integral diverges when $\beta = 1$, we shall show that this is always the case when the underlying domain satisfies the following condition: *Given $\Omega \subseteq \mathbb{R}^n$ and $x_* \in \partial\Omega$, we say that Ω satisfies an enveloping ball condition of radius $R > 0$ near x_* if there exists $\rho > 0$ with the property that for every $x \in B(x_*, \rho) \cap \partial\Omega$ there exists a ball of radius R which contains Ω and whose boundary contains x.*

The relevance of this piece of terminology is apparent from the following result: *Let $\Omega \subseteq \mathbb{R}^n$ be a bounded open set which satisfies an enveloping ball condition of radius $R > 0$ near a point $x_* \in \partial\Omega$. Then, if u denotes the solution of the Saint Venant boundary value problem* (2.1),

$$u(x) \leq n^{-1}R\,\delta_\Omega(x) \quad \text{for every } x \in \Omega \text{ near } x_*, \tag{2.15}$$

so that, in particular,

$$\int_\Omega u(x)^{-1}\,dx = +\infty. \tag{2.16}$$

To prove the above bound on u in terms of the distance to the boundary, consider an arbitrary point $x_0 \in B(x_*, \rho/2) \cap \Omega$ and denote by $x_1 \in \partial\Omega$ a point for which $r := \delta_\Omega(x_0) = |x_1 - x_0|$. Then, necessarily, $x_1 \in B(x_*, \rho) \cap \partial\Omega$. Consider now a ball $B = B(x_2, R)$ which contains Ω and such that $x_1 \in \partial B$. Note that $B(x_0, r) \subseteq \Omega \subseteq B$, so that the balls $B(x_0, r)$ and $B(x_2, R)$ are tangent at x_1. This implies that the points x_1, x_0, x_2 are collinear hence, further, $R - r = |x_0 - x_2|$. Next, use the Maximum Principle to deduce that

$$u(x) \leq (2n)^{-1}(R^2 - |x - x_2|^2) \quad \text{for every } x \in \Omega, \tag{2.17}$$

which, when specialised to $x = x_0$, gives

$$u(x_0) \leq n^{-1}R(R - |x_0 - x_2|) = n^{-1}Rr = n^{-1}R\,\delta_\Omega(x_0). \tag{2.18}$$

Since $x_0 \in B(x_*, \rho/2) \cap \Omega$ was arbitrary, (2.15) follows. As far as (2.16) is concerned, we first note that $B(x_*, \rho/2) \cap \Omega$ is convex, hence Lipschitz (a formal definition is given later, in (3.3)). In turn, this and (2.15) give that

$$\int_\Omega u(x)^{-1}\,dx \geq \int_{B(x_*, \rho/2)\cap\Omega} u(x)^{-1}\,dx \geq nR^{-1} \int_{B(x_*, \rho/2)\cap\Omega} \delta_\Omega(x)^{-1}\,dx = +\infty, \tag{2.19}$$

where the last step is a simple consequence of the fact that $B(x_*, \rho/2) \cap \Omega$ is Lipschitz (a more general result of this nature is discussed later, in Remark 2.10).

A more refined analysis proves Proposition 2.13, stated below. As a preamble, we first recall some definitions, as well as several results of independent interest.

Definition 2.7. *Given an open set $\Omega \subseteq \mathbb{R}^n$, the upper and lower γ-dimensional Minkowski contents of $\partial\Omega$ with respect to Ω are defined as*

$$M^*_\gamma(\partial\Omega) := \limsup_{r\to 0^+} \omega_{n-\gamma}(r), \qquad M_{*,\gamma}(\partial\Omega) := \liminf_{r\to 0^+} \omega_{n-\gamma}(r), \tag{2.20}$$

where, for every $\alpha \in \mathbb{R}$, we have set

$$\omega_\alpha(r) := \frac{|\{x \in \Omega : \delta_\Omega(x) < r\}|}{r^\alpha}. \tag{2.21}$$

The upper and lower Minkowski dimensions of $\partial\Omega$ with respect to Ω are then given by

$$\dim^*_{Minkowski}(\partial\Omega) := \inf\{\gamma \geq 0 : M^*_\gamma(\partial\Omega) < +\infty\} = \sup\{\gamma \geq 0 : M^*_\gamma(\partial\Omega) = +\infty\}$$

$$= \inf\{\gamma \geq 0 : M^*_\gamma(\partial\Omega) = 0\} = \sup\{\gamma \geq 0 : M^*_\gamma(\partial\Omega) > 0\}, \tag{2.22}$$

$$\dim_{*,Minkowski}(\partial\Omega) := \sup\{\gamma \geq 0 : M_{*,\gamma}(\partial\Omega) > 0\} = \inf\{\gamma \geq 0 : M_{*,\gamma}(\partial\Omega) = 0\}$$

$$= \inf\{\gamma \geq 0 : M_{*,\gamma}(\partial\Omega) < +\infty\} = \sup\{\gamma \geq 0 : M_{*,\gamma}(\partial\Omega) = +\infty\}, \tag{2.23}$$

with the convention that $\inf \emptyset := +\infty$ and $\sup \emptyset := -\infty$. When $\dim^*_{Minkowski}(\partial\Omega) = \dim_{*,Minkowski}(\partial\Omega)$, the common value is referred to as the Minkowski dimension of $\partial\Omega$ with respect to Ω, and is denoted by $\dim_{Minkowski}(\partial\Omega)$.

Next, we recall the Coarea Formula (see, e.g., [15], [14, Theorem 2, p. 117]). Given a fixed number $n \in \mathbb{N}$, denote by $\mathcal{L}^n$ the n-dimensional Lebesgue measure in $\mathbb{R}^n$ (occasionally we shall use the notation $\mathcal{L}^n(E)$ in place of $|E|$) and, for each $k \in \mathbb{N}$, $k \leq n$, let $\mathcal{H}^k$ stand for the k-dimensional Hausdorff measure in $\mathbb{R}^n$.

Proposition 2.8. *Assume that $n \geq m$ and that $f : \mathbb{R}^n \to \mathbb{R}^m$ is a given Lipschitz function. Then, for any $A \subseteq \mathbb{R}^n$ which is $\mathcal{L}^n$-measurable and $g \in L^1(A)$,*

$$g\Big|_{A \cap f^{-1}(\{y\})} \quad \text{is } \mathcal{H}^{n-m}\text{-summable for } \mathcal{L}^m\text{-a.e. } y \in \mathbb{R}^m \tag{2.24}$$

and it holds that

$$\int_A g(x)|(Jf)(x)| \, d\mathcal{L}^n(x) = \int_{\mathbb{R}^m} \left(\int_{A \cap f^{-1}(\{y\})} g \, d\mathcal{H}^{n-m} \right) d\mathcal{L}^m(y), \tag{2.25}$$

where $Jf = \sqrt{\det\left[(Df)(Df)^\top\right]}$ is the Jacobian of f.

Recall the definition of ω_α from (2.21).

Lemma 2.9. *Let $\Omega \subseteq \mathbb{R}^n$ be a bounded open set which is Jordan measurable (i.e., a bounded open set whose boundary has Lebesgue measure zero). Suppose that $\alpha > 0$ is such that ω_α vanishes at the origin and satisfies a Dini integrability condition, i.e.,*

$$\lim_{r \to 0^+} \omega_\alpha(r) = 0 \quad \text{and} \quad \int_0 \frac{\omega_\alpha(r)}{r} \, dr < +\infty. \tag{2.26}$$

Then for every $t > 0$ one has

$$\int_{\Omega \setminus \Omega_t} \delta_\Omega(x)^{-\alpha} \, dx = \omega_\alpha(t) + \alpha \int_0^t \frac{\omega_\alpha(r)}{r} \, dr. \tag{2.27}$$

In particular, for every $t > 0$ there holds

$$\int_\Omega \delta_\Omega(x)^{-\alpha} \, dx \leq t^{-\alpha} \mathcal{L}^n(\Omega) + \alpha \int_0^t \frac{\omega_\alpha(r)}{r} \, dr < +\infty. \tag{2.28}$$

Proof. Given that we are assuming that the bounded open set $\Omega \subseteq \mathbb{R}^n$ is Jordan measurable it follows that

$$\mathcal{L}^n(\partial\Omega) = 0 \tag{2.29}$$

For each $t > 0$, set $\Omega_t := \{x \in \Omega : \delta_\Omega(x) \geq t\}$, and apply the coarea formula (2.25) with $A := \Omega \setminus \Omega_t$, $g \in L^1(\Omega \setminus \Omega_t)$ arbitrary, and $f : \mathbb{R}^n \to \mathbb{R}$ given by $f(x) := \delta_\Omega(x)$ for each $x \in \mathbb{R}^n$. Then $\partial\Omega_t = \Omega \cap \delta_\Omega^{-1}(\{t\})$ for every $t > 0$ and

$$Jf(x) = \begin{cases} 1 & \text{for } \mathcal{L}^n\text{-a.e. } x \in \Omega, \\ 0 & \text{for } \mathcal{L}^n\text{-a.e. } x \in \mathbb{R}^n \setminus \overline{\Omega}, \end{cases} \tag{2.30}$$

hence for every $t > 0$ we have (making use of (2.29))

$$\int_{\Omega \setminus \Omega_t} g(x) \, dx = \int_0^t \left(\int_{\partial\Omega_r} g \, d\mathcal{H}^{n-1} \right) dr. \tag{2.31}$$

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In particular,

$$\frac{d}{dt}\left(\int_{\Omega\setminus\Omega_t} g(x)\,dx\right) = \int_{\partial\Omega_t} g\,d\mathcal{H}^{n-1} \quad \text{for } \mathcal{H}^1\text{-a.e. } t > 0, \tag{2.32}$$

which, in the case when $g = 1$ yields

$$\mathcal{H}^{n-1}(\partial\Omega_t) = \frac{d}{dt}\left(\mathcal{L}^n(\Omega\setminus\Omega_t)\right) \quad \text{for } \mathcal{H}^1\text{-a.e. } t > 0. \tag{2.33}$$

Specialise now (2.31) to the case when, for some fixed $\alpha > 0$ and $M > 0$, we take

$$g(x) := \min\left\{\delta_\Omega(x)^{-\alpha},\, M\right\}, \qquad \forall\, x \in \Omega\setminus\Omega_t. \tag{2.34}$$

Then $g \in L^1(\Omega\setminus\Omega_t)$ so this choice yields

$$\int_{\Omega\setminus\Omega_t} \min\left\{\delta_\Omega(x)^{-\alpha},\, M\right\} dx = \int_0^t \min\left\{r^{-\alpha},\, M\right\}\mathcal{H}^{n-1}(\partial\Omega_r)\,dr, \tag{2.35}$$

hence, ultimately,

$$\int_{\Omega\setminus\Omega_t} \delta_\Omega(x)^{-\alpha}\,dx = \int_0^t r^{-\alpha}\mathcal{H}^{n-1}(\partial\Omega_r)\,dr, \tag{2.36}$$

after letting $M \nearrow +\infty$ and invoking Lebesgue's Monotone Convergence Theorem. Thus, from (2.33) and (2.36) we obtain

$$\int_{\Omega\setminus\Omega_t} \delta_\Omega(x)^{-\alpha}\,dx = \int_0^t r^{-\alpha}\frac{d}{dr}\left(\mathcal{L}^n(\Omega\setminus\Omega_r)\right) dr. \tag{2.37}$$

Integrating by parts in the right-hand side of (2.37) then gives

$$\int_{\Omega\setminus\Omega_t} \delta_\Omega(x)^{-\alpha}\,dx = t^{-\alpha}\,\mathcal{L}^n(\Omega\setminus\Omega_t) - \lim_{r\to 0^+}\left(r^{-\alpha}\,\mathcal{L}^n(\Omega\setminus\Omega_r)\right)$$
$$+\alpha\int_0^t r^{-\alpha-1}\mathcal{L}^n(\Omega\setminus\Omega_r)\,dr, \tag{2.38}$$

so that

$$\int_{\Omega\setminus\Omega_t} \delta_\Omega(x)^{-\alpha}\,dx = \omega_\alpha(t) - \lim_{r\to 0^+}\omega_\alpha(r) + \alpha\int_0^t \frac{\omega_\alpha(r)}{r}\,dr. \tag{2.39}$$

Now (2.27) readily follows from this, granted (2.26). Finally, (2.28) is an immediate consequence of (2.27), the crude estimate $\int_{\Omega_t} \delta_\Omega(x)^{-\alpha}\,dx \leq t^{-\alpha}\mathcal{L}^n(\Omega_t)$, and the fact that $t^{-\alpha}\mathcal{L}^n(\Omega_t) + \omega_\alpha(t) = t^{-\alpha}\mathcal{L}^n(\Omega)$. $\qquad\square$

Remark 2.10. *Assume that $\Omega \subseteq \mathbb{R}^n$ is a Jordan measurable, bounded open set. Then an inspection of the proof of Lemma 2.9 reveals that*

$$\int_\Omega \delta_\Omega(x)^{-\alpha}\,dx < +\infty \qquad \text{whenever } \dim^*_{Minkowski}(\partial\Omega) < n - \alpha, \tag{2.40}$$

and

$$\int_\Omega \delta_\Omega(x)^{-\alpha}\,dx = +\infty \quad \text{if } \sup\{\gamma \geq 0 : M^*_\gamma(\partial\Omega) < +\infty,\, M_{*,\gamma}(\partial\Omega) > 0\} > n - \alpha. \tag{2.41}$$

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Definition 2.11. *The set $\Sigma \subseteq \mathbb{R}^n$ is said to be Ahlfors regular if there exist finite constants $C_0, C_1 > 0$ as well as a number $R > 0$ such that*

$$C_0 \, r^{n-1} \le \mathcal{H}^{n-1}(B(x,r) \cap \Sigma) \le C_1 \, r^{n-1}, \qquad \forall x \in \Sigma, \ \forall r \in (0, R), \tag{2.42}$$

The triplet C_0, C_1, R makes up what will henceforth be referred to as the Ahlfors character of Σ.

Lemma 2.12. *Let $\Omega \subseteq \mathbb{R}^n$ be a bounded open set whose boundary is Ahlfors regular. Then, for any exponent $\alpha \in [0,1)$ there exists a finite constant $C > 0$, which depends only n, α and the Ahlfors character of $\partial \Omega$, such that*

$$\int_\Omega \delta_\Omega(x)^{-\alpha} \, dx \le C \left[\mathcal{L}^n(\Omega) \right]^{1-\alpha} \left[\mathcal{H}^{n-1}(\partial \Omega) \right]^\alpha. \tag{2.43}$$

This implies that the following generalised isoperimetric inequality holds:

$$\int_\Omega \delta_\Omega(x)^{-\alpha} \, dx \le C \left[\mathcal{H}^{n-1}(\partial \Omega) \right]^{\frac{n-\alpha}{n-1}}. \tag{2.44}$$

In particular, under the same hypotheses,

$$\int_\Omega \delta_\Omega(x)^{-\alpha} \, dx \le C \operatorname{diam}(\Omega)^{n-\alpha}. \tag{2.45}$$

Proof. The version of the isoperimetric inequality proved by H. Federer (cf. 3.2.43-3.2.44 on p. 278 of [15]) reads

$$E \subset \mathbb{R}^n \ \text{ with } \ \mathcal{L}^n(\overline{E}) < +\infty \implies \mathcal{L}^n(\overline{E}) \le \frac{1}{n(\omega_{n-1})^{1/(n-1)}} \left[\mathcal{H}^{n-1}(\partial E) \right]^{\frac{n}{n-1}}, \tag{2.46}$$

where ω_{n-1} denotes the surface area of S^{n-1}. Of course, (2.46) covers the case $\alpha = 0$ of (2.43), so we will assume in what follows that $0 < \alpha < 1$.

To proceed, we note two consequences of the assumption that that $\Omega \subseteq \mathbb{R}^n$ is a bounded open set whose boundary is Ahlfors regular. First, it is clear that (2.29) holds and, hence, Ω is Jordan measurable. Second, it has been proved in [17] that

$$\mathcal{L}^n(\Omega \setminus \Omega_r) \le C r \, \mathcal{H}^{n-1}(\partial \Omega), \qquad \forall r > 0, \tag{2.47}$$

where $C > 0$ depends only on the Ahlfors character of $\partial \Omega$ and, as before, for each $r > 0$ we have set $\Omega_r = \{x \in \Omega : \delta_\Omega(x) \ge r\}$. (Parenthetically, we wish to point out that this estimate implies that $\dim^*_{Minkowski}(\partial \Omega) \le n-1$ whenever $\Omega \subseteq \mathbb{R}^n$ is a bounded open set whose boundary is Ahlfors regular.) In particular, (2.47) entails

$$\omega_\alpha(r) \le C r^{1-\alpha} \, \mathcal{H}^{n-1}(\partial \Omega), \qquad \forall r > 0, \tag{2.48}$$

and, given that $\alpha \in (0,1)$, it follows that the conditions in (2.26) are satisfied. On the basis of this discussion, (2.28) then gives

$$
\begin{aligned}
\int_\Omega \delta_\Omega(x)^{-\alpha} \, dx &\le t^{-\alpha} \mathcal{L}^n(\Omega) + C \, \alpha \left(\int_0^t r^{-\alpha} \, dr \right) \mathcal{H}^{n-1}(\partial \Omega) \\
&= t^{-\alpha} \mathcal{L}^n(\Omega) + \frac{C \, \alpha}{1-\alpha} t^{1-\alpha} \, \mathcal{H}^{n-1}(\partial \Omega),
\end{aligned}
\tag{2.49}
$$

for every $t > 0$. Choosing $t := \mathcal{L}^n(\Omega) / \mathcal{H}^{n-1}(\partial \Omega)$ then readily yields (2.43). Having justified (2.43), then (2.44) follows from this after observing that (2.46) implies

$$\left[\mathcal{L}^n(\Omega) \right]^{1-\alpha} \le C_{n,\alpha} \left[\mathcal{H}^{n-1}(\partial \Omega) \right]^{\frac{n(1-\alpha)}{n-1}}. \tag{2.50}$$

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As regards (2.45), this is going to be a consequence of (2.44) and the fact that for any set $E \subseteq \mathbb{R}^n$ whose boundary is Ahlfors regular there holds

$$\mathcal{H}^{n-1}(\partial E) \leq C \left[\operatorname{diam}(E)\right]^{n-1}, \tag{2.51}$$

where $C > 0$ depends only on the Ahlfors character of ∂E (in fact only the upper estimate in the Ahlfors regularity condition is really needed for this purpose). At this stage there remains to prove (2.51) and, given the dilation and translation invariant nature of this estimate, there is no loss of generality in assuming that $\operatorname{diam}(E) = 1$ and that, in fact, $\overline{E} \subseteq (-1,1)^n$. Partition the cube $(-1,1)^n$ into a grid of congruent subcubes, call them $\{Q\}_{Q \in \mathcal{J}}$, of side-length $R/(2\sqrt{n})$, where $R \in (0,1)$ is such that there exists $C > 0$ for which

$$\mathcal{H}^{n-1}(B(x,r) \cap \partial E) \leq C\, r^{n-1}, \qquad \forall x \in \partial E, \ \forall r \in (0,R). \tag{2.52}$$

Consider $\mathcal{J}_* := \{Q \in \mathcal{J} : Q \cap \partial E \neq \emptyset\}$ and, for each $Q \in \mathcal{J}_*$, select $x_Q \in Q \cap \partial E$. Then, clearly,

$$\partial E \subseteq \bigcup_{Q \in \mathcal{J}_*} B(x_Q, R/2) \tag{2.53}$$

which, when used in conjunction with (2.52) and the fact that $\mathcal{H}^{n-1}$ is an outer measure, gives

$$\begin{aligned} \mathcal{H}^{n-1}(\partial E) &\leq \sum_{Q \in \mathcal{J}_*} \mathcal{H}^{n-1}(\partial E \cap B(x_Q, R/2)) \\ &\leq C(R/2)^{n-1} \cdot \#\mathcal{J}_* \leq 2\, C\, n^{n/2} R^{-1}. \end{aligned} \tag{2.54}$$

This, of course, suits our purposes, so the proof of (2.51) is complete. $\qquad\square$

Here is the proposition alluded to a while ago.

Proposition 2.13. *(i) If $\Omega \subseteq \mathbb{R}^n$ is a bounded domain whose boundary has a finite upper γ-dimensional Minkowski content, where $\gamma < n$, then (1.6) holds for any $\beta \in (0, (n-\gamma)/2)$.*

(ii) If $\Omega \subseteq \mathbb{R}^n$ is a bounded domain whose boundary is Ahlfors regular, then (1.6) holds for any $\beta \in (0, 1/2)$ and moreover

$$\int_\Omega u(x)^{-\beta}\, dx \leq C \left[\mathcal{L}^n(\Omega)\right]^{1-2\beta} \left[\mathcal{H}^{n-1}(\partial\Omega)\right]^{2\beta} \leq C \operatorname{diam}(\Omega)^{n-2\beta} \quad \textit{if} \ \ \beta < 1/2, \tag{2.55}$$

where C depends only on the Ahlfors character of $\partial\Omega$, n and β.

(iii) As far as (2.55) is concerned, the critical value $\beta = 1/2$ is in the nature of best possible in the sense that for every $\beta \in (1/2, 1)$ there exists a bounded domain $\Omega_\beta \subseteq \mathbb{R}^2$ which is regular for the Dirichlet problem and has an Ahlfors regular boundary and with the property that if u solves (2.1) then

$$\int_{\Omega_\beta} u(x)^{-\beta}\, dx = +\infty. \tag{2.56}$$

Proof. The claims in (i)-(ii) follow from Lemma 2.9 and Lemma 2.12, respectively, as soon as we show that the solution of the problem (2.1) is bounded from below by a multiple of the square of the distance function to the boundary. The latter property can be established in a variety of ways. One such approach involves $\delta_{\Omega,\mathrm{reg}}$, the regularised distance function to $\partial\Omega$ (in the sense of Theorem 2, p. 171 in [31]). Recall that this is a C^∞ function in $\mathbb{R}^n$ satisfying $\delta_{\Omega,\mathrm{reg}} \approx \delta_\Omega$ and which has the property that for each multi-index α there exists $C_\alpha > 0$ such that

$$|\partial^\alpha \delta_{\Omega,\mathrm{reg}}(x)| \leq C_\alpha\, \delta_\Omega(x)^{1-|\alpha|}, \quad \forall x \in \Omega. \tag{2.57}$$

In particular, there exists a finite dimensional constant $C > 0$ with the property that $|\Delta(\delta^2_{\Omega,\mathrm{reg}})(x)| \leq C$ for all $x \in \Omega$. This implies that $u - C^{-1}\delta^2_{\Omega,\mathrm{reg}}$ is superharmonic in Ω, continuous on $\overline{\Omega}$, and vanishes on

$\partial\Omega$. Hence, by the Maximum Principle, $u(x) \geq C_n\,\delta_\Omega(x)^2$ for all $x \in \Omega$. However, the sharp version of this estimate is

$$u(x) \geq (2n)^{-1}\delta_\Omega(x)^2, \qquad \text{for every } x \in \Omega, \tag{2.58}$$

and this is established as follows. Fix an arbitrary point $x_0 \in \Omega$ and abbreviate $r := \delta_\Omega(x_0)$. Then for every $\varepsilon \in (0, r)$ we have that $\overline{B(x_0, r - \varepsilon)} \subseteq \Omega$ and $u \in C^\infty(\overline{B(x_0, r - \varepsilon)})$. Next, consider the standard barrier

$$v(x) := (2n)^{-1}\big((r - \varepsilon)^2 - |x - x_0|^2\big), \qquad x \in \overline{B(x_0, r - \varepsilon)}, \tag{2.59}$$

and note that, by the Maximum Principle and the properties of u, we have $u \geq v$ in $\overline{B(x_0, r - \varepsilon)}$. In particular, $u(x_0) \geq v(x_0)$ which gives $u(x_0) \geq (2n)^{-1}(r - \varepsilon)^2$. Hence, after sending ε to zero we obtain $u(x_0) \geq (2n)^{-1}r^2 = (2n)^{-1}\delta_\Omega(x_0)^2$. Given that $x_0 \in \Omega$ has been chosen arbitrarily, (2.58) is proved. This concludes the proof of (i)-(ii).

Concerning (iii), the task is to construct a counterexample to the statement (1.6) in the case when $\beta \in (1/2, 1)$ in the class of bounded domains which are regular for the Dirichlet problem and have Ahlfors regular boundaries. To this end, fix $\beta \in (1/2, 1)$ and consider the curvilinear triangle Ω_β in $\mathbb{R}^2$ given by

$$\Omega_\beta := \big\{(x, y) \in \mathbb{R}^2 : 0 < x < 1 \ \text{ and } \ 0 < y < \varepsilon\,x^{1/(2\beta - 1)}\big\}, \tag{2.60}$$

where $\varepsilon = \varepsilon(\beta)$ is a sufficiently small positive constant, to be specified momentarily. Clearly, the function

$$v(x, y) := y\big(\varepsilon x^{1/(2\beta - 1)} - y\big), \qquad \forall\,(x, y) \in \Omega_\beta, \tag{2.61}$$

is positive in the domain Ω_β and is nonnegative on its boundary. In addition, for all $(x, y) \in \Omega_\beta$ we have

$$
\begin{aligned}
-(\Delta v)(x, y) &= 2 - \frac{2\varepsilon(1 - \beta)}{(2\beta - 1)^2}\,x^{(3 - 4\beta)/(2\beta - 1)}\,y \\[2mm]
&\geq 2 - \frac{2\varepsilon^2(1 - \beta)}{(2\beta - 1)^2}\,x^{(4 - 4\beta)/(2\beta - 1)} \\[2mm]
&\geq 2 - 2\varepsilon^2(1 - \beta)(2\beta - 1)^{-2}, \tag{2.62}
\end{aligned}
$$

where the last step makes essential use of the fact that $\beta \in (1/2, 1)$. At this stage, pick $\varepsilon > 0$ sufficiently small so that the last expression in (2.62) is ≥ 1. Such a choice forces $u - v$ to be subharmonic in Ω_β, if u is the solution of the Saint Venant problem in Ω_β (cf. (2.1)). In addition, $u - v \leq 0$ on $\partial\Omega_\beta$ by design. The Maximum Principle then gives that $u \leq v$ in Ω_β. Consequently, we may estimate

$$
\begin{aligned}
\int_{\Omega_\beta} u(x, y)^{-\beta}\,dxdy &\geq \int_{\Omega_\beta} v(x, y)^{-\beta}\,dxdy \\[2mm]
&= \int_0^1 \bigg(\int_0^{\varepsilon x^{1/(2\beta - 1)}} y^{-\beta}\big(\varepsilon x^{1/(2\beta - 1)} - y\big)^{-\beta}\,dy\bigg)dx \\[2mm]
&= \varepsilon^{1 - 2\beta}\bigg(\int_0^1 x^{-1}\,dx\bigg)\bigg(\int_0^1 t^{-\beta}(1 - t)^{-\beta}\,dt\bigg) = +\infty, \tag{2.63}
\end{aligned}
$$

after making the change of variables $y = \varepsilon x^{1/(2\beta - 1)}\,t$ in the inner integral in the second line. This completes the proof of the proposition. $\qquad\square$

3 Barrier functions and domains satisfying a cone condition

Here and elsewhere S^{n-1} stands for the unit sphere in $\mathbb{R}^n$. We denote by $\Gamma_\theta^R(x_0, \eta)$ the open, one-component circular cone in $\mathbb{R}^n$ with vertex at $x_0 \in \mathbb{R}^n$, half-aperture $\theta \in (0, \pi)$, axis along $\eta \in S^{n-1}$, and (roundly) truncated at $R > 0$, i.e.,

$$\Gamma_\theta^R(x_0, \eta) := \{x \in \mathbb{R}^n : (x - x_0) \cdot \eta > |x - x_0| \cos\theta \ \text{ and } \ |x - x_0| < R\}, \tag{3.1}$$

When $R = +\infty$ (that is, the cone is infinite) we agree to simply write $\Gamma_\theta(x_0, \eta)$. Furthermore, we use the abbreviation Γ_θ^R (respectively, Γ_θ, when $R = +\infty$) whenever $x_0 = 0 \in \mathbb{R}^n$ and $\eta = e_n := (0, ..., 0, 1) \in S^{n-1}$.

Of course, $\Gamma_\theta \cap S^{n-1}$ is the spherical cap with centre at the north pole and (spherical) radius θ. More generally, given an open, connected subset $\mathfrak{G}$ of S^{n-1}, we denote by $\Gamma_\mathfrak{G}$ the open cone in $\mathbb{R}^n$ with vertex at the origin and shape $\mathfrak{G}$, i.e.,

$$\Gamma_\mathfrak{G} := \{\rho\,\omega : \rho > 0 \ \text{ and } \ \omega \in \mathfrak{G}\}. \tag{3.2}$$

Going further, we let $\Delta_{S^{n-1}}$ stand for the Laplace-Beltrami operator on S^{n-1} and fix an open, connected subset $\mathfrak{G}$ of S^{n-1} with the property that $\partial_{S^{n-1}}\mathfrak{G}$, the boundary of $\mathfrak{G}$ relative to S^{n-1}, is sufficiently regular. In this setting, we let $\Lambda_\mathfrak{G} > 0$ be the first positive eigenvalue of the nonnegative operator $-\Delta_{S^{n-1}}$ equipped with (homogeneous) Dirichlet boundary condition on $\mathfrak{G}$ and denote by $\phi_\mathfrak{G}$ an eigenfunction corresponding to the eigenvalue $\Lambda_\mathfrak{G}$. Hence,

$$-\Delta_{S^{n-1}}\phi_\mathfrak{G} = \Lambda_\mathfrak{G}\,\phi_\mathfrak{G} \ \text{ in } \ \mathfrak{G}, \quad \text{and} \quad \phi_\mathfrak{G} = 0 \ \text{ on } \ \partial_{S^{n-1}}\mathfrak{G}. \tag{3.3}$$

Recall that any eigenfunction corresponding to $\Lambda_\mathfrak{G}$ does not change sign in $\mathfrak{G}$ (see, e.g., the discussion on p. 42-43 in [10] in the case of a spherical cap), Since $\phi_\mathfrak{G}$ is uniquely determined only up to a re-normalisation, it follows that there is no loss of generality in assuming that

$$\phi_\mathfrak{G} > 0 \ \text{ in } \ \mathfrak{G}, \quad \text{and} \ \sup_\mathfrak{G} \phi_\mathfrak{G} = 1. \tag{3.4}$$

For further reference let us also record here that, granted sufficient regularity for $\partial_{S^{n-1}}\mathfrak{G}$, the function $\phi_\mathfrak{G}$ behaves essentially like the distance to the boundary of $\mathfrak{G}$. More precisely, if $\partial_{S^{n-1}}\mathfrak{G}$ is of class C^1, then the following estimate (which is going to be useful in §5) holds

$$\phi_\mathfrak{G}(\omega) \approx \text{dist}_{S^{n-1}}\big(\omega, \partial_{S^{n-1}}\mathfrak{G}\big), \quad \text{uniformly for} \ \omega \in \mathfrak{G}, \tag{3.5}$$

where $\text{dist}_{S^{n-1}}(\omega, \omega') := \arccos(\omega \cdot \omega')$, for $\omega, \omega' \in S^{n-1}$, denotes the geodesic distance on S^{n-1}. This property is a consequence of the fact that $\phi_\mathfrak{G}$ is the solution of the boundary value problem (3.3) and that $\partial_{S^{n-1}}\mathfrak{G}$ is a C^1 submanifold of codimension one in the unit sphere.

Corresponding to the case when $\mathfrak{G}$ is a spherical cap, say $\mathfrak{G} = S^{n-1} \cap \Gamma_\theta$ for some $\theta \in (0, \pi)$, we agree to write ϕ_θ and Λ_θ in place of $\phi_{S^{n-1} \cap \Gamma_\theta}$ and $\Lambda_{S^{n-1} \cap \Gamma_\theta}$, respectively. Hence, in particular,

$$-\Delta_{S^{n-1}}\phi_\theta = \Lambda_\theta\,\phi_\theta \quad \text{in} \ \ S^{n-1} \cap \Gamma_\theta \ \text{ and } \ \phi_\theta = 0 \ \text{ on } \ S^{n-1} \cap \partial\Gamma_\theta.$$
$$\phi_\theta > 0 \ \text{ in } \ S^{n-1} \cap \Gamma_\theta \quad \text{and} \quad \sup_{S^{n-1} \cap \Gamma_\theta} \phi_\theta = 1 \quad \text{for each} \ \ \theta \in (0, \pi). \tag{3.6}$$

Definition 3.1. *Given an open, connected subset $\mathfrak{G}$ of S^{n-1}, with a sufficiently regular boundary (relative to S^{n-1}), we associate the index $\alpha_\mathfrak{G}$ defined by*

$$\alpha_\mathfrak{G} := -\tfrac{n-2}{2} + \sqrt{\tfrac{(n-2)^2}{4} + \Lambda_\mathfrak{G}}, \tag{3.7}$$

that is, the unique positive root of the equation

$$\alpha_\mathfrak{G}(\alpha_\mathfrak{G} + n - 2) = \Lambda_\mathfrak{G}. \tag{3.8}$$

(Note that since $\Lambda_{\mathfrak{G}} > 0$, these considerations are meaningful.) Finally, for each $\theta \in (0, \pi)$, we abbreviate $\alpha_{S^{n-1} \cap \Gamma_\theta}$ by α_θ. Hence, in this notation,

$$\alpha_\theta = -\frac{n-2}{2} + \sqrt{\frac{(n-2)^2}{4} + \Lambda_\theta} \quad and \quad \Lambda_\theta = \alpha_\theta(\alpha_\theta + n - 2). \tag{3.9}$$

for any $\theta \in (0, \pi)$.

The format of (3.8) is suggested by the formula for the Euclidean Laplacian in spherical polar coordinates $x = \rho\,\omega \in \mathbb{R}^n \setminus \{0\}$, with $\rho := |x| > 0$ and $\omega := x/|x| \in S^{n-1}$, i.e.,

$$\Delta f = \rho^{1-n}\partial_\rho(\rho^{n-1}\partial_\rho f) + \rho^{-2}\Delta_{S^{n-1}}f = \partial_\rho^2 f + (n-1)\rho^{-1}\partial_\rho f + \rho^{-2}\Delta_{S^{n-1}}f. \tag{3.10}$$

Indeed, introducing the barrier function $v_{\mathfrak{G}} : \Gamma_{\mathfrak{G}} \to \mathbb{R}$ by setting

$$v_{\mathfrak{G}}(x) := \rho^{\alpha_{\mathfrak{G}}}\phi(\omega) = |x|^{\alpha_{\mathfrak{G}}}\phi_{\mathfrak{G}}\left(\frac{x}{|x|}\right) \quad \text{for} \quad \omega = \frac{x}{|x|} \in \mathfrak{G} \subseteq S^{n-1} \text{ and } \rho = |x| > 0, \tag{3.11}$$

it follows that, for each $\omega \in \mathfrak{G}$ and $\rho > 0$,

$$(\Delta v_{\mathfrak{G}})(\rho\,\omega) = [\alpha_{\mathfrak{G}}(\alpha_{\mathfrak{G}} + n - 2) - \Lambda_{\mathfrak{G}}]\rho^{-2}\phi_{\mathfrak{G}}(\omega) = 0 \tag{3.12}$$

precisely for the choice (3.8). This ensures that the function $v_{\mathfrak{G}}$ is harmonic in the cone $\Gamma_{\mathfrak{G}}$. In summary, taking $\alpha_{\mathfrak{G}}$ as in (3.7) ensures that

$$\Delta v_{\mathfrak{G}} = 0 \text{ in } \Gamma_{\mathfrak{G}}, \quad v_{\mathfrak{G}} = 0 \text{ on } \partial\Gamma_{\mathfrak{G}}, \quad v_\theta > 0 \text{ in } \Gamma_{\mathfrak{G}}. \tag{3.13}$$

In the axially symmetric case, i.e., when $\mathfrak{G} = S^{n-1} \cap \Gamma_\theta$ for some $\theta \in (0, \pi)$, a good deal is known about the properties enjoyed by the exponent α_θ introduced in (3.9) (see, e.g., Theorem 3 on p. 44, Theorem 6 on p. 50 in [10] and the discussion on p. 112 of [4]). Specifically, for each $n \geq 2$ one has

$$(0, \pi) \ni \theta \mapsto \alpha_\theta \in (0, +\infty) \quad \text{is strictly decreasing and continuous,} \tag{3.14}$$

$$\alpha_{\pi/2} = 1 \quad \text{and} \quad \lim_{\theta \searrow 0} \alpha_\theta = +\infty, \tag{3.15}$$

$$\lim_{\theta \nearrow \pi} \alpha_\theta = 0 \quad \text{if } n \geq 3, \tag{3.16}$$

$$\alpha_\theta = \frac{\pi}{2\theta} \quad \text{if } n = 2, \quad \text{and} \quad \alpha_\theta = \frac{\pi}{\theta} - 1 \quad \text{if } n = 4, \tag{3.17}$$

$$\alpha_\theta \in \left(\tfrac{1}{2}, +\infty\right) \quad \text{and} \quad \lim_{\theta \nearrow \pi} \alpha_\theta = \tfrac{1}{2} \quad \text{if } n = 2, \tag{3.18}$$

$$\alpha_\theta = 2 \iff \theta = \arccos\left(1/\sqrt{n}\right). \tag{3.19}$$

The computations in the case $n = 2$ are particularly simple. Indeed, the eigenvalue problem for the Dirichlet-Laplacian on the one-dimensional arc $\{e^{i\omega} : -\theta < \omega < \theta\}$ in the unit circle becomes (with 'prime' denoting the angular derivative $d/d\omega$) $\phi''(\omega) + \Lambda\phi(\omega) = 0$ for $-\theta < \omega < \theta$, $\phi(-\theta) = \phi(\theta) = 0$. The smallest positive eigenvalue is then $\Lambda = \Lambda_\theta = \frac{\pi^2}{(2\theta)^2}$ which, in light of (3.9), gives the first formula in (3.17). In the higher dimensional setting, the eigenvalue problem on a spherical cap leads to a less transparent equation. To describe this, recall that the so-called Gegenbauer functions, $\mathcal{C}_\alpha^\nu(z)$, are the solutions of Gegenbauer's differential equation

$$(z^2 - 1)\frac{d^2g}{dz^2} + (2\nu + 1)z\frac{dg}{dz} - \alpha(\alpha + 2\nu)g = 0, \qquad z, \nu, \alpha \in \mathbb{C}. \tag{3.20}$$

When considered with the variable z restricted to the interval $(-1, 1)$ on the real axis, the above second-order ODE is endowed with the initial conditions

$$g(-1) = 1 \quad \text{and} \quad \frac{dg}{dz}(-1) = -\frac{\alpha(\alpha + 2\nu)}{2\nu + 1}. \tag{3.21}$$

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For more details on this subject see, e.g., [26]. In the present context, the key feature of the Gegenbauer functions is that the exponent α_θ from (3.9) coincides with the first positive zero of the mapping $\alpha \mapsto \mathcal{C}_\alpha^{\frac{n-2}{2}}(-\cos\theta)$; compare with Lemma 6.6.3 in [22]. For example, the continuity of (3.14) follows from this representation and classical results on the dependence of the solution of ODE's on parameters. For related material see also [29] (especially Theorem 2, p. 308), and [25] (where, in lieu of (3.20), the authors work with an ODE for $f_{n,\alpha}(\theta) := \mathcal{C}_\alpha^{\frac{n-2}{2}}(-\cos\theta)$).

We continue by recording the definition of the class of nontangentially accessible domains (introduced by Jerison and Kenig in [18]), and by making a couple of remarks.

Definition 3.2. *A nonempty, proper open subset Ω of $\mathbb{R}^n$ is called an NTA domain provided Ω satisfies both an interior and an exterior corkscrew condition (with constants M, r_* as in Definition 2.2) and Ω satisfies a Harnack chain condition, defined as follows (with reference to M as above).*

If $x_1, x_2 \in \Omega$ are such that $\delta_\Omega(x_i) \geq \varepsilon$ for $i = 1, 2$, and $|x_1 - x_2| \leq 2^k \varepsilon > 0$, then there exist Mk balls $B_j \subseteq \Omega$, $1 \leq j \leq Mk$, such that

(i) $x_1 \in B_1$, $x_2 \in B_{Mk}$ and $B_j \cap B_{j+1} \neq \emptyset$ for $1 \leq j \leq Mk - 1$;

(ii) each ball B_j has a radius r_j satisfying

$$M^{-1} r_j \leq \mathrm{dist}(B_j, \partial\Omega) \leq M r_j \quad and \quad r_j \geq M^{-1} \min\{\delta_\Omega(x_1), \delta_\Omega(x_2)\}. \tag{3.22}$$

Two comments are going to be of importance for us later on. First, the relevance of the Harnack chain condition is that, thanks to Harnack's inequality, if w is a positive harmonic function in Ω then, in the context of the second part of Definition 3.2,

$$M^{-k} w(x_1) \leq w(x_2) \leq M^k w(x_1). \tag{3.23}$$

Second, any bounded NTA domain is regular for the Dirichlet problem (it suffices to recall that any such domain satisfies an exterior corkscrew condition).

Moving on, a bounded domain Ω in $\mathbb{R}^n$ is called a Lipschitz domain provided Ω and its boundary $\partial\Omega$ locally coincide with, respectively, upper-graphs and graphs of Lipschitz function functions. In this vein, recall that a function $f : D \to \mathbb{R}$ where, say, $D \subseteq \mathbb{R}^{n-1}$, is called Lipschitz provided there exists $M > 0$ so that $|f(x) - f(y)| \leq M|x - y|$ for any $x, y \in D$. A formal definition is given below.

Definition 3.3. *A bounded domain $\Omega \subset \mathbb{R}^n$ is called Lipschitz (of class C^α, respectively) if for any $x_0 \in \partial\Omega$ there exist $r, h > 0$ and a coordinate system $\{x_1, \ldots, x_n\}$ in $\mathbb{R}^n$ (isometric to the canonical one) with origin at x_0 along with a function $\varphi : \mathbb{R}^{n-1} \to \mathbb{R}$ which is Lipschitz (of class C^α, respectively) and for which the following property holds. If $C(r, h)$ denotes the open cylinder $\{x = (x', x_n) : |x'| < r \text{ and } -h < x_n < h\} \subset \mathbb{R}^n$, then*

$$\begin{aligned} \partial\Omega \cap C(r, h) &= \{x = (x', x_n) : |x'| < r \text{ and } x_n = \varphi(x_1, \ldots, x_{n-1})\}, \\ \Omega \cap C(r, h) &= \{x = (x', x_n) : |x'| < r \text{ and } \varphi(x_1, \ldots, x_{n-1}) < x_n < h\}. \end{aligned} \tag{3.24}$$

Fix an atlas for $\partial\Omega$, i.e. a finite collection of cylinders $\{C_k(r_k, h_k)\}_{1 \leq k \leq N}$ (with associated Lipschitz maps $\{\varphi_k\}_{1 \leq k \leq N}$) covering $\partial\Omega$. The Lipschitz constant of Ω, denoted in what follows by κ_Ω, is defined as the infimum of $\max\{\|\nabla\varphi_k\|_{L^\infty} : 1 \leq k \leq N\}$ taken over all possible atlases of $\partial\Omega$.

Clearly, any bounded Lipschitz domain is NTA (hence regular for the Dirichlet problem), and has an Ahlfors regular boundary. For further reference, let us also remark here that

$$\Omega \subseteq \mathbb{R}^n \text{ bounded } C^1 \text{ domain} \implies \kappa_\Omega = 0. \tag{3.25}$$

For an open set $\Omega \subseteq \mathbb{R}^n$ and a number $R > 0$, define Ω_R as the collection of points in Ω at distance at least R from the boundary, i.e.,

$$\Omega_R := \{x \in \Omega : \delta_\Omega(x) > R\}. \tag{3.26}$$

Definition 3.4. *We say that an open set $\Omega \subseteq \mathbb{R}^n$ satisfies an (axially symmetric) inner cone condition with half-aperture $\theta \in (0, \pi/2)$ provided there exists $R \in (0, \operatorname{diam}(\Omega))$ with the property that*

$$\forall\, x \in \overline{\Omega} \setminus \Omega_R \quad \exists\, \eta \in S^{n-1} \quad such\ that\ \ \Gamma_\theta^R(x, \eta) \subseteq \Omega. \tag{3.27}$$

More generally, given an open connected C^1 subdomain $\mathfrak{G}$ of S^{n-1}, we say that $\Omega \subseteq \mathbb{R}^n$ satisfies an inner cone condition with smooth profile $\mathfrak{G}$ provided there exists $R \in (0, \operatorname{diam}(\Omega))$ so that

$$\begin{aligned} &\forall\, x \in \overline{\Omega} \setminus \Omega_R\ \exists\, U\ isometry\ of\ \mathbb{R}^n\ for\ which\\ &U(0) = x\ and\ U(\Gamma_{\mathfrak{G}} \cap B(0, R)) \subseteq \Omega. \end{aligned} \tag{3.28}$$

Definition 3.5. *Given a bounded Lipschitz domain $\Omega \subseteq \mathbb{R}^n$, with Lipschitz constant $\kappa_\Omega \in [0, +\infty)$, define α_Ω to be the index associated as in Definition 3.1 for the angle*

$$\theta = \theta_\Omega := \arctan\left(\frac{1}{\kappa_\Omega}\right) \in (0, \tfrac{\pi}{2}), \tag{3.29}$$

that is, $\alpha_\Omega = \alpha_{\theta_\Omega}$.

In the context of the above definitions, it is illuminating to point out that, in the class of bounded Lipschitz domains,

$$\alpha_\Omega \geq 1 \text{ and, in fact, } \alpha_\Omega \searrow 1 \text{ as } \kappa_\Omega \searrow 0. \tag{3.30}$$

Indeed, this follows readily from (3.14) and (the first formula in) (3.15). On the other hand, by the second formula in (3.15),

$$\alpha_\Omega \nearrow +\infty \ \text{ as } \ \kappa_\Omega \nearrow +\infty. \tag{3.31}$$

It is also straightforward to check that

$$\begin{aligned} &\text{every bounded Lipschitz domain } \Omega \subseteq \mathbb{R}^n \text{ satisfies an}\\ &\text{inner cone condition with half-aperture } \theta, \text{ for any } \theta \in (0, \theta_\Omega). \end{aligned} \tag{3.32}$$

In the opposite direction we note that there exist bounded NTA domains which satisfy an inner cone condition but which are not necessarily Lipschitz (take, for example, the union on two disjoint, open truncated circular cones with a common vertex).

We now proceed to discuss a useful bound from below for the Green function associated with the Dirichlet Laplacian in bounded NTA domains satisfying an inner cone condition. It should be noted that in the class of bounded Lipschitz domains and for a more restrictive concept of cone condition, [25, Proposition 2, p. 272] contains such an estimate. Most unsatisfactorily for us, though, is the fact that this estimate is written in [25] in a form which is not sufficiently explicit for our purposes. Here, we largely follow the approach in [33] with the goal of monitoring how the geometrical characteristics of Ω enter the final estimate.

Proposition 3.6. *Assume that $\Omega \subseteq \mathbb{R}^n$ is a bounded NTA domain which satisfies an inner cone condition with smooth profile $\mathfrak{G} \subseteq S^{n-1}$. Let $\alpha_{\mathfrak{G}}$ be the index associated with the subdomain $\mathfrak{G}$ of S^{n-1} as in Definition 3.1, and fix $R \in (0, \operatorname{diam}(\Omega)/4)$ such that (3.28) holds.*

Then, if $n \geq 3$, there exists a finite constant $c = c(n, \mathfrak{G}) > 0$ with the property that the Green function $G(\cdot, \cdot)$ for the Dirichlet Laplacian in Ω satisfies the dilation invariant estimate

$$G(x, y) \geq c(n, \mathfrak{G}) \left(\frac{\delta_\Omega(x)}{\operatorname{diam}(\Omega)}\right)^{\alpha_{\mathfrak{G}}} \left(\frac{R}{\operatorname{diam}(\Omega)}\right)^m R^{2-n}, \ \ for\ every\ x \in \Omega\ and\ y \in \Omega_R, \tag{3.33}$$

where $m > 0$ depends only on the NTA constants of Ω. Furthermore, a similar conclusion holds in the case when $n = 2$ provided the factor R^{2-n} in the right-hand side of (3.33) is replaced by $\log\left(\operatorname{diam}(\Omega)/R\right)$.

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Proof. We shall only consider the case $n \geq 3$, since the two-dimensional case is treated analogously. The proof is divided into several steps, starting with

Step 1. Assume that $\mathfrak{G}$ is a connected, C^1 smooth subdomain of S^{n-1} and recall the barrier function $v_{\mathfrak{G}}$ from (3.11). Also, fix $z \in \mathfrak{G}$. Then there exists a finite constant $C(\mathfrak{G}, z) > 0$ with the property that for every $r > 0$ one has

$$w(rz)\, v_{\mathfrak{G}}(x) \leq C(\mathfrak{G}, z)\, r^{\alpha_{\mathfrak{G}}} w(x), \qquad \forall\, x \in \Gamma_{\mathfrak{G}} \cap B(0, r), \tag{3.34}$$

for every function

$$w \in C^0(\overline{\Gamma_{\mathfrak{G}} \cap B(0, 2r)}) \;\; \text{satisfying} \;\; w > 0 \;\; \text{and} \;\; \Delta w = 0 \;\; \text{in} \;\; \Gamma_{\mathfrak{G}} \cap B(0, 2r). \tag{3.35}$$

It suffices to establish the above claim in the case when $r = 1$, since then (3.34) follows by rescaling. If this is the case, by considering $x \mapsto w(x)/w(z)$ in place of $w(x)$, there is also no loss of generality in assuming that $w(z) = 1$. In this scenario, the desired conclusion follows from the Maximum Principle as soon as we show that there exists some finite constant $C(\mathfrak{G}, z) > 0$ such that

$$v_{\mathfrak{G}}(x) \leq C(\mathfrak{G}, z)\, w(x), \qquad \forall\, x \in \Gamma_{\mathfrak{G}} \cap \partial B(0, 1), \tag{3.36}$$

for every positive function $w \in C^0(\overline{\Gamma_{\mathfrak{G}} \cap B(0, 2)})$ which is harmonic in $\Gamma_{\mathfrak{G}} \cap B(0, 2)$ and satisfies $w(z) = 1$. With this goal in mind, we then observe that, by Harnack's inequality and the smoothness of $\mathfrak{G}$, there exists $C = C(\mathfrak{G}, z) > 0$ with the property that

$$w(\omega) \geq C \operatorname{dist}_{S^{n-1}}\big(\omega, \partial_{S^{n-1}} \mathfrak{G}\big), \qquad \forall\, \omega \in \mathfrak{G}, \tag{3.37}$$

whereas, by virtue of (3.5),

$$v_{\mathfrak{G}}(\omega) \approx \operatorname{dist}_{S^{n-1}}\big(\omega, \partial_{S^{n-1}} \mathfrak{G}\big), \quad \text{uniformly for} \;\; \omega \in \mathfrak{G}. \tag{3.38}$$

In concert, (3.37) and (3.38) establish estimate (3.36), thus concluding the proof of the claim in Step 1.

Step 2. Suppose that $\Omega \subseteq \mathbb{R}^n$ is a bounded NTA domain. Then there exists a dimensional constant $C_n > 0$ and some $m > 0$ which depends only on the NTA constants of Ω with the property that for each $R \in (0, \operatorname{diam}(\Omega))$ the Green function associated with the Dirichlet Laplacian in Ω satisfies

$$G(x, y) \geq C_n \Big(\frac{R}{\operatorname{diam}(\Omega)}\Big)^m R^{2-n}, \quad \text{for every} \;\; x, y \in \Omega_R. \tag{3.39}$$

To justify this claim, recall the constant M from Definition 3.2 and pick $m > 0$ such that $M = 2^m$. Going further, fix $x, y \in \Omega_R$ and select a point $y_o \in B(y, R/2) \setminus B(y, R/4)$. Consider now a Harnack chain of balls joining x and y_o in Ω. More specifically, pick a natural number $k \sim \log_2(|x - y|/R)$ and suppose $B_j \subseteq \Omega$, $1 \leq j \leq Mk$, is a family of balls such that x is the centre of B_1, y_o is the centre of B_{Mk}, $B_j \cap B_{j+1} \neq \emptyset$ for $1 \leq j \leq Mk-1$, each ball B_j has a radius r_j satisfying $M^{-1} r_j \leq \operatorname{dist}(B_j, \partial\Omega) \leq M r_j$, as well as $r_j \geq M^{-1} \min\{\delta_\Omega(x), \delta_\Omega(y_o)\}$. Then, by repeated applications of Harnack's inequality (compare with (3.23)), we obtain

$$G(x, y) \geq C_n M^{-k} G(y_o, y) \geq C_n \Big(\frac{R}{|x - y_o|}\Big)^m R^{2-n}, \tag{3.40}$$

by the choice of m, k, and y_o, and thanks to (2.5). Since $|x - y_o| \leq \operatorname{diam}(\Omega)$, (3.39) follows.

A moment's reflection shows that (3.39) implies (3.33) in the case when $x, y \in \Omega_R$. We continue with:

Step 3. Here we prove the inequality stated in (3.33) in the case when $0 < R < \operatorname{diam}(\Omega)/4$ and when $y \in \Omega_{4R}$ and $x \in \Omega \setminus (\Omega_R \cup B(y, 2R))$. Assuming that two such points have been fixed, pick $x_* \in \partial\Omega$ such that $\delta_\Omega(x) = |x - x_*|$, and introduce $x_0 := \frac{1}{2}(x + x_*)$. Also, choose an isometry U of $\mathbb{R}^n$ with $U(0) = x_0$ and $U(\Gamma_{\mathfrak{G}} \cap B(0, 2R)) \subseteq \Omega$. It follows that if $z \in \mathfrak{G}$ is fixed, then there exists a finite constant $C = C(\mathfrak{G}, z) \geq 1$ such that $RU(z) \in \Omega_{R/C}$. For the reader's convenience, the special case when $\mathfrak{G}$ is a spherical cap on S^{n-1} with half-angle $\theta \in (0, \pi/2)$ and when $U(\Gamma_{\mathfrak{G}} \cap B(0, 2R)) = \Gamma_\theta^{2R}(x_0, \eta)$ for some $\eta \in S^{n-1}$ is sketched in the picture below:
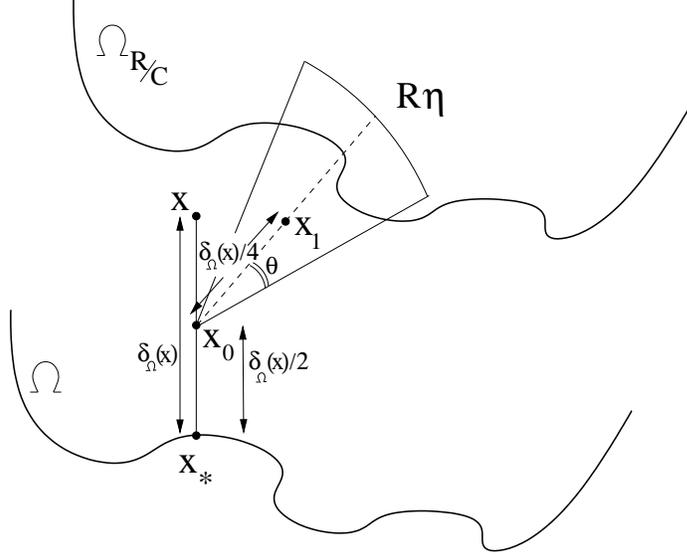
To continue, introduce $x_1 := U(\delta_\Omega(x)z/4)$ which, given that $\delta_\Omega(x) \leq R$, belongs to the cone $U(\Gamma_{\mathfrak{G}} \cap B(0, 2R))$. Then, on the one hand, Harnack's inequality gives

$$G(x, y) \approx G(x_0, y) \approx G(x_1, y), \tag{3.41}$$

with universal comparability constants, while on the other hand, (3.34) applied to the function $w := G(\cdot, y)$ yields

$$
\begin{aligned}
G(x_1, y) &\geq C(n, \mathfrak{G}) R^{-\alpha_{\mathfrak{G}}} |x_1 - x_0|^{\alpha_{\mathfrak{G}}} G(RU(z), y) \\
&\geq C(n, \mathfrak{G}) R^{-\alpha_{\mathfrak{G}}} \delta_\Omega(x)^{\alpha_{\mathfrak{G}}} \left(\frac{R}{\operatorname{diam}(\Omega)} \right)^m R^{2-n},
\end{aligned}
\tag{3.42}
$$

where the last inequality utilises (3.39) and the fact that $RU(z) \in \Omega_{R/C}$. Now (3.33) follows in the case we are currently considering from (3.41) and (3.42) (here we also use the fact that $0 < R/\operatorname{diam}(\Omega) < 1$ and that $\alpha_{\mathfrak{G}} > 0$).

The final arguments in the proof of (3.33) are contained in:

Step 4. When $y \in \Omega_{4R}$ and $x \in (\Omega \setminus \Omega_R) \cap B(y, 2R)$ we have $\delta_\Omega(y)/2 \geq 2R \geq |x - y|$, so (2.5) gives $G(x, y) \geq C_n |x - y|^{2-n} \geq C_n R^{2-n}$. This is good enough to justify (3.33) in this case. Granted this and the cases treated in Steps 2-3, it follows that (3.33) has been proved whenever $y \in \Omega_{4R}$ and $x \in \Omega$. After relabeling, we may therefore conclude that (3.33) holds as stated. □

The estimate in Proposition 3.6 plays a basic role in our next theorem, which is the main result in this section.

Theorem 3.7. *Assume that $\Omega \subseteq \mathbb{R}^n$ is a bounded NTA domain, with an Ahlfors regular boundary, and which satisfies an inner cone condition with smooth profile $\mathfrak{G} \subseteq S^{n-1}$. As usual, we denote by $\alpha_{\mathfrak{G}}$ the index associated with the subdomain $\mathfrak{G}$ of S^{n-1} as in Definition 3.1.*

Next, let $R \in (0, \operatorname{diam}(\Omega)/4)$ be such that (3.28) holds and suppose that $0 < \beta < 1/\alpha_{\mathfrak{G}}$. Also, recall that Ω_R has been introduced in (3.26). Then, if $n \geq 3$, the solution u of (2.1) satisfies the dilation

invariant estimate

$$
\begin{aligned}
\int_{\Omega} u(x)^{-\beta}\,dx \;\le\; & \; C_{\Omega}(n,\mathfrak{G},\beta)\Big(\frac{\operatorname{diam}(\Omega)}{R}\Big)^{(m+\alpha_{\mathfrak{G}})\beta}\Big(\frac{\mathcal{H}^{n-1}(\partial\Omega)}{R^{n-1}}\Big)^{\alpha_{\mathfrak{G}}\beta}\times \\
& \times\Big(\frac{|\Omega|}{R^{n}}\Big)^{2\beta/n-\alpha_{\mathfrak{G}}\beta}\Big(\frac{R^{n}}{|\Omega_{R}|}\Big)^{\beta}|\Omega|^{1-2\beta/n} \\
\le\; & \; C_{\Omega}(n,\mathfrak{G},\beta)\Big(\frac{\operatorname{diam}(\Omega)}{R}\Big)^{n+m\beta}\Big(\frac{R^{n}}{|\Omega_{R}|}\Big)^{1+(n-2)\beta/n}|\Omega|^{1-2\beta/n} \qquad (3.43)
\end{aligned}
$$

where $m > 0$ depends only on the NTA constants of Ω, and $C_{\Omega}(n,\mathfrak{G},\beta) > 0$ is a finite constant which depends only on the Ahlfors character of $\partial\Omega$, the dimension n, the profile $\mathfrak{G}$ and the parameter β.

Corresponding to $n = 2$, assume that $\Omega \subseteq \mathbb{R}^2$ is a bounded NTA domain with an Ahlfors regular boundary, satisfies the inner cone condition (3.27) with half-aperture $\theta \in (0,\pi/2)$ and height $R \in (0,\operatorname{diam}(\Omega)/4)$. Then if $0 < \beta < 2\theta/\pi$, the solution u of (2.1) satisfies the version of (3.43) written for $n = 2$, i.e., the dilation invariant estimate

$$
\begin{aligned}
\int_{\Omega} u(x)^{-\beta}\,dx \;\le\; & \; C_{\Omega}(\theta,\beta)\Big(\frac{\operatorname{diam}(\Omega)}{R}\Big)^{(m+\alpha_{\mathfrak{G}})\beta}\Big(\frac{\mathcal{H}^{1}(\partial\Omega)}{R}\Big)^{\alpha_{\mathfrak{G}}\beta}\times \\
& \times\Big(\frac{|\Omega|}{R^{2}}\Big)^{\beta-\alpha_{\mathfrak{G}}\beta}\Big(\frac{R^{2}}{|\Omega_{R}|}\Big)^{\beta}|\Omega|^{1-\beta} \qquad (3.44)
\end{aligned}
$$

where, as before, $m > 0$ depends only on the NTA constants of Ω, and $C_{\Omega}(\theta,\beta) > 0$ is a finite constant which depends only on the Ahlfors character of $\partial\Omega$, the angle θ, and the parameter β.

Proof. Suppose that $n \ge 3$. The representation in (2.2), together with the nonnegativity of the Green function and estimate (3.33) give

$$
\begin{aligned}
u(x) \;=\; & \int_{\Omega} G(x,y)\,dy \ge \int_{\Omega_{R}} G(x,y)\,dy \\
\ge\; & \; c(n,\mathfrak{G})\Big(\frac{\delta_{\Omega}(x)}{\operatorname{diam}(\Omega)}\Big)^{\alpha_{\mathfrak{G}}}\Big(\frac{R}{\operatorname{diam}(\Omega)}\Big)^{m} R^{2-n}|\Omega_{R}|, \qquad \forall\, x \in \Omega. \qquad (3.45)
\end{aligned}
$$

With (3.45) in hand, (3.43) follows from Lemma 2.12 (recall that $0 < \beta < 1/\alpha_{\mathfrak{G}}$), after some simple algebra (and using the fact that $|\Omega_{R}| \le |\Omega|$). The case $n = 2$ is similar. More specifically, the same type of argument as above yields the bound

$$
\begin{aligned}
\int_{\Omega} u(x)^{-\beta}\,dx \;\le\; & \; C(\theta,\beta)\Big(\frac{\operatorname{diam}(\Omega)}{R}\Big)^{(m+\alpha_{\mathfrak{G}})\beta}\Big(\frac{\mathcal{H}^{1}(\partial\Omega)}{R}\Big)^{\alpha_{\mathfrak{G}}\beta}\times \\
& \times\Big(\frac{|\Omega|}{R^{2}}\Big)^{\beta-\alpha_{\mathfrak{G}}\beta}\Big(\frac{R^{2}}{|\Omega_{R}|}\Big)^{\beta}|\Omega|^{1-\beta}\Big(\log\Big(\frac{\operatorname{diam}(\Omega)}{R}\Big)\Big)^{-\beta} \qquad (3.46)
\end{aligned}
$$

and, given that $\operatorname{diam}(\Omega)/R > 4$, the logarithmic factor can be bounded by $(\log 4)^{-\beta}$. This gives (3.44). $\qquad\square$

We continue by recording the following corollary.

Corollary 3.8. *If $\Omega \subseteq \mathbb{R}^n$, $n \ge 2$, is a bounded Lipschitz domain and if α_{Ω} is the critical exponent associated with Ω as in Definition 3.5, then the finiteness condition (1.6) holds granted that*

$$
0 < \beta < \frac{1}{\alpha_{\Omega}}. \qquad (3.47)
$$

In particular, (1.6) holds for any $\beta \in (0,1)$ in the case when Ω is a bounded C^1 domain.

Proof. The claim in the first part of the statement is an immediate consequence of our previous theorem, whereas (3.43) and (3.30) readily yield the claim in the second part of the statement. □

The principle emerging from Theorem 3.7 is that, for a bounded NTA domain $\Omega \subseteq \mathbb{R}^n$ with an Ahlfors regular boundary, the ratio

$$\left(\int_{\Omega} u(x)^{-\beta} \, dx \right) \Big/ |\Omega|^{1-2\beta/n} \tag{3.48}$$

can be controlled in terms of the proportion of the size of the cone (involved in cone condition (3.28)) relative to the size of the domain Ω itself (assuming that $\beta \in (0,1)$ relates favourably to the spherical profile of the cone).

An example of this principle at work in a concrete case of interest is as follows. For each $\kappa \in (0,1)$ and $N \geq 3$, denote by $\mathcal{P}(\kappa, N)$ the set of polygons with N sides inscribed in the unit circle S^1 and having the property that the ratio of any two sides belongs to the interval (κ, κ^{-1}). Theorem 3.7 then gives that, once $\kappa \in (0,1)$ and $\beta \in (0,1)$ have been fixed, there exist $N_{\beta,\kappa} \in \mathbb{N}$ and $C_{\beta,\kappa} > 0$ with the property that

$$\int_{\Omega} u(x)^{-\beta} \, dx \leq C_{\beta,\kappa}, \qquad \text{whenever } \Omega \in \mathcal{P}(\kappa, N) \text{ with } N \geq N_{\beta,\kappa}. \tag{3.49}$$

In particular, if Ω_N denotes the regular polygon with N sides ($N \in \mathbb{N}$, $N \geq 3$) inscribed in S^1 and u_N is the solution of (2.1) for $\Omega = \Omega_N$, then (3.49) gives that for every fixed $\beta \in (0,1)$ we have

$$\int_{\Omega_N} u_N(x)^{-\beta} \, dx = \mathcal{O}(1), \qquad \text{as } N \to \infty. \tag{3.50}$$

Improvements of (3.49)-(3.50) (vis-à-vis the range of β's, the shape of the polygon and the nature of the estimate for the β-integral) are presented in §4. For the time being, we wish to point out that (for any $0 < \beta < 1$)

$$\sup_{\Omega} \left(\int_{\Omega} u(x)^{-\beta} \, dx \right) \Big/ |\Omega|^{1-2\beta/n} = +\infty, \tag{3.51}$$

when the supremum is taken over all bounded convex sets in $\mathbb{R}^n$. With this goal in mind, for a fixed, small $\varepsilon > 0$, consider the thin rectangular domain

$$\Omega := \left\{ x = (x', x_n) \in \mathbb{R}^{n-1} \times \mathbb{R} : x' \in (0,1)^{n-1}, \ |x_n| < \varepsilon \right\} \subseteq \mathbb{R}^n, \tag{3.52}$$

and set

$$v(x', x_n) := \tfrac{1}{2}(\varepsilon^2 - x_n^2), \qquad \forall \, x = (x', x_n) \in \Omega. \tag{3.53}$$

Then $v \in C^0(\overline{\Omega})$, $-\Delta v = 1$ on Ω, and $v \geq 0$ on $\partial\Omega$. Therefore, if u solves (1.5) for Ω as in (3.52), we have $u(x) \leq v(x)$ for every $x \in \Omega$, on account of the Maximum Principle. As a result, for every $\beta \in (0,1)$ we may estimate

$$\begin{aligned} \int_{\Omega} u(x)^{-\beta} \, dx \ &\geq \ \int_{\Omega} v(x)^{-\beta} \, dx = 2^{\beta} \int_{-\varepsilon}^{\varepsilon} (\varepsilon^2 - x_n^2)^{-\beta} \, dx_n \\ &= \ 2^{1+\beta} \left(\int_0^1 (1 - t^2)^{-\beta} \, dt \right) \varepsilon^{1-2\beta} = C_{\beta} \, \varepsilon^{1-2\beta}, \\ &= \ C_{\beta} \, \varepsilon^{-2\beta(n-1)/n} \, |\Omega|^{1-2\beta/n}, \end{aligned} \tag{3.54}$$

from which (3.51) readily follows.

A natural end-point version of the estimate (1.6) is the 'weak-type' inequality

$$|\{x \in \Omega \,|\, u(x) < \lambda\}| \leq C(\Omega, \tilde{\beta})\lambda^{\tilde{\beta}} < +\infty \quad \text{for all } \lambda > 0. \tag{3.55}$$

The two conditions (1.6) and (3.55) are closely related, in that if (1.6) holds for some $\beta > 0$ then (3.55) holds for $0 < \tilde{\beta} \leq \beta$, and if (3.55) holds for some $\tilde{\beta} > 0$ then (1.6) holds for $0 < \beta < \tilde{\beta}$. These two statements follow easily from Chebyshev's inequality and the equality

$$\int_\Omega u(x)^{-\beta} dx = \beta \int_0^\infty \lambda^{-(\beta+1)} |\{x \in \Omega \,|\, u(x) < \lambda\}| d\lambda, \tag{3.56}$$

respectively.

In the case of a ball, the calculation in Remark 2.5 shows that (3.55) holds if and only if $\tilde{\beta} \leq 1$ and so, while (1.6) fails when $\beta = 1$ (see (2.14)), (3.55) holds when $\tilde{\beta} = 1$.

More general examples may be established via the same methods we have employed above. For example, an immediate consequence of (2.47) is

$$|\{x \in \Omega \,|\, \delta_\Omega(x)^{\tilde{\beta}} < \lambda\}| \leq C\lambda^{1/\tilde{\beta}} \mathcal{H}^{n-1}(\partial\Omega), \tag{3.57}$$

for each $\tilde{\beta} \geq 0$ and all $\lambda > 0$. This can be used to prove the following theorem.

Theorem 3.9. *(i) If $\Omega \subseteq \mathbb{R}^n$ is a bounded domain whose boundary is Ahlfors regular, then (3.55) holds provided that $0 < \tilde{\beta} \leq 1/2$.*

(ii) If $\Omega \subseteq \mathbb{R}^n$, $n \geq 2$, is a bounded Lipschitz domain and if α_Ω is the critical exponent associated with Ω as in Definition 3.5, then (3.55) holds provided that

$$0 < \beta \leq \frac{1}{\alpha_\Omega}. \tag{3.58}$$

Part (i) of Theorem 3.9 contains the appropriate end-point version of Proposition 2.13 (ii) corresponding to $\beta = 1/2$ and is proved using (3.57) and (2.58). Proposition 2.13 (iii) and the above discussion show this is sharp. Part (ii) corresponds to the end-point $\beta = 1/\alpha_\Omega$ of Corollary 3.8 and is proved again using (3.57) and, this time, (3.45).

We end the current section by recording a special case of Theorem 3.7 of independent interest. This requires that we first make the following definition.

Definition 3.10. *We say that an open set $\Omega \subseteq \mathbb{R}^n$ satisfies an inner ball condition with radius $r_0 \in (0, +\infty)$ provided*

$$\forall\, x \in \Omega \;\; \exists\, y \in \Omega \;\; \text{such that} \;\; x \in B(y, r_0) \subseteq \Omega. \tag{3.59}$$

In other words, an open set $\Omega \subseteq \mathbb{R}^n$ satisfies an inner ball condition with radius r_0 provided Ω can be written as the union of all balls of radius r_0 contained in Ω.

Obviously, an open set $\Omega \subseteq \mathbb{R}^n$ satisfying an inner ball condition with radius r_0 also satisfies an axially symmetric inner cone condition with any half-aperture $\theta \in (0, \pi/2)$ and any height $\leq 2r_0 \cos\theta$. This observation and Theorem 3.7 then readily yield the following corollary.

Corollary 3.11. *Assume that $\Omega \subseteq \mathbb{R}^2$ is a bounded NTA domain with an Ahlfors regular boundary, which satisfies an inner ball condition with radius r_0. Fix $\beta \in (0, 1)$ and select $\theta \in (\beta\pi/2, \pi/2)$. Then the solution u of (2.1) satisfies*

$$\begin{aligned} \int_\Omega u(x)^{-\beta}\, dx &\leq C_\Omega(\theta, \beta)\Big(\frac{\operatorname{diam}(\Omega)}{r_0}\Big)^{m\beta}\Big(\frac{\operatorname{diam}(\Omega)^2}{|\Omega_{r_0}|}\Big)|\Omega|^{1-\beta} \\ &\leq C_\Omega(\theta, \beta)\Big(\frac{\operatorname{diam}(\Omega)}{r_0}\Big)^{m\beta+2}|\Omega|^{1-\beta}, \end{aligned} \tag{3.60}$$

where $m > 0$ depends only on the NTA constants of Ω, and $C_\Omega(\theta, \beta) > 0$ is a finite constant which depends only on the Ahlfors character of $\partial\Omega$, the angle θ, and the parameter β.

In fact, a result similar in spirit holds in the case when $\Omega \subseteq \mathbb{R}^n$ with $n \geq 3$ as well.

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4   The case of a polygon in the plane

In this section we focus on the finiteness of the β-integral (cf. (1.6)) in the situation when Ω is a polygonal domain in $\mathbb{R}^2$. Some preparations are necessary. Given $\theta \in (0, \pi)$, consider the infinite sector

$$S_\theta := \{z \in \mathbb{C} = \mathbb{R}^2 : |\arg(z)| < \theta\} \tag{4.1}$$

and, for each $r > 0$, consider its truncated version

$$S_{\theta, r} = \{z \in S_\theta : |z| < r\}. \tag{4.2}$$

Hence, in polar coordinates $x = (\rho \cos \omega, \rho \sin \omega) \in \mathbb{R}^2$ with $(\rho, \omega) \in (0, \infty) \times (-\pi, \pi)$,

$$S_{\theta, r} = \{(\rho, \omega) : 0 < \rho < r \ \text{ and } \ -\theta < \omega < \theta\}. \tag{4.3}$$

One basic technical result in this section is contained in the next proposition below. To be able to formulate it, we will need the Gamma and Beta functions which, for the convenience of the reader, we now briefly recall. As is well-known, they are respectively given by

$$\Gamma(z) := \int_0^\infty t^{z-1} e^{-t} \, dt, \qquad z \in \mathbb{C}, \ \ \text{Re } z > 0, \tag{4.4}$$

and

$$\mathfrak{B}(z_1; z_2) := \int_0^1 t^{z_1 - 1}(1 - t)^{z_2 - 1} \, dt, \qquad z_j \in \mathbb{C}, \ \ \text{Re } z_j > 0, \ \ j = 1, 2. \tag{4.5}$$

with both integrals convergent under the specified conditions. It will also be useful to recall that an alternative formula for the Beta function is

$$\mathfrak{B}(z_1; z_2) = 2 \int_0^{\pi/2} (\sin \alpha)^{2z_1 - 1}(\cos \alpha)^{2z_2 - 1} \, d\alpha, \qquad z_j \in \mathbb{C}, \ \ \text{Re } z_j > 0, \ \ j = 1, 2. \tag{4.6}$$

Here is the proposition alluded to above.

Proposition 4.1. *Suppose that Ω is a bounded domain in $\mathbb{R}^2$ for which there exist $r > 0$ and $\theta \in (0, \pi)$ such that*

$$\Omega \cap B(0, r) = S_{\theta, r}. \tag{4.7}$$

Also, let u be the function defined by (2.1). Then for every $\beta \in (0, 1)$ there holds

$$\int_{S_{\theta, r}} u(x)^{-\beta} \, dx \leq C(\theta, \beta) \, r^{2(1 - \beta)}, \tag{4.8}$$

where $C(\theta, \beta) > 0$ is the finite constant described as

$$C(\theta, \beta) := \begin{cases} \left(\left(\frac{\pi}{2\theta} \right)^2 - 4 \right)^\beta \frac{(2\theta)^2}{\pi(\pi - 4\theta)} \mathfrak{B}\left(\frac{\theta(1 - \beta)}{\pi/4 - \theta}; 1 - \beta \right) \mathfrak{B}\left(\frac{1}{2}; \frac{1 - \beta}{2} \right) & \text{if } \theta \in \left(0, \frac{\pi}{4} \right), \\[2ex] 2^{3\beta - 2} \left(\frac{1}{1 - \beta} \right)^{1 - \beta} \Gamma(1 - \beta) \mathfrak{B}\left(\frac{1}{2}; \frac{1 - \beta}{2} \right) & \text{if } \theta = \frac{\pi}{4}, \\[2ex] \left(4 - \left(\frac{\pi}{2\theta} \right)^2 \right)^\beta \frac{(2\theta)^2}{\pi(4\theta - \pi)} \mathfrak{B}\left(\frac{\theta - \beta\pi/2}{\theta - \pi/4}; 1 - \beta \right) \mathfrak{B}\left(\frac{1}{2}; \frac{1 - \beta}{2} \right) & \text{if } \theta \in \left(\frac{\pi}{4}, \pi \right]. \end{cases} \tag{4.9}$$

As a consequence, for every $\beta \in (0, 1)$,

$$\int_{S_{\theta, r}} u(x)^{-\beta} \, dx \leq C_\beta \, \theta^{1 - 2\beta} \, r^{2(1 - \beta)}. \tag{4.10}$$

The bound in (4.10) is in the nature of best possible, in the sense that if $\Omega := S_{\theta,r}$ for some $\theta \in (0,\pi)$ and $r > 0$, and if u is the solution of (2.1) for this domain then, in fact, for every $\beta \in (0,1)$,

$$\int_{S_{\theta,r}} u(x)^{-\beta}\, dx \approx \theta^{1-2\beta}\, r^{2(1-\beta)}, \qquad \text{uniformly for } \theta \in (0,\pi) \text{ and } r > 0, \tag{4.11}$$

with comparability constants which depend exclusively on β.

To get a better feel for the constant defined in (4.9), a few comments are in order. Since, as is well-known, we have the following asymptotic formula

$$\mathfrak{B}(x,y) \sim \Gamma(y)\, x^{-y} \quad \text{when } x > 0 \text{ is large, for each fixed } y > 0, \tag{4.12}$$

we deduce from this and (4.9) that, for each $\beta \in (0,1)$ fixed,

$$C(\theta,\beta) \sim \pi^{-1}(1-\beta)^{\beta-1}\theta^{1-\beta}(\pi+4\theta)^{\beta}\,\Gamma(1-\beta)\mathfrak{B}\Big(\frac{1}{2};\frac{1-\beta}{2}\Big) \quad \text{for } \theta \text{ close to } \tfrac{\pi}{4}. \tag{4.13}$$

In particular, this shows that the functions $\theta \mapsto C(\theta,\beta)$ from (4.9) are continuous at $\pi/4$. In this vein, let us also remark here that since

$$\lim_{x \to 0^+} \Big(x\mathfrak{B}(x,y)\Big) = 1 \qquad \text{for each fixed } y > 0, \tag{4.14}$$

it follows from (4.9) that, for each $\beta \in (0,1)$ fixed,

$$C(\theta,\beta) \sim \theta^{1-2\beta}, \qquad \text{for } \theta \text{ close to zero.} \tag{4.15}$$

Altogether, the above analysis shows that for each fixed $\beta \in (0,1)$, the quantity $C(\theta,\beta)$ depends continuously on $\theta \in (0,\pi]$ and satisfies

$$C(\theta,\beta) \approx \theta^{1-2\beta}, \qquad \text{uniformly for } \theta \in (0,\pi]. \tag{4.16}$$

The proof of the above proposition requires further preparations. For each $\theta \in (0,\pi)$ and $r > 0$ consider the barrier function $v_{\theta,r} : S_{\theta,r} \to \mathbb{R}$ which, in polar coordinates (ρ,ω), is given by

$$v_{\theta,r}(\rho,\omega) := \begin{cases} \Big(\big(\tfrac{\pi}{2\theta}\big)^2 - 4\Big)^{-1} r^2 \Big[\big(\tfrac{\rho}{r}\big)^2 - \big(\tfrac{\rho}{r}\big)^{\pi/(2\theta)}\Big] \cos\big(\tfrac{\pi\omega}{2\theta}\big) & \text{if } \theta \in (0,\pi) \setminus \{\tfrac{\pi}{4}\}, \\[2mm] \tfrac{\rho^2}{4}\log\big(\tfrac{r}{\rho}\big)\cos(2\omega) & \text{if } \theta = \tfrac{\pi}{4}. \end{cases} \tag{4.17}$$

It is reassuring to observe that, because of the differentiation quotient present in (4.17), which can be highlighted by writing

$$
\begin{aligned}
v_{\theta,r}(\rho,\omega) &= \Big(\big(\tfrac{\pi}{2\theta}\big)^2 - 4\Big)^{-1} r^2 \Big[\big(\tfrac{\rho}{r}\big)^2 - \big(\tfrac{\rho}{r}\big)^{\pi/(2\theta)}\Big] \cos\big(\tfrac{\pi\omega}{2\theta}\big) \\[2mm]
&= -r^2\Big(\tfrac{\pi}{2\theta}+2\Big)^{-1} \frac{\big(\tfrac{\rho}{r}\big)^2 - \big(\tfrac{\rho}{r}\big)^{\pi/(2\theta)}}{2 - \tfrac{\pi}{2\theta}} \cos\big(\tfrac{\pi\omega}{2\theta}\big),
\end{aligned}
\tag{4.18}$$

the formula for $v_{\theta,r}$ corresponding to the special value $\theta = \pi/4$ (i.e., second line of (4.17)) is the natural limit case of the formula for $v_{\theta,r}$ in the first line of (4.17) as $\theta \to \pi/4$.

The above barrier function has been designed precisely as to satisfy, for any $\theta \in (0,\pi)$ and any $r > 0$

$$
\begin{aligned}
&v_{\theta,r} = 0 \ \text{ on } \ \partial S_{\theta,r}, \qquad v_{\theta,r} > 0 \ \text{ in } \ S_{\theta,r}, \quad \text{and} \\[2mm]
&-(\Delta v_{\theta,r})(\rho,\omega) = \cos\big(\pi\omega/(2\theta)\big) \ \text{ in } \ S_{\theta,r}.
\end{aligned}
\tag{4.19}$$

The last property is verified by means of an elementary calculation based on the fact that, in polar coordinates in the plane, the Laplacian can be written as $\Delta = d^2/d\rho^2 + \rho^{-1}d/d\rho + \rho^{-2}d^2/d\omega^2$. The normalisation constants in (4.17) have been selected so that the right-hand side in the second line of (4.19) is precisely a cosine (this will be of relevance shortly; cf. (4.30) below). While checking the last formula in (4.19) it also helps to notice that the function

$$\widetilde{v}_\theta(\rho,\omega) := \rho^{\pi/(2\theta)} \cos\left(\frac{\pi\omega}{2\theta}\right) = \mathrm{Im}\,[z^{\pi/(2\theta)}], \qquad z = \rho\,e^{i\omega}, \tag{4.20}$$

is harmonic for every $\theta \in (0,\pi)$ and every $r > 0$. The primary reason for which we prefer the more intricate barrier function $v_{\theta,r}$ from (4.17) over the simpler barrier function $\widetilde{v}_\theta$ from (4.20) is that for small values of θ (more precisely, for $\theta \in (0,\pi/4)$) the former exhibits a better behaviour than the latter near the origin. This aspect plays a crucial role in our subsequent analysis.

With the above notation and conventions we have:

Lemma 4.2. *For any $\beta \in (0,1)$, any $r > 0$ and any $\theta \in (0,\pi)$, we have*

$$\int_{S_{\theta,r}} v_{\theta,r}(x)^{-\beta}\,dx = C(\theta,\beta)\,r^{2(1-\beta)} \tag{4.21}$$

where $C(\theta,\beta) > 0$ is the finite constant given in (4.9).

Proof. Fix $\beta \in (0,1)$, let $r > 0$ be arbitrary and first assume that $\theta \in (0,\pi/4)$. We have

$$\begin{aligned}
\int_{S_{\theta,r}} v_{\theta,r}(x)^{-\beta}\,dx &= \left(\left(\frac{\pi}{2\theta}\right)^2 - 4\right)^\beta r^{-2\beta} \int_{S_{\theta,r}} \left[\left(\frac{\rho}{r}\right)^2 - \left(\frac{\rho}{r}\right)^{\pi/(2\theta)}\right]^{-\beta} \left[\cos\left(\frac{\pi\omega}{2\theta}\right)\right]^{-\beta}dx \\
&= \left(\left(\frac{\pi}{2\theta}\right)^2 - 4\right)^\beta r^{-2\beta} I \cdot II, \tag{4.22}
\end{aligned}$$

where

$$I := \int_0^r \left[\left(\frac{\rho}{r}\right)^2 - \left(\frac{\rho}{r}\right)^{\pi/(2\theta)}\right]^{-\beta}\rho\,d\rho \quad \text{and} \quad II := \int_{-\theta}^\theta \left[\cos\left(\frac{\pi\omega}{2\theta}\right)\right]^{-\beta}d\omega. \tag{4.23}$$

Making two changes of variables, first introducing $t := \rho/r$, then substituting s for t^ε, where $\varepsilon := \pi/(2\theta) - 2 > 0$, yields

$$\begin{aligned}
I &= \varepsilon^{-1}r^2 \int_0^1 s^{2(1-\beta)/\varepsilon-1}(1-s)^{-\beta}\,ds = \frac{r^2}{\varepsilon}\mathfrak{B}\left(\frac{2(1-\beta)}{\varepsilon};1-\beta\right) \\
&= r^2\left(\frac{\pi}{2\theta}-2\right)^{-1}\mathfrak{B}\left(\frac{\theta(1-\beta)}{\pi/4-\theta};1-\beta\right), \tag{4.24}
\end{aligned}$$

after unravelling notation. On the other hand, making the change of variables $\alpha := \pi\omega/(2\theta)$ and using the parity of the cosine function permits one to write (after a reference to (4.6))

$$II = \frac{4\theta}{\pi}\int_0^{\pi/2}(\cos\alpha)^{-\beta}\,d\alpha = \frac{2\theta}{\pi}\mathfrak{B}\left(\frac{1}{2};\frac{1-\beta}{2}\right). \tag{4.25}$$

Collectively, (4.22)-(4.25) establish the validity of (4.21) in the case when $\theta \in (0,\pi/4)$ and with a constant $C(\theta,\beta)$ as in the first line in the right-hand side of (4.9). The case when $\theta \in (\pi/4,\pi)$ is treated in a most analogous manner and we omit it.

Finally, corresponding to $\theta = \pi/4$, we have

$$\begin{aligned}
\int_{S_{\pi/4,r}} v_{\pi/4,r}(x)^{-\beta}\,dx &= 4^\beta \int_{S_{\pi/4,r}} \rho^{-2\beta}\left[\log\left(\frac{r}{\rho}\right)\right]^{-\beta}\left[\cos\left(2\omega\right)\right]^{-\beta}dx \\
&= 4^\beta r^{2-2\beta}III \cdot IIV, \tag{4.26}
\end{aligned}$$

where, after the changes of variables $t = \rho/r$ and $\alpha = 2\omega$,

$$III := \int_0^1 t^{1-2\beta} \left[\log\left(\frac{1}{t}\right) \right]^{-\beta} dt \quad \text{and} \quad IV := 2^{-1} \int_{-\pi/2}^{\pi/2} (\cos \alpha)^{-\beta} d\alpha. \tag{4.27}$$

One more change of variables, substituting $\exp(-s/(2-2\beta))$ for t in III, transforms this term into

$$III = \left(\frac{1}{2-2\beta}\right)^{1-\beta} \int_0^\infty e^{-s} s^{-\beta} \, ds = 2^{\beta-1} \left(\frac{1}{1-\beta}\right)^{1-\beta} \Gamma(1-\beta), \tag{4.28}$$

whereas, much as before,

$$IV = 2^{-1} \mathfrak{B}\left(\frac{1}{2}; \frac{1-\beta}{2}\right). \tag{4.29}$$

In concert, (4.26)-(4.29) justify (4.21) in the case when $\theta = \pi/4$, with a constant $C(\theta, \beta)$ as in the middle line in the right-hand side of (4.9). This completes the proof of the lemma. $\square$

After this preamble, here is the end-game in the following proof.

Proof of Proposition 4.1. Let $r > 0$, $\theta \in (0, \pi)$ be as in the statement of the proposition. With $v_{\theta,r}$ as in (4.17), it follows that the function $w := u - v_{\theta,r} : S_{\theta,r} \to \mathbb{R}$ satisfies (recall that we are assuming that $\Delta u = -1$ in $\Omega \supseteq S_{\theta,r}$)

$$(\Delta w)(\rho, \omega) = -1 + \cos\left(\pi\omega/(2\theta)\right) \leq 0 \quad \text{for each} \quad \rho\, e^{i\omega} \in S_{\theta,r}, \tag{4.30}$$

i.e., w is superharmonic in $S_{\theta,r}$. In addition, w is continuous in $\overline{S_{\theta,r}}$ and

$$w\big|_{\partial S_{\theta,r}} = u\big|_{\partial S_{\theta,r}} \geq 0 \tag{4.31}$$

given that, by design, u is nonnegative in $\overline{\Omega}$. Hence, the Maximum Principle applies and yields $w \geq 0$ in $S_{\theta,r}$ or, in other words,

$$u \geq v_{\theta,r} \quad \text{in} \quad S_{\theta,r}. \tag{4.32}$$

The estimate (4.8) now readily follows by combining (4.32) with the result proved in Lemma 4.2. With this in hand, (4.10) is a direct consequence of (4.8) and (4.15) (cf. also the claim following this last equation).

Finally, there remains to prove the equivalence in (4.11). Of course, the left-pointing inequality is contained in (4.10), so we only need to check the right-pointing inequality in (4.11). To this end, suppose in what follows that $\Omega = S_{\theta,r}$ for some $\theta \in (0, \pi)$, $r > 0$, and that $u = u_{\theta,r}$ solves (2.1) for this particular domain. To continue, fix $\theta_0 \in (0, \pi/2)$ and assume first that $\theta \in (\theta_0, \pi)$. Then

$$\int_{S_{\theta,r}} u_{\theta,r}(x)^{-\beta} \, dx = r^{2(1-\beta)} \int_{S_{\theta,1}} u_{\theta,1}(x)^{-\beta} \, dx \geq C_\beta \, \theta_0 \, r^{2(1-\beta)}, \tag{4.33}$$

because the way the first integral scales in the parameter r, and (2.7). Since in the situation we are currently considering $\theta^{1-2\beta}$ behaves like a constant, the desired conclusion follows in this case. We are left with considering the case when $\theta \in (0, \theta_0)$. In such a scenario, set

$$\widetilde{S}_{\theta,r} := \{(x,y) \in \mathbb{R}^2 : 0 < y < x \tan\theta \ \text{and} \ 0 < x < r\}. \tag{4.34}$$

and note that

$$\widetilde{S}_{\theta,r\cos\theta} \subseteq S_{\theta,r} \subseteq \widetilde{S}_{\theta,r}. \tag{4.35}$$

Next, consider the following barrier function, designed to befit the triangular region $\widetilde{S}_{\theta,r}$:

$$\widetilde{v}_{\theta,r} : \widetilde{S}_{\theta,r} \longrightarrow \mathbb{R}, \quad \widetilde{v}_{\theta,r}(x,y) := \tfrac{1}{2}\, y\left(x\tan\theta - y\right), \quad \text{for all} \ (x,y) = x + iy \in \widetilde{S}_{\theta,r}. \tag{4.36}$$

Hence, $\widetilde{v}_{\theta,r}$ is continuous in $\overline{\widetilde{S}_{\theta,r}}$ and satisfies $\Delta\widetilde{v}_{\theta,r} = -1$ in $\widetilde{S}_{\theta,r}$ as well as $\widetilde{v}_{\theta,r} \geq 0$ on $\partial\widetilde{S}_{\theta,r}$. Consequently, the function $\widetilde{v}_{\theta,r} - u$ is harmonic in $S_{\theta,r}(=\Omega)$, continuous on its closure, and $\widetilde{v}_{\theta,r} - u = \widetilde{v}_{\theta,r} \geq 0$ on $\partial S_{\theta,r}(=\partial\Omega)$. Therefore, by the Maximum Principle, $u \leq \widetilde{v}_{\theta,r}$ in $S_{\theta,r}$ which, in turn, gives that

$$
\begin{aligned}
\int_{S_{\theta,r}} u(x)^{-\beta}\, dx \;\; & \geq \;\; \int_{S_{\theta,r}} \widetilde{v}_{\theta,r}(x)^{-\beta}\, dx \geq \int_{\widetilde{S}_{\theta,r\cos\theta}} \widetilde{v}_{\theta,r}(x)^{-\beta}\, dx \\
& = \;\; 2^{-\beta}\int_0^{r\cos\theta}\Big(\int_0^{x\tan\theta} y^{-\beta}(x\tan\theta - y)^{-\beta}\, dy\Big)dx \\
& = \;\; 2^{-\beta}(\tan\theta)^{1-2\beta}\Big(\int_0^{r\cos\theta} x^{1-2\beta}\, dx\Big)\Big(\int_0^1 t^{-\beta}(1-t)^{-\beta}\, dt\Big), \\
& = \;\; C_\beta\,(\tan\theta)^{1-2\beta}(\cos\theta)^{2-2\beta}r^{2-2\beta} = C_\beta\,(\cos\theta)(\sin\theta)^{1-2\beta}r^{2-2\beta} \\
& \approx \;\; \theta^{1-2\beta}r^{2-2\beta}, \quad \text{uniformly for } \theta \in (0,\theta_0), \quad\quad\quad (4.37)
\end{aligned}
$$

(taking into account (4.35) in the second inequality and after making the change of variables $y = t\,x\tan\theta$ in the inner integral in the second line). This concludes the justification of (4.11) and finishes the proof of the proposition. $\qquad\square$

We are now in a position to formulate the first main result in this section. This shows that, whenever Ω is a polygon, (4.8) holds for every $\beta < 1$ which is remarkable since, as opposed to the situation discussed in Corollary 3.8, this time $\partial\Omega$ is far from being regular.

Theorem 4.3. *Assume that Ω is a polygon in $\mathbb{R}^2$ and that u is the solution of* (2.1). *Then, for every $\beta \in (0,1)$,*

$$
\int_\Omega u(x)^{-\beta}\, dx \leq C(\Omega,\beta) < +\infty. \quad\quad\quad (4.38)
$$

Proof. Let $\{P_1, ..., P_N\}$ be the vertices of the polygon Ω and, for each $i \in \{1, ..., N\}$, denote by $\theta_i \in (0,\pi)$ the half-measure of the angle corresponding to P_i, and by L'_i and L''_i the lengths of the two sides of Ω emerging from P_i. Also, for each $i \in \{1, ..., N\}$, introduce

$$
r_i := \min\Big\{L'_i,\, L''_i,\, (L'_i)^{1/(2-2\beta)}\theta_i^{(2\beta-1)/(2-2\beta)},\, (L''_i)^{1/(2-2\beta)}\theta_i^{(2\beta-1)/(2-2\beta)}\Big\}. \quad\quad\quad (4.39)
$$

Parenthetically, we wish to note that

$$
\beta = 1/2 \implies r_i = \min\{L'_i, L''_i\}. \quad\quad\quad (4.40)
$$

Writing a formula similar to (4.10) at each vertex (note that the problem (2.1) transforms naturally under rigid motions of the plane) and summing up all contributions obtained from integrating $u^{-\beta}$ near each vertex gives

$$
\begin{aligned}
\sum_{i=1}^N \int_{\Omega\cap B(P_i,r_i)} u(x)^{-\beta}\, dx \;\; & \leq \;\; C_\beta \sum_{i=1}^N \theta_i^{1-2\beta} r_i^{2(1-\beta)} \\
& \leq \;\; C_\beta \sum_{i=1}^N \min\{L'_i, L''_i\} \leq C_\beta\,\mathcal{H}^1(\partial\Omega), \quad\quad\quad (4.41)
\end{aligned}
$$

with a finite constant $C_\beta > 0$ which depends only on β.

Having estimated the contribution from the vertices, construct now a C^1 domain $\Omega_* \subseteq \Omega$ by rounding off each vertex P_i with a suitably small circular arc contained within $B(P_i, r_i/2)$. This ensures that Ω_β has the property that

$$
\Omega_* \cup \Big(\bigcup_{1\leq i\leq N}(\Omega\cap B(P_i,r_i))\Big) = \Omega. \quad\quad\quad (4.42)
$$

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Next, consider u_* such that

$$\begin{cases} \Delta u_* = -1 \;\; \text{in} \;\; \Omega_*, \\ u_* = 0 \;\; \text{on} \;\; \partial\Omega_*, \\ u_* \in C^0(\overline{\Omega_*}). \end{cases} \tag{4.43}$$

Since $u - u_*$ is a continuous function in $\overline{\Omega_*}$ which satisfies $\Delta(u - u_*) = 0$ in Ω_β and $u - u_* = u \geq 0$ on $\partial\Omega_* \subseteq \overline{\Omega}$, it follows from the Maximum Principle that $u \geq u_*$ in Ω_*. Thus, for every $\beta \in (0,1)$,

$$\int_{\Omega_*} u(x)^{-\beta}\, dx \leq \int_{\Omega_*} u_*(x)^{-\beta}\, dx \leq C(\Omega_*, \beta) < +\infty, \tag{4.44}$$

by virtue of Corollary 3.8. Now, (4.38) follows readily from (4.41), (4.42) and (4.44). $\quad\square$

In the class of convex polygons in $\mathbb{R}^2$ it is possible to further clarify the nature of the constant $C(\Omega, \beta)$ in (4.38). This is done in Theorem 4.7, stated later in this section. As a preamble, a key technical result used in the proof of this theorem is isolated in the proposition below.

Proposition 4.4. *Assume that Ω is a bounded domain in $\mathbb{R}^2$ for which there exist $r > 0$ and $\theta \in (0, \pi/2)$ with the property that*

$$\Omega \cap B(0,r) = S_{\theta,r} \quad\text{and}\quad B\big((r(\cos\theta)^{-1}, 0), r\tan\theta\big) \subseteq \Omega. \tag{4.45}$$

As usual, let u be the function defined by (2.1). Then

$$0 < \beta < \min\left\{1, \tfrac{4\theta}{\pi}\right\} \implies \int_{S_{\theta,r}} u(x)^{-\beta}\, dx \leq C_\beta \left(\frac{\theta^{2(1-\beta)}}{4\theta - \pi\beta}\right) r^{2(1-\beta)} (\cos\theta)^\beta, \tag{4.46}$$

where $C_\beta > 0$ is the finite constant which depends only on β.

Furthermore, retaining (4.45) it follows that for every $\beta \in (0,1)$ there holds

$$\int_{S_{\theta,r}} u(x)^{-\beta}\, dx \leq C_\beta\, \theta^{1-2\beta}\, r^{2(1-\beta)} (\cos\theta)^\beta. \tag{4.47}$$

Proof. Consider the function defined by

$$v(x) := \tfrac{1}{4}\left[r^2 \tan^2\theta - \left(x_1 - r(\cos\theta)^{-1}\right)^2 - x_2^2 \right], \qquad x = (x_1, x_2). \tag{4.48}$$

In polar coordinates $x = (x_1, x_2) = (\rho\cos\omega, \rho\sin\omega)$ this takes the form

$$v(\rho, \omega) = \tfrac{1}{4}\left[2r\rho\cos\omega(\cos\theta)^{-1} - r^2 - \rho^2 \right]. \tag{4.49}$$

Let us also define the harmonic function

$$w(\rho, \omega) := \rho^{\pi/(2\theta)} \cos\left(\frac{\pi\omega}{2\theta}\right), \qquad \rho > 0, \; |\omega| < \theta. \tag{4.50}$$

Then, w vanishes on the straight sides of $\partial S_{\theta,r}$ (i.e., for $\omega = \pm\theta$), while on the rounded portion of the boundary of $\partial S_{\theta,r}$ (i.e., the arc described by $\rho = r$ and $|\omega| < \theta$) we have

$$\frac{w(r, \omega)}{v(r, \omega)} = 2(\cos\theta)\, r^{\pi/(2\theta)-2} \frac{\cos\left(\frac{\pi\omega}{2\theta}\right)}{\cos\omega - \cos\theta}, \qquad |\omega| < \theta. \tag{4.51}$$

Note that

$$\sup_{|\omega|<\theta} \left(\frac{\cos\left(\frac{\pi\omega}{2\theta}\right)}{\cos\omega - \cos\theta}\right) = \sup_{0<\omega<\theta} \left(\frac{\cos\left(\frac{\pi\omega}{2\theta}\right)}{\cos\omega - \cos\theta}\right) \tag{4.52}$$

and, for each $\omega \in (0, \theta)$,

$$\frac{\cos\left(\frac{\pi\omega}{2\theta}\right)}{\cos\omega - \cos\theta} = \frac{\sin\left(\frac{\pi}{2} - \frac{\pi\omega}{2\theta}\right)}{\cos\omega - \cos\theta} = \frac{\sin\left(\frac{\pi(\theta-\omega)}{2\theta}\right)}{\cos\omega - \cos\theta}$$

$$= \frac{\pi}{2\theta}\left(\frac{\sin\left(\frac{\pi(\theta-\omega)}{2\theta}\right)}{\frac{\pi(\theta-\omega)}{2\theta}}\right)\frac{\theta-\omega}{\cos\omega - \cos\theta}. \tag{4.53}$$

Since $\pi(\theta - \omega)/(2\theta) \in (0, \pi/2)$ whenever $\omega \in (0, \theta)$, it follows that the fraction in parentheses in the right-most expression in (4.53) is ≤ 1 for each $\omega \in (0, \theta)$. Also, elementary calculus shows that there exists a universal constant $c > 0$ such that $(\theta - \omega)/(\cos\omega - \cos\theta) \leq c/\theta$ if $0 < \omega < \theta$ (recall that $\theta \in (0, \pi/2)$). Consequently,

$$\frac{w(r, \omega)}{v(r, \omega)} \leq c\,\theta^{-2}(\cos\theta)\,r^{\pi/(2\theta)-2}, \qquad \text{for every } \omega \in (-\theta, \theta). \tag{4.54}$$

Next, given that by design

$$v = 0 \text{ on } \partial B\big((r(\cos\theta)^{-1}, 0), r\tan\theta\big), \quad \Delta v = -1 \text{ in } B\big((r(\cos\theta)^{-1}, 0), r\tan\theta\big), \tag{4.55}$$

and that, by assumption, $B\big((r(\cos\theta)^{-1}, 0), r\tan\theta\big) \subseteq \Omega$, the Maximum Principle ensures that

$$v \leq u \text{ in } B\big((r(\cos\theta)^{-1}, 0), r\tan\theta\big). \tag{4.56}$$

In particular,

$$v(r, \omega) \leq u(r, \omega) \text{ for all } \omega \in (-\theta, \theta). \tag{4.57}$$

From this and (4.54) we may therefore conclude that

$$c^{-1}\,\theta^2(\cos\theta)^{-1}\,r^{2-\pi/(2\theta)}w(r, \omega) \leq u(r, \omega) \qquad \text{for every } \omega \in (-\theta, \theta). \tag{4.58}$$

Granted this, as well as the properties of w recorded just after (4.50), the Maximum Principle applies again and yields that

$$c^{-1}\,\theta^2(\cos\theta)^{-1}\,r^{2-\pi/(2\theta)}w(x) \leq u(x) \qquad \text{for every } x \in S_{\theta,r}. \tag{4.59}$$

Hence, for every $\beta \in (0, 4\theta/\pi)$, a familiar (by now) computation gives that

$$\int_{S_{\theta,r}} u(x)^{-\beta}\,dx \leq c^\beta\theta^{-2\beta}(\cos\theta)^\beta\,r^{(\pi/(2\theta)-2)\beta}\int_{S_{\theta,r}} w(x)^{-\beta}\,dx$$

$$= c^\beta\theta^{-2\beta}(\cos\theta)^\beta\,r^{(\pi/(2\theta)-2)\beta}\,\mathfrak{B}\left(\frac{1}{2}; \frac{1-\beta}{2}\right)\left(\frac{4\theta^2}{\pi(4\theta - \pi\beta)}\right)r^{2-\pi\beta/(2\theta)}$$

$$= C_\beta\left(\frac{\theta^{2(1-\beta)}}{4\theta - \pi\beta}\right)r^{2(1-\beta)}(\cos\theta)^\beta, \tag{4.60}$$

proving (4.46).

Finally, when $\theta \in (\pi/3, \pi/2)$, it follows from (4.46) that for every $\beta \in (0, 1)$ we have

$$\int_{S_{\theta,r}} u(x)^{-\beta}\,dx \leq C_\beta\,r^{2(1-\beta)}(\cos\theta)^\beta, \tag{4.61}$$

which further implies (4.47) in the case we are considering. When $\theta \in (0, \pi/3)$, then (4.47) is a direct consequence of (4.10). $\qquad\square$

Continuing the buildup to Theorem 4.7, we now make several definitions and comment on their significance and how they interrelate.

Definition 4.5. *The eccentricity of an open, bounded convex set Ω in $\mathbb{R}^n$ is defined as*

$$\mathrm{ecc}\,(\Omega) := \frac{\inf\{R_1 > 0 : \exists\, x \in \mathbb{R}^n \quad such\ that\ \ \Omega \subseteq B(x, R_1)\}}{\sup\{R_2 > 0 : \exists\, x \in \mathbb{R}^n \quad such\ that\ \ B(x, R_2) \subseteq \Omega\}}. \tag{4.62}$$

It follows that

$$\begin{aligned} &\mathrm{ecc}\,(\Omega)\ \ \text{controls both the NTA constants of } \Omega \\ &\text{as well as the Ahlfors character of } \partial\Omega, \end{aligned} \tag{4.63}$$

uniformly in the class of open, bounded and convex subsets Ω of $\mathbb{R}^n$. Furthermore, there exists a dimensional constant c_n with the property that for every open, bounded convex set $\Omega \subseteq \mathbb{R}^n$ we have

$$\mathrm{diam}\,(\Omega) \leq c_n\,\mathrm{ecc}\,(\Omega)\,|\Omega|^{1/n}. \tag{4.64}$$

Definition 4.6. *Let Ω be a convex polygon in $\mathbb{R}^2$. Call $R > 0$ an admissible radius for Ω provided for each side of Ω there exists a ball of radius R contained in Ω which is tangent to that side. Then define the maximal admissible radius of Ω as*

$$R_\Omega := \sup\{R > 0 : R\ \ is\ an\ admissible\ radius\ for\ \Omega\}, \tag{4.65}$$

and set

$$\Omega^\# := \bigcup_{x \in \Omega,\, \delta_\Omega(x) > R_\Omega} B(x, R_\Omega). \tag{4.66}$$

Straight from definitions it can be seen that

$$\Omega \subseteq \mathbb{R}^2\ \ \text{convex polygon} \implies \Omega^\#\ \ \text{satisfies an inner ball condition of radius } R_\Omega. \tag{4.67}$$

Also, it is not too difficult to show that

$$\begin{aligned} \Omega \subseteq \mathbb{R}^2\ \ \text{convex polygon} \implies &\Omega^\#\ \ \text{is a } C^{1,1}\ \text{convex domain} \\ &\text{satisfying ecc}\,(\Omega^\#) \leq \mathrm{ecc}\,(\Omega). \end{aligned} \tag{4.68}$$

Shortly we will also need the readily verified claim that

$$\begin{aligned} \Omega \subseteq \mathbb{R}^2\ \ \text{convex polygon} \implies &\text{the angles of } \Omega \text{ are} \geq 2\theta_*, \text{ where} \\ &\theta_* \in (0, \pi/2) \text{ depends only on ecc}\,(\Omega). \end{aligned} \tag{4.69}$$

After this prelude, we are now prepared to state and prove a refined version of (4.38) in the class of convex polygons in $\mathbb{R}^2$.

Theorem 4.7. *Assume that Ω is a convex polygon in $\mathbb{R}^2$. Then for every $\beta \in (0, 1)$ the solution u of (2.1) satisfies*

$$\int_\Omega u(x)^{-\beta}\,dx \leq C(\mathrm{ecc}\,(\Omega), \beta)\Big(\frac{\mathrm{diam}\,(\Omega)}{R_\Omega}\Big)^{m\beta+2}|\Omega|^{1-\beta}, \tag{4.70}$$

where $m > 0$ depends only on $\mathrm{ecc}\,(\Omega)$.

Proof. Denote by $\{P_1, ..., P_N\}$ the vertices of the polygon Ω and, for each $i \in \{1, ..., N\}$, let $\theta_i \in (0, \pi/2)$ be the half-measure of the angle corresponding to P_i. Set

$$r_i := R_\Omega \cos\theta_i, \qquad 1 \leq i \leq N, \tag{4.71}$$

and fix $\beta \in (0,1)$. Also, let $\theta_* \in (0, \pi/2)$ be as in (4.69). Then, thanks to (4.47), for each $i \in \{1, ..., N\}$ we have

$$\int_{\Omega \cap B(P_i, r_i)} u(x)^{-\beta} \, dx \leq C(\theta_*, \beta) \, r_i^{2(1-\beta)} (\cos \theta_i)^{\beta} = C(\theta_*, \beta) \, R_{\Omega}^{-\beta} r_i^{2-\beta}. \tag{4.72}$$

Consequently,

$$\sum_{i=1}^{N} \int_{\Omega \cap B(P_i, r_i)} u(x)^{-\beta} \, dx \quad \leq \quad C(\theta_*, \beta) \, R_{\Omega}^{-\beta} \sum_{i=1}^{N} r_i^{2-\beta}$$

$$\leq \quad C(\theta_*, \beta) \, R_{\Omega}^{-\beta} \Big(\sum_{i=1}^{N} r_i \Big)^{2-\beta}, \tag{4.73}$$

since $2 - \beta > 1$. Given that $\sum_{i=1}^{N} r_i$ is controlled by the perimeter of Ω, this yields

$$\sum_{i=1}^{N} \int_{\Omega \cap B(P_i, r_i)} u(x)^{-\beta} \, dx \quad \leq \quad C(\theta_*, \beta) \, R_{\Omega}^{-\beta} \big[\mathcal{H}^1(\partial \Omega) \big]^{2-\beta}$$

$$\leq \quad C(\theta_*, \beta) \, R_{\Omega}^{-\beta} (\operatorname{diam}(\Omega))^{2-\beta}$$

$$= \quad C(\theta_*, \beta) \Big(\frac{\operatorname{diam}(\Omega)}{R_{\Omega}} \Big)^{\beta} (\operatorname{diam}(\Omega))^{2(1-\beta)}. \tag{4.74}$$

In light of (4.70) (and keeping (4.69) in mind), this bound suits our purposes.

To continue, we note that

$$\Omega^{\#} \cup \Big(\bigcup_{i=1}^{N} (\Omega \cap B(P_i, r_i)) \Big) = \Omega, \tag{4.75}$$

and observe that, thanks to (4.67) and (3.60),

$$\int_{\Omega^{\#}} u(x)^{-\beta} \, dx \quad \leq \quad C_{\Omega^{\#}}(\beta) \Big(\frac{\operatorname{diam}(\Omega^{\#})}{R_{\Omega}} \Big)^{m\beta+2} |\Omega^{\#}|^{1-\beta}$$

$$\leq \quad C_{\Omega^{\#}}(\beta) \Big(\frac{\operatorname{diam}(\Omega)}{R_{\Omega}} \Big)^{m\beta+2} (\operatorname{diam}(\Omega))^{2(1-\beta)}, \tag{4.76}$$

where $m > 0$ depends only on the NTA constants of $\Omega^{\#}$, and $C_{\Omega^{\#}}(\beta) > 0$ is a finite constant which depends only on the Ahlfors character of $\partial \Omega^{\#}$ and β. Hence, by (4.63) and (4.68), $C_{\Omega^{\#}}(\beta)$ can be controlled in terms of $\operatorname{ecc}(\Omega)$ and β. Estimate (4.70) now follows from this observation, (4.74), (4.75), (4.76) and (4.64). $\qquad \square$

Remark 4.8. *In regard to (4.70), it should be pointed out that, in the class of convex polygons in $\mathbb{R}^2$, the maximal admissible radius cannot be controlled in terms of the diameter and the eccentricity. A simple example is as follows. Let Ω be the triangle whose vertices have coordinates $(-1, 0)$, $(1, 0)$, $(0, 1)$ and, for each $j \geq 2$, consider the convex quadrilateral $\Omega_j := \{(x, y) \in \Omega : y < 1 - 1/j\}$. It is then clear that while $\operatorname{ecc}(\Omega_j)$ and $\operatorname{diam}(\Omega_j)$ stay bounded, $R_{\Omega_j} \to 0$ as $j \to \infty$.*

We conclude this section by giving an asymptotic formula for the β-integral of a regular polygon, as the number of vertices increases. This augments earlier estimates in (3.49)-(3.50).

Proposition 4.9. *For each $N \in \mathbb{N}$, $N \geq 3$, let $\Omega_N \subseteq \mathbb{R}^2$ denote the regular polygon with N sides, circumscribed by $B(0, 1)$. Denote by u_N the solution u of (2.1) when $\Omega = \Omega_N$. Then for each $\beta \in (0, 1)$ the following asymptotic formula holds*

$$\int_{\Omega_N} u_N(x)^{-\beta} \, dx = \frac{4^{\beta} \pi}{1 - \beta} + \mathcal{O}(N^{\beta-1}) \qquad as \;\; N \to \infty. \tag{4.77}$$

Proof. We specialise part of the proof of Theorem 4.7 to the present case. In the current setting, using notation introduced on that occasion, we have:

$$R_{\Omega_N} = 1, \quad \theta_i = \frac{\pi}{2} - \frac{\pi}{N}, \quad r_i = R_{\Omega_N} \cos\theta_i = \sin(\pi/N), \quad 1 \leq i \leq N. \tag{4.78}$$

Also, as before, we let $P_1, ..., P_N$ be the vertices of Ω_N. Hence, the first inequality in (4.73) gives

$$\sum_{i=1}^{N} \int_{\Omega_N \cap B(P_i, r_i)} u_N(x)^{-\beta}\, dx \leq C_\beta \sum_{i=1}^{N} \left(\sin(\pi/N)\right)^{2-\beta} \leq C_\beta\, N^{\beta-1}. \tag{4.79}$$

Since the Maximum Principle and (2.12) imply $u(x) \geq \frac{1}{4}(1 - |x|^2)$ for all $x \in B(0,1)$, we therefore obtain the asymptotic estimate

$$
\begin{aligned}
\int_{\Omega_N} u_N(x)^{-\beta}\, dx &\leq \sum_{i=1}^{N} \int_{\Omega_N \cap B(P_i, r_i)} u_N(x)^{-\beta}\, dx + \int_{B(0,1)} u_N(x)^{-\beta}\, dx \\
&\leq 4^\beta \int_{B(0,1)} \frac{dx}{(1-|x|^2)^\beta} + \mathcal{O}(N^{\beta-1}) \\
&= \frac{4^\beta \pi}{1-\beta} + \mathcal{O}(N^{\beta-1}).
\end{aligned}
\tag{4.80}
$$

On the other hand, (2.12) and the Maximum Principle give

$$u_N(x) \leq \tfrac{1}{4}\left((\cos(\pi/N))^{-2} - |x|^2\right) \quad \text{for all } x \in \Omega_N, \tag{4.81}$$

which then forces

$$
\begin{aligned}
\int_{\Omega_N} u_N(x)^{-\beta}\, dx &\geq 4^\beta \int_{B(0,1)} \frac{dx}{\left((\cos(\pi/N))^{-2} - |x|^2\right)^\beta} = 4^\beta \pi \int_0^1 \frac{dt}{\left((\cos(\pi/N))^{-2} - t\right)^\beta} \\
&= \frac{4^\beta \pi}{1-\beta}\left((\cos(\pi/N))^{-2(1-\beta)} - \left((\cos(\pi/N))^{-2} - 1\right)^{1-\beta}\right) \\
&= \frac{4^\beta \pi}{1-\beta}(\cos(\pi/N))^{-2(1-\beta)}\left(1 - \left(\sin(\pi/N)\right)^{2(1-\beta)}\right) \\
&= \frac{4^\beta \pi}{1-\beta}\left(1 - \mathcal{O}(N^{2(\beta-1)})\right).
\end{aligned}
\tag{4.82}
$$

Now, (4.77) follows from (4.80) and (4.82). $\qquad\square$

5 Piecewise smooth domains with conical singularities in $\mathbb{R}^n$

In this section we shall work in the general n-dimensional case, $n \geq 2$. Throughout, we retain notation introduced in §3. Given $r > 0$ and an open connected subset $\mathfrak{G}$ of S^{n-1} with a sufficiently regular boundary (relative to S^{n-1}), define the truncated cone

$$S_{\mathfrak{G},r} := B(0,r) \cap \Gamma_{\mathfrak{G}} = \{(\rho, \omega) : 0 < \rho < r, \ \omega \in \mathfrak{G}\}, \tag{5.1}$$

where $(\rho, \omega) \in (0, \infty) \times S^{n-1}$ are the standard polar coordinates in $\mathbb{R}^n$. Associated with this truncated cone, consider the barrier function $v_{\mathfrak{G},r} : S_{\mathfrak{G},r} \to \mathbb{R}$ which, in polar coordinates is given by

$$
v_{\mathfrak{G},r}(\rho, \omega) := \begin{cases} \dfrac{r^2}{\Lambda_{\mathfrak{G}} - 2n}\left[\left(\dfrac{\rho}{r}\right)^2 - \left(\dfrac{\rho}{r}\right)^{\alpha_{\mathfrak{G}}}\right]\phi_{\mathfrak{G}}(\omega) & \text{if } \Lambda_{\mathfrak{G}} \neq 2n, \\[4mm] \dfrac{\rho^2}{n+2} \log\left(\dfrac{r}{\rho}\right)\phi_{\mathfrak{G}}(\omega) & \text{if } \Lambda_{\mathfrak{G}} = 2n, \end{cases}
\tag{5.2}
$$

for each $\omega \in \mathfrak{G}$ and $\rho \in (0, r)$. Given that, by (3.8), we have

$$\Lambda_{\mathfrak{G}} - 2n = (\alpha_{\mathfrak{G}} - 2)(\alpha_{\mathfrak{G}} + n), \tag{5.3}$$

it is worth noting that

$$\Lambda_{\mathfrak{G}} = 2n \iff \alpha_{\mathfrak{G}} = 2. \tag{5.4}$$

In particular, the formula for $v_{\mathfrak{G},r}(\rho, \omega)$ in the second line of (5.2) is the limiting case of the formula for $v_{\mathfrak{G},r}(\rho, \omega)$ in the first line of (5.2) as $\Lambda_{\mathfrak{G}}$ becomes $2n$. Much as before, in axially symmetric case, i.e., for $\mathfrak{G} = S^{n-1} \cap \Gamma_\theta$ for some $\theta \in (0, \pi)$, we agree to abbreviate $S_{S^{n-1} \cap \Gamma_\theta, r}$ and $v_{S^{n-1} \cap \Gamma_\theta, r}$ by $S_{\theta,r}$ and $v_{\theta,r}$, respectively. In this scenario, we therefore have

$$v_{\theta,r}(\rho, \omega) = \begin{cases} \frac{r^2}{\Lambda_\theta - 2n}\left[\left(\frac{\rho}{r}\right)^2 - \left(\frac{\rho}{r}\right)^{\alpha_\theta}\right]\phi_\theta(\omega), & \text{if } \theta \neq \theta_n, \\[3mm] \frac{\rho^2}{n+2}\log\left(\frac{r}{\rho}\right)\phi_{\theta_n}(\omega), & \text{corresponding to } \theta = \theta_n, \end{cases} \tag{5.5}$$

where $\theta_n \in (0, \pi)$ is the unique angle for which $\alpha_{\theta_n} = 2$. Note that since the assignment $\theta \mapsto \alpha_\theta$ is strictly decreasing (cf. (3.14)) and since $\alpha_{\pi/2} = 1$ and $\alpha_\theta \nearrow +\infty$ as $\theta \searrow 0$ (cf. (3.15)), there exists precisely one angle $\theta_n \in (0, \pi/2)$ for which $\alpha_{\theta_n} = 2$. In fact, from (3.19) we know that

$$\theta_n = \arccos\left(1/\sqrt{n}\right), \qquad n \geq 2, \tag{5.6}$$

so that, in particular,

$$\theta_n = \begin{cases} \frac{\pi}{4} & \text{when } n = 2, \\[2mm] \frac{\pi}{3} & \text{when } n = 4, \end{cases} \qquad \text{and} \quad \theta_n \nearrow \frac{\pi}{2} \quad \text{as} \quad n \to \infty. \tag{5.7}$$

This discussion shows that $\Lambda_\theta \neq 2n$ if $\theta \neq \theta_n$ and, hence, $v_{\theta,r}(\rho, \omega)$ is well-defined for every $\theta \in (0, \pi)$ and $v_{\theta_n, r}(\rho, \omega)$ is the limit of $v_{\theta,r}(\rho, \omega)$ as $\theta \to \theta_n$.

Lemma 5.1. *Assume that $\mathfrak{G}$ is an open, connected subset of S^{n-1} whose relative boundary is a C^1 submanifold of codimension one in S^{n-1}. Then for every $\beta \in (0, 1)$ and $r > 0$, the barrier function $v_{\mathfrak{G},r}$ from (5.2) satisfies*

$$\int_{S_{\mathfrak{G},r}} v_{\mathfrak{G},r}(x)^{-\beta}\, dx = c_n(\mathfrak{G}, \beta)\, r^{n-2\beta}\left(\int_{\mathfrak{G}} \phi_{\mathfrak{G}}(\omega)^{-\beta}\, d\omega\right) < +\infty, \tag{5.8}$$

where

$$c_n(\mathfrak{G}, \beta) := \begin{cases} \frac{(\Lambda_{\mathfrak{G}} - 2n)^\beta}{\alpha_{\mathfrak{G}} - 2}\mathfrak{B}\left(\frac{n-2\beta}{\alpha_{\mathfrak{G}} - 2}; 1 - \beta\right), & \text{if } 2 < \alpha_{\mathfrak{G}} < +\infty, \\[3mm] (n - 2\beta)^{\beta - 1}(n + 2)^\beta\Gamma(1 - \beta), & \text{if } \alpha_{\mathfrak{G}} = 2, \\[3mm] \frac{(2n - \Lambda_{\mathfrak{G}})^\beta}{2 - \alpha_{\mathfrak{G}}}\mathfrak{B}\left(\frac{n-2\beta}{2 - \alpha_{\mathfrak{G}}}; 1 - \beta\right) & \text{if } 0 < \alpha_{\mathfrak{G}} < 2. \end{cases} \tag{5.9}$$

Proof. The proof largely parallels that of Lemma 4.2. We include it primarily to indicate how the right-hand side of (5.8) shapes up. Fix $\beta \in (0, 1)$, let $r > 0$ be arbitrary and first assume that $\alpha_{\mathfrak{G}} \in (2, \infty)$. This forces $\Lambda_{\mathfrak{G}} > 2n$ and we have

$$\begin{aligned} \int_{S_{\mathfrak{G},r}} v_{\mathfrak{G},r}(x)^{-\beta}\, dx &= r^{-2\beta}(\Lambda_{\mathfrak{G}} - 2n)^\beta\int_{S_{\mathfrak{G},r}}\left[\left(\frac{\rho}{r}\right)^2 - \left(\frac{\rho}{r}\right)^{\alpha_{\mathfrak{G}}}\right]^{-\beta}\phi_{\mathfrak{G}}(\omega)^{-\beta}\, dx \\[2mm] &= r^{-2\beta}(\Lambda_{\mathfrak{G}} - 2n)^\beta I \cdot II, \end{aligned} \tag{5.10}$$

where

$$I := \int_0^r \left[\left(\frac{\rho}{r} \right)^2 - \left(\frac{\rho}{r} \right)^{\alpha_\theta} \right]^{-\beta} \rho^{n-1} \, d\rho \quad \text{and} \quad II := \int_{\mathfrak{G}} \phi_{\mathfrak{G}}(\omega)^{-\beta} \, d\omega. \tag{5.11}$$

As in the past, we make two changes of variables, first letting $t := \rho/r$, then replacing t^ε by s where, this time, we set $\varepsilon := \alpha_{\mathfrak{G}} - 2 > 0$. This yields

$$\begin{aligned} I &= \varepsilon^{-1} r^n \int_0^1 s^{(n-2\beta)/\varepsilon - 1} (1-s)^{-\beta} \, ds = \frac{r^n}{\varepsilon} \mathfrak{B} \left(\frac{n-2\beta}{\varepsilon}; 1-\beta \right) \\ &= \frac{r^n}{\alpha_{\mathfrak{G}} - 2} \mathfrak{B} \left(\frac{n-2\beta}{\alpha_{\mathfrak{G}} - 2}; 1-\beta \right), \end{aligned} \tag{5.12}$$

Thus, (5.10)-(5.12) prove (5.8) in the case when $\alpha_{\mathfrak{G}} \in (2, \infty)$ and with a constant $c_n(\mathfrak{G}, \beta)$ as in the first line in the right-hand side of (5.9). The case when $\alpha_{\mathfrak{G}} \in (0, 2)$ is treated similarly and we omit it. Moving on, in the case corresponding to $\alpha_{\mathfrak{G}} = 2$ we write

$$\begin{aligned} \int_{S_{\mathfrak{G}, r}} v_{\mathfrak{G}, r}(x)^{-\beta} \, dx &= (n+2)^\beta \int_{S_{\mathfrak{G}, r}} \rho^{-2\beta} \left[\log \left(\frac{r}{\rho} \right) \right]^{-\beta} \phi_{\mathfrak{G}}(\omega)^{-\beta} \, dx \\ &= (n+2)^\beta r^{2-2\beta} III \cdot IV, \end{aligned} \tag{5.13}$$

where (after a natural change of variables)

$$III := \int_0^1 t^{n-1-2\beta} \left[\log \left(\frac{1}{t} \right) \right]^{-\beta} \, dt \quad \text{and} \quad IV := \int_{\mathfrak{G}} \phi_{\mathfrak{G}}(\omega)^{-\beta} \, d\omega. \tag{5.14}$$

Substituting $\exp(-s/(n-2\beta))$ for t in III, further transforms this term into

$$III = (n-2\beta)^{\beta-1} \int_0^\infty e^{-s} s^{-\beta} \, ds = (n-2\beta)^{\beta-1} \Gamma(1-\beta). \tag{5.15}$$

Together, (5.13)-(5.15) justify (5.8) in the case when $\alpha_{\mathfrak{G}} = 2$, with a constant $c_n(\mathfrak{G}, \beta)$ as in the middle line in the right-hand side of (5.9).

The last thing left to justify, in order to complete the proof of the lemma, is the finiteness condition in (5.8). This, however, is a direct consequence of (3.5) (cf. also Lemma 2.12). $\square$

The main result in this section is the following higher-dimensional analogue of Theorem 4.3. Essentially, this asserts that (1.6) holds for every $\beta \in (0, 1)$ in the class of bounded piecewise C^1 domains in $\mathbb{R}^n$ with conical singularities. While this constitutes a subclass of the larger class of bounded Lipschitz domains, it is worth recalling that Theorem 3.7 establishes (1.6) only for a smaller range of values for the parameter β (described in (3.47)). This is surprising especially since direct pointwise estimates (from below) on the solution u of (2.1) indicate that the regularity of u at a given conical singularity deteriorates as the aperture of the cone tends to zero. In this vein, note that the conclusion in (the first part of) Corollary 3.8 progressively weakens precisely when the Lipschitz constant of a domain becomes large (cf. (3.31)). Thus, heuristically speaking, for the type of domains considered in this section, the size of the reciprocal of the solution of the Saint Venant problem (2.1) exhibits a milder character (than perhaps originally suspected) in the process of taking an integral average.

Theorem 5.2. *Assume that Ω is a bounded open set in $\mathbb{R}^n$, $n \geq 2$, whose boundary is of class C^1 with the exception of finitely many points $P_1, ..., P_N \in \partial\Omega$, and such that for each $i \in \{1, ..., N\}$ there exist an open, connected subset $\mathfrak{G}_i$ of S^{n-1} whose relative boundary is a C^1 submanifold of codimension one in S^{n-1} and $r_i > 0$ with the property that*

$$\Omega \cap B(P_i, r_i) \text{ and } S_{\mathfrak{G}_i, r_i} \text{ coincide, modulo a rigid transformation of } \mathbb{R}^n. \tag{5.16}$$

Let u be the solution of (2.1). Then for every $\beta \in (0, 1)$ there holds

$$\int_\Omega u(x)^{-\beta} \, dx \leq C(\Omega, \beta) < +\infty. \tag{5.17}$$

Proof. Pick $P \in \partial\Omega$ with the property that there exist an open, connected subset $\mathfrak{G}$ of S^{n-1} whose relative boundary is a C^1 submanifold of codimension one in S^{n-1} and $r > 0$ such that $\Omega \cap B(P, r)$ and $S_{\mathfrak{G}, r}$ are congruent. Without loss of generality, assume that P is the origin in $\mathbb{R}^n$ and that, in fact, $\Omega \cap B(0, r) = S_{\mathfrak{G}, r}$. Bring in the barrier function $v_{\mathfrak{G}, r}$ from (5.2) and note that, by design,

$$v_{\mathfrak{G}, r} = 0 \ \text{ on } \ \partial S_{\mathfrak{G}, r}, \qquad v_{\mathfrak{G}, r} > 0 \ \text{ in } \ S_{\mathfrak{G}, r},$$
$$\text{and } \ (\Delta v_{\mathfrak{G}, r})(\rho, \omega) = -\phi_{\mathfrak{G}}(\omega) \ \text{ in } \ S_{\mathfrak{G}, r}. \tag{5.18}$$

Indeed, these are direct consequences of (5.2), (3.10), (3.6), (3.8), and (3.4). The normalisation constants in (5.5) have been selected so that the right-hand side in the second line of (5.18) is always a number belonging to the interval $[-1, 0]$ (for this, (3.4) is crucial). As such, $\Delta(u - v_{\mathfrak{G}, r})(\rho, \omega) = -1 + \phi_{\mathfrak{G}}(\omega) \leq 0$ for any $\rho\omega \in S_{\mathfrak{G}, r}$. Since $u - v_{\mathfrak{G}, r}$ is continuous in $\overline{S_{\mathfrak{G}, r}}$ and is equal to (the nonnegative function) u on $\partial S_{\mathfrak{G}, r}$, the Maximum Principle gives that

$$u \geq v_{\mathfrak{G}, r} \ \text{ in } \ S_{\mathfrak{G}, r}. \tag{5.19}$$

In turn, this and (5.8) permit us to estimate

$$\int_{\Omega \cap B(P, r)} u(x)^{-\beta} \, dx \leq c_n(\mathfrak{G}, \beta) \, r^{n-2\beta} \left(\int_{\mathfrak{G}} \phi_{\mathfrak{G}}(\omega)^{-\beta} \, d\omega \right) < +\infty, \tag{5.20}$$

where the constant $c_n(\mathfrak{G}, \beta)$ is as in (5.9). Once this local estimate near a conical point $P \in \partial\Omega$ has been established, the remainder of the proof follows along the lines of the proof of Theorem 4.3. $\qquad\square$

6 Results for other classes of nonsmooth domains

In this section we study the nature of β-integrals associated with other important classes of non-smooth domains, starting with

6.1 The case of polyhedral domains

Consider the case when $\Omega \subseteq \mathbb{R}^3$ is a polyhedral domain[3]. Pick a vertex $x_0 \in \partial\Omega$ and, for a sufficiently small $\varepsilon > 0$, set

$$\mathfrak{G} := \big\{ (x - x_0)/|x - x_0| : x \in B(x_0, \varepsilon) \cap \Omega \big\} \subseteq S^2. \tag{6.1}$$

Hence, the spherical polygon $\mathfrak{G}$ is the profile of the cone which agrees with $\partial\Omega$ is a small neighbourhood of x_0. In this setting, a good portion of our earlier analysis carries through verbatim. In particular, we may consider the eigenvalue problem (3.3) which continues to have a solution which satisfies (3.4). The key feature which is lost in the present setting (in which $\mathfrak{G}$ no longer has a smooth boundary in S^2) is the equivalence (3.5). Recall that this played a basic role in the finiteness condition in (5.8). We nonetheless have the following result.

Theorem 6.1. *Suppose that u is the solution of the Saint Venant problem (2.1) in the case when Ω is a polyhedral domain in $\mathbb{R}^3$. Then*

$$\int_{\Omega} u(x)^{-\beta} \, dx \leq C(\Omega, \beta) < +\infty \tag{6.2}$$

for every $\beta \in (0, 1)$.

---

[3]Throughout, a polyhedral domain is understood to have finitely many faces, edges and vertices.

The remainder of this section is devoted to presenting a proof of this result. Dealing with the first eigenfunction $\phi_\mathfrak{G}$ in the case when $\mathfrak{G}$ is a spherical polyhedral domain requires the following asymptotic representation of $\phi_\mathfrak{G}$ near a corner point of the spherical polygon $\mathfrak{G}$ (in the spirit of work in [21])

$$\phi_\mathfrak{G}(\psi, \varphi) = C_\mathfrak{G} (\sin \varphi)^{\pi/(2\theta)} \cos(\pi\psi/(2\theta)) + \mathcal{O}\big((\sin \varphi)^{\pi/(2\theta)+\varepsilon}\big), \tag{6.3}$$

for some $\varepsilon > 0$, where $C_\mathfrak{G}$ is a constant depending on the global shape of $\mathfrak{G}$ and ψ, φ are local polar coordinates near a corner vertex $O \in \partial_{S^2}\mathfrak{G}$, i.e., $0 < \varphi \ll 1$, $|\psi| < \theta$ where θ is the half-aperture of the spherical angle at O. In addition, the coefficient $C_\mathfrak{G}$ in (6.3) is given by the following formula, itself a special case of closely related results proved in [28],

$$C_\mathfrak{G} = \Lambda_\mathfrak{G} \int_\mathfrak{G} \phi_\mathfrak{G}(\omega)\zeta(\omega)\,d\omega. \tag{6.4}$$

Above, ζ is a positive function in $\mathfrak{G}$, harmonic (in the sense of Laplace-Beltrami) in $\mathfrak{G}$ and vanishing on $\partial_{S^2}\mathfrak{G} \setminus \{O\}$, and which exhibits a prescribed singularity at the vertex O, namely

$$\zeta(\psi, \varphi) \sim (2/\pi)(\sin \varphi)^{-\pi/(2\theta)} \cos(\pi\psi/(2\theta)), \quad \text{uniformly for } |\psi| < \theta, \text{ as } \varphi \searrow 0. \tag{6.5}$$

Together with (6.4) and the fact that $\phi_\mathfrak{G} > 0$, this analysis shows that $C_\mathfrak{G} > 0$.

Unfortunately, formula (6.3) is not sufficiently refined in order to allow us to estimate

$$\int_\mathfrak{G} \phi_\mathfrak{G}(\omega)^{-\beta}\,d\omega. \tag{6.6}$$

An asymptotic expansion of $\phi_\mathfrak{G}$ near a corner $O \in \partial_{S^2}\mathfrak{G}$ which better suits our purposes is contained in the lemma below.

Lemma 6.2. *The remainder in the asymptotic formula*

$$\phi_\mathfrak{G}(\psi, \varphi) = C_\mathfrak{G}(\sin \varphi)^{\pi/(2\theta)} \cos\big(\pi\psi/(2\theta)\big) + A(\psi, \varphi), \tag{6.7}$$

where $0 < \varphi \ll 1$ and $|\psi| < \theta$, obeys the estimate

$$|A(\psi, \varphi)| \leq C(\sin \varphi)^{\pi/(2\theta)+\varepsilon} \cos(\pi\psi/(2\theta)), \tag{6.8}$$

for some $\varepsilon > 0$.

Proof. Without loss of generality assume that the direction of the edge inducing the spherical angle with opening 2θ is along the x_3-axis and set

$$x = (x', x_3), \quad x' = (x_1, x_2) = (\rho \sin \varphi \cos \psi, \rho \sin \varphi \sin \psi),$$
$$x_3 = \rho \cos \varphi, \quad \rho = \sqrt{x_1^2 + x_2^2 + x_3^2}. \tag{6.9}$$

Let $\alpha_\mathfrak{G} > 0$ solve (3.8) so that, in particular, the function described in polar coordinates by $\rho^{\alpha_\mathfrak{G}}\phi_\mathfrak{G}(\psi, \varphi)$ is harmonic. Consequently,

$$0 = \Delta_x\big(\rho^{\alpha_\mathfrak{G}}\phi_\mathfrak{G}(\psi, \varphi)\big) = C_\mathfrak{G}\Delta_{x'}\big(\rho^{\alpha_\mathfrak{G}-\pi/(2\theta)}(\rho \sin \varphi)^{\pi/(2\theta)} \cos(\pi\psi/(2\theta))\big)$$
$$+ \Delta_x\big(\rho^{\alpha_\mathfrak{G}} A(\psi, \varphi)\big). \tag{6.10}$$

Given that

$$\Delta_{x'}\big((\rho \sin \varphi)^{\pi/(2\theta)} \cos(\pi\psi/(2\theta))\big) = 0, \tag{6.11}$$

we can express the last term in the first line of (6.10) as

$$C_\mathfrak{G}\Delta_{x'}\big(\rho^{\alpha_\mathfrak{G}-\pi/(2\theta)}\big) \cdot (\rho \sin \varphi)^{\pi/(2\theta)} \cos(\pi\psi/(2\theta))$$
$$+ 2C_\mathfrak{G}\nabla_{x'}\big(\rho^{\alpha_\mathfrak{G}-\pi/(2\theta)}\big) \cdot \nabla_{x'}\big((\rho \sin \varphi)^{\pi/(2\theta)} \cos(\pi\psi/(2\theta))\big). \tag{6.12}$$

Fix $\delta \in (0,1)$ small enough and restrict ρ to the interval $[1-\delta, 1]$. On this range, we have $|\nabla_{x'}\rho| \leq C \sin\varphi$ and $|\Delta_{x'}\rho| \leq C$, where the constant C depends only on δ. As a result, the absolute value of the expression in (6.12) does not exceed

$$C(\sin\varphi)^{\pi/(2\theta)}. \tag{6.13}$$

Using this and (6.10) we may then write

$$\Upsilon := \Delta_x\big(\rho^{\alpha_{\mathfrak{G}}}A(\psi,\varphi)\big) = \mathcal{O}\Big((\sin\varphi)^{\pi/(2\theta)}\Big). \tag{6.14}$$

Next, since $\rho^{\alpha_{\mathfrak{G}}}A(\psi,\varphi)$ vanishes on the sides of the dihedral angle $\psi = \pm\theta$, it follows from the classical local regularity result of Agmon, Douglas and Nirenberg (cf. [1]) that

$$\sum_{0 \leq |\gamma| \leq 2} \delta^{|\gamma|-2}\|\partial_x^\gamma(\rho^{\alpha_{\mathfrak{G}}}A)\|_{L^p(\mathcal{E}_\delta)} \leq C\Big(\|\Upsilon\|_{L^p(\widetilde{\mathcal{E}}_\delta)} + \delta^{3/p-2}\max_{\widetilde{\mathcal{E}}_\delta}|\rho^{\alpha_{\mathfrak{G}}}A|\Big), \tag{6.15}$$

where $p \in (1,\infty)$ is fixed and we have set

$$\begin{aligned}
\mathcal{E}_\delta &:= \big\{x \in \mathbb{R}^3 : 1-\delta \leq \rho \leq 1, \ |\psi| \leq \theta, \ \delta \leq \sin\varphi \leq 2\delta\big\}, \\
\widetilde{\mathcal{E}}_\delta &:= \big\{x \in \mathbb{R}^3 : 1-2\delta \leq \rho \leq 1+\delta, \ |\psi| \leq \theta, \ \delta/2 \leq \sin\varphi \leq 4\delta\big\}.
\end{aligned} \tag{6.16}$$

On account of this and the Sobolev embedding theorem we therefore obtain for $p > 3$

$$\delta^{3/p-1}\max_{\mathcal{E}_\delta}|\nabla_x(\rho^{\alpha_{\mathfrak{G}}}A)| \leq C\Big(\|\Upsilon\|_{L^p(\widetilde{\mathcal{E}}_\delta)} + \delta^{3/p-2}\max_{\widetilde{\mathcal{E}}_\delta}|\rho^{\alpha_{\mathfrak{G}}}A|\Big). \tag{6.17}$$

On the other hand, recall that (cf. (6.14) and (6.3))

$$\Upsilon = \mathcal{O}\Big((\sin\varphi)^{\pi/(2\theta)}\Big) \quad \text{and} \quad A = \mathcal{O}\Big((\sin\varphi)^{\pi/(2\theta)+\varepsilon}\Big). \tag{6.18}$$

In concert with (6.17) this yields

$$\delta^{3/p-1}\max_{\mathcal{E}_\delta}|\nabla_x(\rho^{\alpha_{\mathfrak{G}}}A)| \leq C\Big(\delta^{3/p+\pi/(2\theta)} + \delta^{3/p-2+\pi/(2\theta)+\varepsilon}\Big) \tag{6.19}$$

so that, ultimately,

$$\max_{\mathcal{E}_\delta}|\nabla_x(\rho^{\alpha_{\mathfrak{G}}}A)| \leq C\delta^{\pi/(2\theta)+\varepsilon-1}. \tag{6.20}$$

Using this and the fact that $\rho^{\alpha_{\mathfrak{G}}}A$ vanishes on the sides of the dihedral angle $\psi = \pm\theta$ we then obtain

$$\begin{aligned}
\rho^{\alpha_{\mathfrak{G}}}|A(\psi,\varphi)| &\leq C\cos(\pi\psi/(2\theta)) \cdot \delta \cdot \max_{\mathcal{E}_\delta}|\nabla_x(\rho^{\alpha_{\mathfrak{G}}}A)| \\
&\leq C\delta^{\pi/(2\theta)+\varepsilon}\cos(\pi\psi/(2\theta)).
\end{aligned} \tag{6.21}$$

This proves that (6.8) holds for small values of φ. $\qquad\square$

Corollary 6.3. *Retaining notation introduced above we have*

$$\phi_{\mathfrak{G}}(\psi,\varphi) = \Big(C_{\mathfrak{G}} + \mathcal{O}\big((\sin\varphi)^\varepsilon\big)\Big)(\sin\varphi)^{\pi/(2\theta)}\cos(\pi\psi/(2\theta)), \tag{6.22}$$

as $\varphi \to 0$. In addition,

$$\phi_{\mathfrak{G}}(x) \geq C\operatorname{dist}_{S^2}\big(x,\partial_{S^2}\mathfrak{G}\big), \quad \text{uniformly for } x \in \mathfrak{G} \text{ away from the vertices of } \mathfrak{G}. \tag{6.23}$$

Proof. Formula (6.22) is simply a re-writing of (6.7)-(6.8), whereas formula (6.23) follows from a simple barrier argument and the Maximum Principle. $\qquad\square$

Let O be a point on one of the edges such that its distance from the nearest vertex of our polyhedron is $r > 0$. In such a scenario, $\mathfrak{G}$ equals the diangle on the unit sphere

$$\mathfrak{G} = \{\omega = (\psi, \varphi) : 0 < \varphi < \pi, \, |\psi| < \theta\}. \tag{6.24}$$

In this case, the normalised eigenfunction (in spherical polar coordinates) is

$$\phi_{\mathfrak{G}}(\omega) = (\sin \varphi)^{\pi/(2\theta)} \cos(\pi\psi/(2\theta)), \qquad \omega = (\psi, \varphi). \tag{6.25}$$

After this preamble, we are ready to present the

Proof of Theorem 6.1. Let $\mathfrak{G} \subseteq S^2$ be the spherical polygon from (6.1) and consider the truncated cone $S_{\mathfrak{G},1}$ (cf. (5.1)) with edges meeting at a generic vertex $O \in \mathbb{R}^3$ of Ω. Without loss of generality, assume that O is the origin in $\mathbb{R}^3$. Throughout, $\alpha_{\mathfrak{G}}$ retains its earlier significance, and we denote the openings of the dihedral angles of $S_{\mathfrak{G},1}$ by $2\theta_j$, $1 \leq j \leq N$. We wish to show that $u^{-\beta}$ is integrable near O, and divide the subsequent analysis into several cases, starting with:

Case (i): Assume that the vertex O is such that $\alpha_{\mathfrak{G}} > 2$ (our analysis also applies to "fictitious vertices", i.e., for points of edges with $2\theta_j < \pi/2$). Let $\Phi_{\mathfrak{G}}$ be the (unique) variational solution of the Dirichlet problem

$$-\Delta_{S^2} \Phi_{\mathfrak{G}}(\omega) - 6\,\Phi_{\mathfrak{G}}(\omega) = 1 \quad \text{on} \quad \mathfrak{G}, \qquad \Phi_{\mathfrak{G}}\big|_{\partial_{S^2}\mathfrak{G}} = 0. \tag{6.26}$$

Since $6 < \alpha_{\mathfrak{G}}(\alpha_{\mathfrak{G}} + 1)$ in the current case, it follows from the Maximum Principle that $\Phi_{\mathfrak{G}} > 0$ on $\mathfrak{G}$. Furthermore, there exists $\delta > 0$ small with the property that the following asymptotic representations hold for $\Phi_{\mathfrak{G}}$ near the j-th angle vertex on $\partial_{S^2}\mathfrak{G}$ (cf. [21] and [28] for closely related results):

$$\Phi_{\mathfrak{G}}(\omega) = \begin{cases} \frac{\varphi_j^2}{4}(1 + \mathcal{O}(\varphi_j^\delta))\left(\frac{\cos(2\psi_j)}{\cos(2\theta_j)} - 1\right), & \text{if } \theta_j < \pi/4, \\[2mm] \varphi_j^2(1 + \mathcal{O}(\varphi_j^\delta))\left(\log\frac{1}{\varphi_j}\right)\cos(2\psi_j), & \text{if } \theta_j = \pi/4, \\[2mm] C_j \varphi_j^{\pi/2\theta_j}(1 + \mathcal{O}(\varphi_j^\delta))\cos\left(\frac{\pi\psi_j}{2\theta_j}\right), & \text{if } \theta_j > \pi/4, \end{cases} \tag{6.27}$$

where $C_j > 0$ and (φ_j, ψ_j) are the polar coordinates of the point $\omega \in S^2$, near the j-th angle vertex of $\partial_{S^2}\mathfrak{G}$. In turn, this yields

$$\int_{\mathfrak{G}} \frac{d\omega}{\Phi_{\mathfrak{G}}(\omega)^\beta} < +\infty \quad \text{whenever} \quad \beta < 1. \tag{6.28}$$

To proceed, recall that, in general, for any function w in $\mathbb{R}^3$ we have (with $\rho := |x|$)

$$\Delta w = \rho^{-2}(\rho^2 w_\rho)_\rho + \rho^{-2}\Delta_{S^2} w. \tag{6.29}$$

Hence, if we now introduce the function

$$v(x) := \left(|x|^2 - |x|^{\alpha_{\mathfrak{G}}}\right)\Phi_{\mathfrak{G}}\left(\frac{x}{|x|}\right) \quad \text{on} \quad S_{\mathfrak{G},1}, \tag{6.30}$$

it follows that

$$-\Delta v(x) = 1 + \left(\alpha_{\mathfrak{G}}(\alpha_{\mathfrak{G}} + 1) - 6\right)\left(\frac{x}{|x|}\right)^{\alpha_{\mathfrak{G}}-2}\Phi_{\mathfrak{G}}\left(\frac{x}{|x|}\right) - |x|^{\alpha_{\mathfrak{G}}-2} \quad \text{on } S_{\mathfrak{G},1} \tag{6.31}$$

$$v\big|_{\partial S_{\mathfrak{G},1}} = 0. \tag{6.32}$$

Since, by assumption, $\alpha_{\mathfrak{G}} > 2$, the right-hand side of (6.31) bounded by some positive finite constant c_0. Based on this and the Maximum Principle we may then conclude that for a sufficiently small $r > 0$ there holds

$$u(x) \geq \frac{1}{c_0}r^2 v\left(\frac{x}{r}\right) \quad \text{on} \quad S_{\mathfrak{G},r}. \tag{6.33}$$

Now, the fact that $u^{-\beta} \in L^1(\Omega \cap B(O, r))$ follows from (6.33) and (6.28), by observing that

$$\int_0^1 \frac{\rho^2 \, d\rho}{(\rho^2 - \rho^{\alpha_{\mathfrak{G}}})^\beta} = \frac{1}{\alpha_{\mathfrak{G}} - 2} \mathfrak{B}\Big(\frac{3 - 2\beta}{\alpha_{\mathfrak{G}} - 2}; 1 - \beta\Big) < +\infty. \tag{6.34}$$

Case (ii): Assume that the vertex O is such that $\alpha_{\mathfrak{G}} = 2$ (our subsequent analysis also applies to "fictitious" vertices, i.e., points of edges with $2\theta_j = \pi/2$). In this case, it follows from (3.8) that the first eigenvalue of $-\Delta_{S^2}$ on $\mathfrak{G}$ is $\Lambda_{\mathfrak{G}} = 6$, and we recall the eigenfunction $\phi_{\mathfrak{G}}$ from (3.3) (with $n = 3$). Also, fix some small $r > 0$ and set

$$v_r(x) := |x|^2 \Big(\log \frac{r}{|x|}\Big) \phi_{\mathfrak{G}}\Big(\frac{x}{|x|}\Big) \quad \text{on} \quad S_{\mathfrak{G}, r}. \tag{6.35}$$

Hence, by (6.29),

$$-\Delta v_r(x) = -\Big(\log \frac{r}{|x|}\Big)(\Delta_{S^2} \phi_{\mathfrak{G}} + 6\phi_{\mathfrak{G}}) + 5\phi_{\mathfrak{G}} = 5\phi_{\mathfrak{G}}. \tag{6.36}$$

Since $v_r\big|_{\partial S_{\mathfrak{G}, r}} = 0$ and $u\big|_{\partial S_{\mathfrak{G}, r}} \geq 0$, and since

$$-\Delta v_r(x) \leq 5 \max_{\mathfrak{G}} \phi_{\mathfrak{G}} = -\Delta(5 \max_{\mathfrak{G}} \phi_{\mathfrak{G}} \cdot u(x)), \tag{6.37}$$

it follows that

$$u(x) \geq (5 \max_{\mathfrak{G}} \phi_{\mathfrak{G}})^{-1} v_r(x) \quad \text{on} \quad S_{\mathfrak{G}, r}. \tag{6.38}$$

As a consequence of this, (6.35) and Corollary 6.23, we therefore obtain

$$u(x) \geq C|x|^2 (\sin \varphi_j)^{\pi/(2\theta_j)} \cos(\pi \psi_j/(2\theta_j)), \qquad \text{where} \quad x/|x| = (\psi_j, \varphi_j), \tag{6.39}$$

for x near O. In turn, the nature of the expression in the right-hand side of (6.39) guarantees the integrability of $u^{-\beta}$ for each $\beta \in (0, 1)$ over small conical neighbourhoods of every edge of a dihedral angle of half-opening $\theta_j \geq \pi/4$, since generally speaking

$$\int_0^1 \frac{\rho^2 \, d\rho}{(\rho^2)^\beta} \int_0^\pi \frac{\sin \varphi_j \, d\varphi_j}{(\sin \varphi_j)^{(\pi/(2\theta_j))\beta}} \int_{|\psi_j| < \theta_j} \frac{d\psi_j}{[\cos(\pi \psi_j/(2\theta_j)]^\beta} < +\infty \tag{6.40}$$

whenever $\beta < \min\{1, \frac{4\theta_j}{\pi}\}$.

The treatment of the case when $\theta_j < \pi/4$ requires a different approach which we now describe. Let us consider a conic neighbourhood, centred at O, of the j-th edge of a dihedral angle of opening $2\theta_j$ with $\theta_j < \pi/4$, i.e.,

$$U_j := \{x = (\rho, \psi_j, \varphi_j) : \rho > 0, \ |\psi_j| \leq \theta_j, \ 0 < \varphi_j < \epsilon\}, \tag{6.41}$$

where $\epsilon > 0$ is a small number. We choose a point P on this edge, set $r := |P|/3$ and apply (6.33) with the role of v played by the function

$$v^{(P)}(x) := \big(|x - P|^2 - |x - P|^{\pi/(2\theta_j)}\big) \Phi^{(j)}\big(\tfrac{x-P}{|x-P|}\big), \tag{6.42}$$

where $\Phi^{(j)}$ is our old $\Phi_{\mathfrak{G}}$ (cf. (6.27)) constructed for the point P. On account of the first asymptotic formula in (6.27), this gives (much as for (6.33)) that on $\{x \in U_j \cap \Omega : |x| = r\}$ we have

$$\begin{aligned} u(x) &\geq c_o^{-1} r^2 \left(\Big(\frac{|x - P|}{r}\Big)^2 - \Big(\frac{|x - P|}{r}\Big)^{\pi/(2\theta_j)} \right) \Phi^{(j)}\big(\tfrac{x-P}{|x-P|}\big) \\[6pt] &\geq C|x - P|^2 \varphi_j^2 \Big(\frac{\cos(2\psi_j)}{\cos(2\theta_j)} - 1\Big) \\[6pt] &\geq Cr^2 \varphi_j^2 \Big(\frac{\cos(2\psi_j)}{\cos(2\theta_j)} - 1\Big). \end{aligned} \tag{6.43}$$

To continue, let $\Gamma_\mathfrak{G}$ solve the boundary value problem

$$-\Delta_{S^2}\Gamma_\mathfrak{G} - 6\Gamma_\mathfrak{G} = 1 - \left(\int_\mathfrak{G}\phi_\mathfrak{G}(\omega)\,d\omega\right)^{-1}\phi_\mathfrak{G} \quad \text{in } \mathfrak{G},$$

$$\Gamma_\mathfrak{G}\Big|_{\partial_{S^2}\mathfrak{G}} = 0, \tag{6.44}$$

Much as in [21], it follows that there exists $\delta > 0$ with the property that

$$\Gamma_\mathfrak{G}(\omega) = 4^{-1}\varphi_j^2\big(1 + \mathcal{O}(\varphi_j^\delta)\big)\Big(\frac{\cos(2\psi_j)}{\cos(2\theta_j)} - 1\Big). \tag{6.45}$$

Also, clearly,

$$\Delta\big(|x|^2\Gamma_\mathfrak{G}(x/|x|)\big) = -(\Delta_{S^2}\Gamma_\mathfrak{G})(x/|x|) - 6\Gamma_\mathfrak{G}(x/|x|) \leq 1. \tag{6.46}$$

To continue, we consider a smooth, nonnegative cutoff function $\eta_j(\omega)$ with the property that $\eta_j = 1$ for $\varphi_j \in (0, \epsilon/2)$ and $\eta_j = 0$ for $\varphi_j \in (\epsilon, \pi)$. Then

$$-\Delta\big(\eta_j(x/|x|)|x|^2\Gamma_\mathfrak{G}(x/|x|)\big) \leq c_1 \quad \text{on} \quad S_{\mathfrak{G},r}. \tag{6.47}$$

Furthermore, by (6.45), on $\{x \in U_j \cap \Omega : |x| = r\}$ we have

$$\eta_j(\omega)|x|^2\Gamma_\mathfrak{G}(\omega) \leq cr^2\varphi_j^2\Big(\frac{\cos(2\psi_j)}{\cos(2\theta_j)} - 1\Big). \tag{6.48}$$

Note that the function in the left-hand side of (6.48) vanishes on the conical side of $\partial S_{\mathfrak{G},r}$. The same is, obviously, true for u. Consequently, on account of this, (6.43) and (6.47)-(6.48), we obtain

$$u(x) \geq c_2\eta_j(\omega)|x|^2\Gamma_\mathfrak{G}(\omega) \quad \text{on} \quad \{x \in U_j \cap \Omega : |x| = r\}. \tag{6.49}$$

Hence, by the Maximum Principle, we ultimately have

$$u(x) \geq c_2\eta_j\Big(\frac{x}{|x|}\Big)|x|^2\Gamma_\mathfrak{G}\Big(\frac{x}{|x|}\Big) \quad \text{for all} \quad x \in S_{\mathfrak{G},r}. \tag{6.50}$$

This, along with (6.45) now shows that $u^{-\beta}$ is integrable for all $\beta \in (0,1)$ over a small conical neighbourhood of the j-th edge, completing the treatment of the situation described in Case (ii).

Case (iii): Assume that the vertex O is such that $\alpha_\mathfrak{G} < 2$ (our analysis also applies to a point on the edge with $2\theta_j > \pi/2$). Hence, in this situation, $6 > \alpha_\mathfrak{G}(\alpha_\mathfrak{G} + 1)$ and, therefore, by a slight variant of results proved in [21] and [28], there exists a positive constant C_0 with the property that, granted that $|x|$ is small,

$$u(x) = C_0|x|^{\alpha_\mathfrak{G}}\phi_\mathfrak{G}\Big(\frac{x}{|x|}\Big) + \mathcal{O}(|x|^{\alpha_\mathfrak{G}+\delta}), \quad \text{for some } \delta > 0, \tag{6.51}$$

where $\phi_\mathfrak{G}$ is the eigenfunction of the Laplace-Beltrami operator Δ_{S^2} with Dirichlet boundary condition on $\mathfrak{G}$ corresponding to the eigenvalue $\alpha_\mathfrak{G}(\alpha_\mathfrak{G} + 1)$ (compare with (3.3) when $n = 3$). Let $\eta(\omega)$ denote a smooth cut-off function with support in a small conical neighbourhood of the edges, such that $\eta = 1$ in a conical neighbourhood of every edge. Furthermore, let $H(x)$ be a smooth cut-off function in $C_0^\infty(B(O,2))$ with small support near O, which is identically equal to 1 on $B(O,1)$.

As in Case (i), we need the solution $\Phi_\mathfrak{G}$ of problem (6.26). We wish to stress that, in the current context, it is not known whether $\Phi_\mathfrak{G}$ is of definite sign, but we shall not make use of this property. The asymptotic formulae for $\Phi_\mathfrak{G}(\omega)$ given in Case (i) for $\theta_j \leq \pi/4$ remain valid here as well. However, for $\theta_j > \pi/4$ we can only say that

$$\Phi_\mathfrak{G}(\omega) = \mathcal{O}\Big(\varphi_j^{\pi/2\theta_j}\cos\frac{\pi\psi_j}{\theta_j}\Big). \tag{6.52}$$

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For some small $r > 0$, let us now introduce the function

$$w(x) := |x|^2 \Phi_{\mathfrak{G}}\left(\frac{x}{|x|}\right) \eta\left(\frac{x}{|x|}\right) H\left(\frac{x}{r}\right). \tag{6.53}$$

Since $-\Delta(|x|^2 \Phi_{\mathfrak{G}}\left(\frac{x}{|x|}\right)) = 1$ on $S_{\mathfrak{G},\infty}$, we have $-\Delta w = 1$ in the intersection of $B(O,r)$ and small conical neighbourhoods of the edges meeting at O. Moreover, $|\Delta w| \le C$ on $S_{\mathfrak{G},r}$. Finally, introduce

$$W(x) := u(x) - w(x), \quad \forall\, x \in S_{\mathfrak{G},r} \tag{6.54}$$

and note that $W = 0$ on the conical side of $\partial S_{\mathfrak{G},r}$. Also,

$$-\Delta W = 1 + \Delta w \quad \text{on} \quad S_{\mathfrak{G},r}, \tag{6.55}$$

and the right-hand side is bounded, and vanishes on small conical neighbourhoods of the edges meeting at O. Consequently, by [21],

$$W(x) = C_1 |x|^{\alpha_{\mathfrak{G}}} \phi_{\mathfrak{G}}\left(\frac{x}{|x|}\right) + \mathcal{O}(|x|^{\alpha_{\mathfrak{G}}+\delta}). \tag{6.56}$$

Comparing this with (6.51) and using (6.52), (6.53) and (6.54), we derive from $\alpha_{\mathfrak{G}} < 2$ that actually $C_1 = C_0$. Using the asymptotics (6.22) of $\phi_{\mathfrak{G}}$ near angle vertices on $\partial_{S^2}\mathfrak{G}$ as well as the harmonicity of the remainder term in (6.56), we see that the remainder term in (6.56) can be replaced by $\mathcal{O}(|x|^{\alpha_{\mathfrak{G}}+\delta})\phi_{\mathfrak{G}}\left(\frac{x}{|x|}\right)$. (The argument is similar to the proof of Lemma 6.2 and we omit it.) Thus,

$$W(x) = C_0 |x|^{\alpha_{\mathfrak{G}}}\left(1 + \mathcal{O}(|x|^{\delta})\right)\phi_{\mathfrak{G}}\left(\frac{x}{|x|}\right), \tag{6.57}$$

and we have by (6.53) and (6.54)

$$\begin{aligned}
u(x) &= w(x) + W(x) \\
&= |x|^2 \Phi_{\mathfrak{G}}\left(\frac{x}{|x|}\right) \eta\left(\frac{x}{|x|}\right) H\left(\frac{x}{r}\right) + C_0 |x|^{\alpha_{\mathfrak{G}}}\left(1 + \mathcal{O}(|x|^{\delta})\right)\phi_{\mathfrak{G}}\left(\frac{x}{|x|}\right).
\end{aligned} \tag{6.58}$$

This means that for $|x| < r$ in a small angular neighbourhood of the edges meeting at O and such that $\theta_j \le \pi/4$ the following holds

$$u(x) \ge c_1 |x|^2 \Phi_{\mathfrak{G}}\left(\frac{x}{|x|}\right) \ge c_2(|x|\varphi_j)^2. \tag{6.59}$$

On the other hand, in a small angular neighbourhood of the edges with $\theta_j > \pi/4$, for $|x| < r$ we have that

$$\begin{aligned}
u(x) &\ge C_0 |x|^{\alpha_{\mathfrak{G}}}\left(1 + \mathcal{O}(|x|^{\delta})\right)\phi_{\mathfrak{G}}\left(\frac{x}{|x|}\right) - C|x|^2 \varphi_j^{\pi/(2\theta_j)} \cos\left(\frac{\pi\psi_j}{2\theta_j}\right) \\
&\ge C_1 |x|^{\alpha_{\mathfrak{G}}} \varphi_j^{\pi/(2\theta_j)} \cos\left(\frac{\pi\psi_j}{2\theta_j}\right),
\end{aligned} \tag{6.60}$$

where C_1 is a finite constant. Finally, when x is at a fixed, positive angular distance to the edges meeting at O, we may conclude in the same way as in (6.43) that

$$u(x) \ge C_2 |x|^{\alpha_{\mathfrak{G}}-1} \text{dist}\,(x, \partial\Omega). \tag{6.61}$$

Collectively, (6.59), (6.60) and (6.61) prove that $u^{-\beta}$ is integrable near O, for each number β in the interval $(0,1)$. $\qquad\square$

6.2 Piecewise C^2 domains with outward cuspidal vertices in $\mathbb{R}^n$

Here we elaborate on the case of domains with exterior cusps. This class of domains consists of bounded open sets $\Omega \subseteq \mathbb{R}^n$ with a piecewise C^2 boundary, exhibiting finitely many exterior cusps. By definition, $x_0 \in \partial\Omega$ is called an exterior cusp if, after a rigid transformation of the space which maps x_0 into the origin of $\mathbb{R}^n$, there exist two small numbers $\varepsilon, \eta > 0$ along with a function $\mathcal{F} \in C^2([0, \eta])$ with $\mathcal{F} > 0$ on $(0, \eta]$ and $(\mathcal{F}^2)''(0) = 0$, and such that $\{x \in \Omega : x_n \leq \eta\}$ coincides with the cuspidal set

$$\{x = (x', x_n) : 0 < x_n \leq \eta, \ |x'| < \varepsilon\mathcal{F}(x_n)\}. \tag{6.62}$$

Theorem 6.4. *Assume that u is the solution of the Saint Venant problem* (2.1) *in the case when $\Omega \subseteq \mathbb{R}^n$, $n \geq 2$, is a domain with exterior cusps. If $n \geq 3$ then for every $\beta \in (0, 1)$ there holds*

$$\int_\Omega u(x)^{-\beta}\,dx \leq C(\Omega, \beta) < +\infty \tag{6.63}$$

The same is true in the two dimensional setting provided $0 < \beta \leq 1/2$. Finally, in the case when $n = 2$ and $\beta \in (1/2, 1)$, then (6.63) *holds if and only the finiteness condition*

$$\int_0^\eta \mathcal{F}(\tau)^{1-2\beta}\,d\tau < +\infty \tag{6.64}$$

holds for every boundary cusp.

Proof. Assume that $0 \in \partial\Omega$ is a cusp, and $\mathcal{F}$ is as in the preamble to this subsection. Without loss of generality assume that $\eta = 1$ (hence $\mathcal{F}$ is defined on $[0, 1]$). The function $v(x) := \varepsilon^2\mathcal{F}^2(x_n) - |x'|^2$ satisfies

$$-\Delta v(x) = 2(n-1) - \varepsilon^2(\mathcal{F}^2)''(x_n) \tag{6.65}$$

and its trace on $\partial\Omega$ is nonnegative and vanishes when $x_n < 1$.

As usual, u stands for the unique solution of

$$-\Delta u = 1 \ \text{ in } \Omega, \quad u = 0 \ \text{ on } \partial\Omega. \tag{6.66}$$

Since $-\Delta v > 1$ on $\{x \in \Omega : x_n < 1\}$ (assuming ε small), we have

$$u \leq v \ \text{ on } \ \{x \in \Omega : x_n < 1\}. \tag{6.67}$$

We are next going to obtain an opposite estimate. First, by the smallness of ε, we have

$$-\Delta v \leq 2n = -2n\Delta u. \tag{6.68}$$

Next, by Hopf's Lemma,

$$\frac{\partial u}{\partial \nu} \geq C > 0 \ \text{ on } \ \{x \in \partial\Omega : x_n = 1\}, \tag{6.69}$$

where ν is the interior normal direction. Therefore,

$$u(x) \geq c_1 \operatorname{dist}(x, \partial\Omega) \ \text{ on } \ \{x \in \Omega : x_n < 1\}. \tag{6.70}$$

The estimate

$$v(x) \leq c_2 \operatorname{dist}(x, \partial\Omega) \ \text{ on } \ \{x \in \Omega : x_n < 1\} \tag{6.71}$$

follows directly from the definition of v. Hence,

$$v(x) \leq \frac{c_2}{c_1}u(x) \ \text{ on } \ \{x \in \Omega : x_n \leq 1\} \tag{6.72}$$

which, together with (6.68) and the Maximum Principle, gives

$$v(x) \leq c\, u(x) \quad \text{on} \quad \{x \in \Omega : x_n \leq 1\}. \tag{6.73}$$

If the origin is the only singularity of $\partial\Omega$, the condition $u^{-\beta} \in L^1(\Omega)$ is equivalent to

$$\int_{\{x \in \Omega : x_n < 1\}} \frac{dx}{v(x)^\beta} < +\infty, \tag{6.74}$$

which is the same as

$$\int_0^1 dx_n \int_{|x| < \varepsilon\mathcal{F}(x_n)} \frac{dx'}{v(x)^\beta} < +\infty. \tag{6.75}$$

This can be written in the form

$$
\begin{aligned}
+\infty \quad > \quad & \int_0^1 dx_n \int_0^{\varepsilon\mathcal{F}(x_n)} \frac{\rho^{n-2}d\rho}{(\varepsilon^2\mathcal{F}^2(x_n) - \rho^2)^\beta} \\
= \quad & \int_0^1 (\varepsilon\mathcal{F}(x_n))^{n-1-2\beta}\, dx_n \int_0^1 \frac{\tau^{n-2}\, d\tau}{(1-\tau^2)^\beta},
\end{aligned}
\tag{6.76}
$$

which is satisfied if and only if $\beta < 1$ for $n \geq 3$, and if and only if $\beta < 1$ and

$$\int_0^1 \mathcal{F}(\tau)^{1-2\beta}\, d\tau < +\infty \tag{6.77}$$

for $n = 2$ (note that (6.77) is always satisfied when $\beta \leq 1/2$; compare also with the example (2.60) in which case we take $\mathcal{F}(\tau) := \tau^{1/(2\beta-1)}$).

Once the contribution from near each boundary cusp has been estimated, the end-game of the argument is similar to that of the proof of Theorem 5.2. $\qquad\square$

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