

Dedicated to V. G. Maz'ya on the occasion of his 80th birthday

SPECTRAL PROPERTIES OF NONASSOCIATIVE ALGEBRAS AND BREAKING REGULARITY FOR NONLINEAR ELLIPTIC TYPE PDES

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In this paper, we address the following question: Why certain nonassociative algebra structures emerge in the regularity theory of elliptic type PDEs and also in constructing nonclassical and singular solutions? Our aim in the paper is twofold. First, to give a survey of diverse examples on nonregular solutions to elliptic PDEs with emphasis on recent results on nonclassical solutions to fully nonlinear equations. Second, to define an appropriate algebraic formalism, which makes the analytic part of the construction of nonclassical solutions more transparent.

§1. Introduction

The first examples of nonassociative algebras (including octonions and Lie algebras) appeared in the mid-19th century. Since then the theory has evolved into an independent branch of algebra, exhibiting many points of contact with other fields of mathematics, physics, mechanics, biology and other sciences. Algebras whose associativity is replaced by identities (as the famous Jacobi identity in Lie algebras or the Jordan identity $[\text{ad}(x), \text{ad}(x^2)] = 0$ in Jordan algebras) were a central topic in mathematics in the 20th century.

Very recently, certain commutative nonassociative algebra structures emerged in a very different context: regularity of viscosity solutions to fully nonlinear elliptic PDEs. More precisely, it turns out that subalgebras of simple rank three formally real Jordan algebras can be employed for constructing truly viscosity solutions to uniformly elliptic Hessian equations. To explain this relationships, we first review some relevant concepts and results.

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1.1. Linear and quasilinear elliptic type PDEs. Starting with the pioneering works of S. N. Bernstein, the maximum principle and *a priori* estimates have provided decisive instruments in proving the existence of solutions of general linear and quasilinear elliptic type PDEs $Lw = 0$. Important progress in this direction was made in the late 1950s by E. De Giorgi, J. Moser, and J. Nash by establishing fundamental *a priori* regularity results. A fundamental result of the regularity theory asserts that every solution w from a ‘natural’ class $W^{1,2}$ of $Lw = 0$ must be Hölder continuous, i.e., $w \in C^{0,\alpha}$ for a certain $\alpha > 0$.

On the other hand, the important counterexamples independently constructed in 1968 for scalar equations by V. G. Maz’ya [31] and for linear elliptic systems by De Giorgi [14] and E. Giusti and M. Miranda [17] showed that there is no natural generalization of De Giorgi’s theorem on the Hölder continuity for generalized solutions to strongly elliptic differential equations of order greater than two.

In particular, in [31] Maz’ya constructed several examples of quasilinear strongly elliptic equations of order $2l \geq 2$ with C^∞ -coefficients whose generalized solutions are not C^∞ regular. The strong ellipticity of an operator $Lu = \sum D^\alpha(a_{\alpha\beta}D^\beta u)$ means that there exists $c > 0$ such that the inequality $\int_{\mathbb{R}^n} \sum a_{\alpha\beta} D^\alpha \phi D^\beta \phi dx \geq c \int_{\mathbb{R}^n} \sum (D^\alpha \phi)^2 dx$ is fulfilled for any $\phi \in C^\infty(\mathbb{R}^n)$ with a compact support, where summation is over all $|\alpha| = |\beta| = l$. Then Maz’ya noted that the equation

$$\nu \Delta^2 w + \varkappa \Delta \left(\frac{x_i x_j}{|x|^2} w_{x_i x_j} \right) + \varkappa \Delta \left(\frac{x_i x_j}{|x|^2} \Delta w \right)_{x_i x_j} + \mu \left(\frac{x_i x_j x_k x_l}{|x|^4} w_{x_i x_j} \right)_{x_k x_l} = 0$$

is strongly elliptic for $\varkappa^2 < \mu\nu$ and the radial symmetric function $w = |x|^a$ with

$$a = 2 - \frac{n}{2} + \sqrt{\frac{n^2}{4} - \frac{(n-1)(\varkappa n + \mu)}{\nu + 2\varkappa + \mu}}$$

is a weak solution of the above equation in $W^{2,2}(B_1)$, where B_1 is the open unit ball in \mathbb{R}^n . For $\varkappa = n(n-2)$, $\mu = n^2$ and $\nu = (n-2)^2 + \varepsilon$, with $\varepsilon < 0$, the above strong ellipticity condition is satisfied and the corresponding solution $w = |x|^a$ is unbounded in B for all dimensions $n \geq 5$ and sufficiently small ε .

A very related problem is the regularity of p -harmonic functions. It is well known that for $p > 1$, $p \neq 2$, a weak (in the distributional sense) p -harmonic function is normally in the Hölder class $C^{1,\alpha}$ [59], but need not to be Hölder continuous or even continuous in a closed domain with nonregular boundary, as it follows from examples constructed by Krol’ and Maz’ya in [26]. On the other hand, if $\text{ess sup } |Du(x)| > 0$ holds locally in a domain, then $u(x)$ is in fact a real analytic function in E [30]. It is interesting whether the converse nonvanishing property holds true. This question naturally leads to homogeneous p -harmonic

functions (also called quasiradial solutions in [4]); see the paper of M. Akman, J. Lewis, A. Vogel [1] in the present volume for the modern status of the problem and further discussion. We only mention that in certain dimensions $n \geq 4$ one can explicitly construct rational homogenous p -harmonic functions for some distinguished values p , $2 < p < n$, based on isoparametric forms and generalizing the angle solutions constructed by Krol' and Maz'ya, see [55].

Another prominent example of breaking regularity is the existence of globally minimal cones and entire solutions of the minimal surface equation in dimensions $n \geq 8$; see [6]. The higher-dimensional examples constructed by Bombieri, De Giorgio, and Giusti [6] make an essential use of the quadratic isoparametric form

$$u_2(x) := x_1^2 + x_2^2 + x_3^2 + x_4^2 - x_5^2 - x_6^2 - x_7^2 - x_8^2, \quad (1)$$

see also [50] and [49] for the general isoparametric case. The function $u_2(x)$ coincides with the norm in the split octonion¹ algebra [5]. It is not clear whether this is merely a coincidence or there is some natural explanation in the context of singular and entire solutions.

On the other hand, Lawson and Osserman [29] constructed nonparametric minimal cones of high codimensions providing examples of Lipschitz but non- C^1 solutions to the minimal surface equations, thereby making sharp contrast to the regularity theorem for minimal graphs of codimension one. The Lawson–Osserman examples based on Hopf foliations corresponding to the classical division algebras \mathbb{A}_d : $\mathbb{A}_2 = \mathbb{C}$, $\mathbb{A}_4 = \mathbb{H}$, and $\mathbb{A}_8 = \mathbb{O}$. More explicitly, for $d = 2, 4, 8$ let

$$w(x) = \frac{1}{2} \sqrt{\frac{2d+1}{d-1}} \frac{\eta(x)}{|x|} : \mathbb{R}^{2d} \rightarrow \mathbb{R}^{d+1}, \quad (2)$$

where

$$\eta(x) = (|z_1|^2 - |z_2|^2, 2z_1\bar{z}_2), \quad x = (z_1, z_2) \in \mathbb{A}_d^2 \cong \mathbb{R}^{2d}. \quad (3)$$

Then the map $w(x) : \mathbb{R}^{2d} \rightarrow \mathbb{R}^{d+1}$ provides a nontrivial Lipschitz solution to the Bernstein problem in codimension $d + 1$.

We finally mention some further classes of semilinear and quasilinear elliptic PDEs with similar phenomena, for example, the recent entire solutions of the Ginzburg–Landau system constructed by A. Farina in [16] by using isoparametric forms; see also Section 3 in [33] for a recent survey of further counterexamples as well as regularity results.

¹The author is grateful to Seidon Alsaody for pointing out this.

1.2. Fully nonlinear elliptic type PDEs. In the general, fully nonlinear case, the regularity and existence issues become more involved and require two principal ingredients: (i) the Harnack inequality for solutions of the 2nd order nondivergence elliptic equations with measurable coefficients established in 1979 by N. Krylov and M. V. Safonov, see [27, 28], and (ii) a suitable extension of the concept of a generalized solution, the so-called *viscosity solution*, an important tool developed by Crandall–Lions, Evans, Jensen, and Ishii. We refer to [11] for exact definitions and self-contained exposition of the basic theory of viscosity solutions and briefly note that this concept generalizes weak solutions for divergence type elliptic equations by utilizing the maximum principle. Viscosity solutions play a central role in many problems from global geometry and analysis, and fit naturally into the contexts of optimal control, differential and stochastic differential games, as well as mathematical finance.

An important model example with a far-ranging set of applications in pure and applied mathematics is the Dirichlet problem for the uniformly elliptic equation of the form

$$\begin{aligned} F(D^2w) &= 0 \quad \text{in } B_1 \\ w &= \phi \quad \text{on } \partial B_1 \end{aligned} \tag{4}$$

with a uniformly elliptic operator F . The latter means that there exist $0 < \lambda \leq \Lambda$ such that

$$\lambda \|N\| \leq F(M + N) - F(M) \leq \Lambda \|N\|$$

whenever M, N are symmetric matrices with N nonnegative semidefinite.

For many geometrical applications, an important particular case of (4) is the *Hessian equations*, i.e., when $F(X)$ depends only on the eigenvalues of the (symmetric) matrix X . The Hessian equations include the classical examples of Laplace’s equation, curvature Weingarten equation, Monge-Ampère equation and special Lagrange equation, and, more recently, the calibrated geometries and Dirichlet duality theory [19, 20]. According to the general theory [11], a fully nonlinear elliptic equation (4) always has a *unique* continuous viscosity solution for any continuous data ϕ . The question whether a viscosity solution is classical, i.e., twice differentiable, has turned out to be very challenging. For certain cases the regularity is known:

- (i) If $n = 2$, then w is a classical ($C^{2,\alpha}$) solution (Nirenberg [44])
- (ii) If $n \geq 2$ and $\Lambda \leq 1 + \epsilon(n)$, then w is $C^{2,\alpha}$ (Cordes [10])
- (iii) $w \in C^{1,\alpha}(B_{1/2})$, $\alpha = \alpha(\Lambda, n)$ (Trudinger [58], Caffarelli [7])
- (iv) F convex (concave) $\Rightarrow w \in C^{2,\alpha}(B_{1/2})$ (Krylov, Evans [15])
- (v) w is $C^{2,\alpha}(B_1 \setminus \Sigma)$, where the Hausdorff dimension $\dim_H \Sigma$ is strictly smaller than $n - \epsilon$, $\epsilon = \epsilon(\Lambda, n)$ (Armstrong–Silvestere–Smart, [3]).

Further regularity results including the modulus of continuity of viscosity solutions can be found in [22, 51]; see also [13] for the VMO-regularity and gradient estimates for viscosity solutions to nonhomogeneous fully nonlinear elliptic equations $F(D^2u) = f$ with $f \in L^p$.

1.3. Truly viscosity solutions. From the above, it follows that in the general case and in dimensions $n \geq 3$ only $C^{1,\alpha}$ -regular viscosity can be expected. Until very recently the very existence of nonclassical solutions for $n \geq 3$ was not known. A major breakthrough in this direction was achieved in a series of papers of Nadirashvili and Vlăduț [37, 38, 41] by establishing first the existence of nonclassical and then also truly singular (Hölder continuous) viscosity solution in any dimension $n \geq 12$. The method of construction makes an essential use of an accumulation property of the spectrum of the Hessian of a certain cubic form u_{12} in \mathbb{R}^{12} . This form is also well known as a *triatlity form* and comes from the Hamilton quaternions \mathbb{H} , [5].

Theorem 1.1 (see [37, 41]). *Let*

$$u_{12}(x) = \operatorname{Re}(z_1 z_2 z_3), \text{ where } x = (z_1, z_2, z_3) \in \mathbb{H}^3 \cong \mathbb{R}^{12} \quad (5)$$

and

$$w_{12,\alpha}(x) = \frac{u_{12}(x)}{|x|^\alpha}. \quad (6)$$

If $1 \leq \alpha < 2$, then $w_{12,\alpha}$ is a viscosity solution in \mathbb{R}^{12} of a uniformly elliptic Hessian equation (4) with a smooth F .

In view of the appearance of the quaternions in (6), it is natural to expect that the corresponding cubic form u_{24} based on the octonions also produces singular viscosity solutions. The paper [40] established that this is actually the case. Furthermore, in [35] a nonclassical viscosity solution w_5 in \mathbb{R}^5 based on the Cartan isoparametric cubic form u_5 was constructed and, shortly after, in [43] Nadirashvili and Vlăduț established that the corresponding $w_{5,\alpha}$ is a singular viscosity solutions in \mathbb{R}^5 . More precisely, one has the following fact.

Theorem 1.2 (see [35]). *There a cubic form $u_5(x)$ such that $w_{5,\alpha}(x) = \frac{u_5(x)}{|x|^\alpha}$ is a viscosity solution in $B_1 \subset \mathbb{R}^5$. If $1 \leq \alpha < 2$, then $w_{5,\alpha}$ is a viscosity solution in \mathbb{R}^{12} of a uniformly elliptic Hessian equation (4) with a smooth F . Furthermore, $w_5 := w_{5,1}$ satisfies (4) with*

$$F(D^2w) = (\Delta w)^5 + 2^8 3^2 (\Delta w)^3 + 2^{12} 3^5 \Delta w + 2^{15} \det D^2(w).$$

This establishes the lowest possible dimension known so far where singular viscosity solutions to uniformly elliptic equations may exist.

The algebraic part of the above construction relies heavily on certain accumulation properties of the spectrum of the cubic forms u_k involved. It was noticed in [39] that to obtain a nonclassical viscosity solution, the corresponding cubic form u_n must be rather exceptional. Furthermore, numerical simulations show that a random ('generic') cubic form does *not* produce a nonclassical solution to (4). This makes it reasonable to ask why u_{12} , u_{24} , and u_5 are so exceptional and how to characterize the cubic forms producing singular/nonclassical viscosity solutions.

We have already mentioned that u_{12} and u_{24} come from the triality concept and are intimately related to exceptional Lie groups and the classical division algebras in dimensions 4 and 8 respectively. From [39], it follows that the corresponding triality cubic form u_6 over the field of complex numbers \mathbb{C} also produces a viscosity solution but only for a degenerate elliptic equation.

The cubic form u_5 is also rather special and has many important relationships in analysis and geometry. In the present context, at least the following properties are essential:

- (a) u_5 is the generic norm on (the trace free subspace of) rank 3 formally real Jordan algebra $\mathcal{H}_3(\mathbb{R})$ of symmetric 3×3 matrices, and
- (b) u_5 is the simplest Cartan isoparametric cubic form.

There exists a natural correspondence between the concepts in (a) and (b), see [52]. Explicitly, the cubic form u_5 is given by the determinant representation

$$u_5(x) = \begin{vmatrix} \frac{1}{\sqrt{3}}x_1 + x_2 & x_3 & x_4 \\ x_2 & \frac{-2}{\sqrt{3}}x_1 & x_5 \\ x_4 & x_5 & \frac{1}{\sqrt{3}}x_1 - x_2 \end{vmatrix} \quad (7)$$

Taking the above observations into account, it is highly desirable to find a conceptual explanation of the existence of nonclassical solutions and relate them to appropriate algebraic structures. In this paper, we discuss some possible approaches to these questions.

The appearance of Jordan algebras in (7) and (a) makes it natural to ask whether these algebras are also relevant for the examples in Subsection 1.1. As we have already mentioned, the isoparametric form u_2 in (1) reminisces the norm of split octonions. Moreover, the Hopf map (3) coincides with the multiplication in Clifford type Jordan algebras (the so-called spin-factors), see [5].

Remark 1.3. In all the above examples of nonregular solutions, including the p -harmonic homogenous examples, the regularity breaks at an isolated point where (the blow-down of) a solution has a Hölder type singularity. Remarkably, the asymptotic cone at the singular point is always a *minimal cone*. It is also interesting to know whether the appearance of the minimality in this context

is mere coincidence or there is a general explanation of this phenomenon. This question also naturally arises in connection with phase transitions and the De Giorgi conjecture, see [46].

The rest of the paper is structured as follows. In §2 we discuss the main ideas underlying the construction of homogeneous nonclassical viscosity solutions of (4). There we also define an algebraic formalism, which connects cubic forms with certain commutative nonassociative algebras. In this setting, the gradient of a cubic form determines the multiplicative algebra structure; in particular, algebra idempotents correspond to stationary points of the cubic form. A cubic form is called special if it generates a nonclassical solution. We reformulate this definition in terms of the Peirce spectrum. We discuss the Peirce decomposition and generic cubic forms in Subsections 2.4 and 2.5. In Subsection 2.6 we apply the algebraic approach to special forms. In the last §3 we revisit Cartan's isoparametric cubics and determine their Peirce spectrum.

§2. Ingredients and definitions

2.1. Preliminaries. Recall that according to the regularity result [3], if w is a nonclassical solution to (4) in the unit ball B_1 , then its singular set must be compactly contained in B_1 . It is natural to consider the simplest single point singularity and look for candidates for nonclassical viscosity solutions in the unit ball $B_1 \subset \mathbb{R}^n$ in the class of homogeneous solutions, i.e., the functions w satisfying

$$w(tx) = t^k w(x), \quad x \in \mathbb{R}^n, \quad t > 0.$$

The ambient dimension must be $n \geq 5$. Indeed, for $n = 2$ from the Nirenberg result (i) mentioned above it follows that *any* (not necessarily homogeneous) viscosity solution is classical, and also if $3 \leq n \leq 4$, then [42] shows that a nonclassical viscosity solution, if exists, is never homogeneous (the case of $n = 3$ follows from a theorem of A. D. Alexandroff in [2]).

Furthermore, it is natural to specify $w(x)$ and assume that it has the following form:

$$w(x) = \frac{u(x)}{|x|^\alpha}, \quad (8)$$

$u(x)$ being a cubic form² in the variable $x \in \mathbb{R}^n$ and $\alpha \geq 1$. Under the made assumptions, it follows from [35] that the class of nonclassical viscosity solutions with $\alpha = 1$ is nonempty for any dimension $n \geq 5$.

²In fact, one can also construct some counterexamples w as in (8) by using isoparametric homogeneous forms u of degree 3, 4, and 6, see Chapter 5 in [36]. Interestingly, there are also *quadratic* isoparametric forms, but they are nonappropriate for constructing nonclassical solutions. In this paper we confine ourselves to the degree 3 case.

In all known cases, the existence of a nonclassical viscosity solution (8) for a certain $u(x)$ with $\alpha = 1$ guarantees that (8) with $1 < \alpha < 2$ will also be a *singular* viscosity solution to a certain F in (4), see Chapter 4 in [36] (it is believed that this observation holds true in general, though no proof is known). In view of this, we shall only consider nonclassical solutions.

2.2. Nonclassical viscosity solutions. We briefly recall the main ideas of the construction along the line of [35, 39]. For more information and the proofs, the reader is referred to Chapters 4 and 5 in [36]. Let w be an arbitrary homogeneous function of order 2 defined on \mathbb{R}^n and smooth in $\mathbb{R}^n \setminus \{0\}$. Then the Hessian $D^2w(x)$ is homogeneous of order 0 at any point $x \in \mathbb{R}^n \setminus \{0\}$ and it induces the Hessian map

$$D^2w : \mathbb{R}^n \setminus \{0\} \rightarrow \text{Sym}_n,$$

where Sym_n denotes the space of symmetric $(n \times n)$ -matrices with real coefficients. By the homogeneity assumption, we have

$$D^2w(tx) = D^2w(x) \quad \text{for any } t > 0,$$

hence the Hessian D^2w is completely determined by its values on the unit sphere ∂B_1 . We denote by

$$H_w = D^2w|_{\partial B_1}$$

the corresponding restriction map. By $O(n)$ we denote the group of isometries in \mathbb{R}^n .

As in [39], we say that a symmetric matrix A is M -hyperbolic for some real $M \geq 1$ if either all λ_i are equal to 0 or

$$\frac{1}{M} \leq \frac{-\lambda_1}{\lambda_n} \leq M, \tag{9}$$

where $\lambda_1 \leq \dots \leq \lambda_n$ denote the eigenvalues of A .

A subset $\mathcal{A} \subset \text{Sym}_n$ is called hyperbolic if there exists a constant M such that for any two matrices $A, B \in \mathcal{A}$, the difference $A - B$ is M -hyperbolic.

The importance of the last concepts is clear from the following result.

Lemma 2.1 (Main Lemma in [37]). *Let a function w be defined by (8) with $\alpha = 1$. Suppose that the restriction map $H_w : \partial B_1 \rightarrow \text{Sym}_n$ is a smooth embedding. Then*

- (i) *if $H_w(\partial B_1)$ is hyperbolic, then w is a viscosity solution of a uniformly elliptic equation (4) in \mathbb{R}^n ;*
- (ii) *if additionally the (larger) subset $\{UAU^{-1} : A \in H_w(\partial B_1), U \in O(n)\}$ is hyperbolic, then w is a solution of a Hessian uniformly elliptic equation (4) in \mathbb{R}^n .*

Thus, the above hyperbolicity condition is crucial for constructing viscosity solutions. Note also that all examples discussed in Subsection 1.3 satisfy in fact the stronger condition (ii). This suggests the following definition.

Definition 2.2. A cubic form $u(x)$ is called *special* if $w = u(x)/|x|$ satisfies hypotheses (ii) in Lemma 2.1.

It is also interesting to characterize the cubic forms u that satisfy the weaker condition (i). For continuity reasons, if u satisfies (i), then all cubic forms ‘sufficiently close’ to u will also satisfy (i).

Remark 2.3. An example of such a form is the determinant $u_9(x) = \det X$, where $X \in \mathbb{R}^{3 \times 3} \cong \mathbb{R}^9$ (X is an arbitrary nonsymmetric matrix over \mathbb{R}). Then it is easy to see that u_9 is the specialization of the triality form u_{12} on the imaginary subspace $(\text{Im } \mathbb{H})^3$, see (5). On the other hand, it is known, see Remark 4.1.5 in [36], that u_9 satisfies the ellipticity criterion (i) and thus produces a nonclassical solution $w = u/|x|$ of a uniformly elliptic fully nonlinear equation in \mathbb{R}^9 . However, a more sophisticated analysis reveals that this function fails to satisfy the sufficient conditions (ii) and thus may be not a solution of a Hessian uniformly elliptic fully nonlinear equation. Remarkably, u_9 is one of the exceptional Hsiang cubic forms emerging in classification of cubic minimal cones [36, Chapter 6], cf. also with the remarks in §4 below.

2.3. Algebras of cubic forms. The main obstruction for identifying which cubic forms u are special is an *a priori* hard problem to control the spectrum of the difference $H_w(x) - H_w(y)$ for all possible pairs $x, y \in \partial B_1$. This problem becomes more tractable and conceptually more transparent if one passes to a certain natural nonassociative algebra attached to the cubic form u . Below, we recall some relevant concepts and results following [53, 54, 56]; see also Section 6.3 in [36].

By an algebra we always mean a (finite dimensional) vector space V over the real numbers \mathbb{R} with a multiplication on V , i.e., a bilinear map (denoted by juxtaposition) $(x, y) \mapsto xy \in V$. By an abuse of notation, we denote by V both the vector space and the corresponding algebra.

In this paper, by an algebra we mean a commutative algebra. A symmetric bilinear form $b : V \times V \rightarrow \mathbb{R}$ on an algebra V is called *associating* if

$$b(xy, z) = b(x, yz), \quad x, y, z \in V. \quad (10)$$

An algebra carrying an associating bilinear form is called *metrised*. The standard example is the generic trace form on a Jordan algebra, see [32].

The most striking consequence of the existence of an associating bilinear form is that the operator of left (=right) multiplication

$$L_x : y \rightarrow xy = yx$$

is selfadjoint, i.e.,

$$\langle L_x y, z \rangle = \langle y, L_x y z \rangle, \quad x, y, z \in V. \quad (11)$$

Now we want to associate an algebra with an arbitrary cubic form. Recall that a cubic form is a function $u : V \rightarrow \mathbb{R}$ such that its full linearization

$$u(x, y, z) = u(x + y + z) - u(x + y) - u(x + z) - u(y + z) + u(x) + u(y) + u(z)$$

is a trilinear form. The trilinear form $u(x, y, z)$ is obviously symmetric and the cubic form is recovered by

$$6u(x) = u(x, x, x). \quad (12)$$

A positive definite symmetric bilinear form $\langle x, y \rangle$ on V is called an inner product. Given an inner product and a cubic form u , its *gradient* $\nabla u(x)$ at x is uniquely determined by the duality

$$\langle \nabla u(x), y \rangle = \frac{1}{2}u(x, x, y), \quad y \in V. \quad (13)$$

From the above definitions, it follows, that thus defined, ∇u coincides with the standard gradient of a function u .

Definition 2.4. Let u be a cubic form on an inner product vector space $(V, \langle \cdot, \cdot \rangle)$. Define the multiplication $(x, y) \rightarrow xy$ as a unique element satisfying

$$\langle xy, z \rangle = u(x, y, z) \text{ for all } z \in V \quad (14)$$

The multiplication is commutative but may be nonassociative. We call thus defined algebra on V the *algebra of the cubic form* u and denote it by $V(u)$.

It is easy to see that $V(u)$ is a zero algebra (i.e., $V(u)V(u) = 0$) if and only if the cubic form u is identically zero.

An important consequence of the above definition and the symmetry of the trilinear form $u(x, y, z)$ is that the inner product on the algebra $V(u)$ is associating, i.e., $V(u)$ is a *metrised commutative algebra*.

Furthermore, it also follows from the definition that, given an inner product vector space $(V, \langle \cdot, \cdot \rangle)$, there is a canonical bijection between the vector space $\mathcal{C}(V)$ of all cubic forms on V and commutative metrised algebra structures $\mathcal{A}(V)$ on V , where the correspondence $\mathcal{C}(V) \rightarrow \mathcal{A}(V)$ is given by (14), and the converse correspondence $\mathcal{A}(V) \rightarrow \mathcal{C}(V)$ is defined by

$$u(x) := \frac{1}{6}\langle x^2, x \rangle. \quad (15)$$

With the above concepts at hand, we are able to identify the standard calculus operations on V as appropriate algebraic concepts on $V(u)$. First note

that (13) combined with (14) yields immediately the fact that the gradient of $u(x)$ is essentially the square of the element x in $V(u)$:

$$\nabla u(x) = \frac{1}{2}xx = \frac{1}{2}x^2, \quad (16)$$

and, similarly, the multiplication in $V(u)$ is explicitly recovered by

$$(D^2u(x))y = xy, \quad (17)$$

where $D^2u(x)$ is the Hessian map of u at x . This implies that the Hessian map of u at x is none other else than the multiplication operator L_x :

$$D^2u(x) = L_x. \quad (18)$$

Then from (10) it follows that L_x is a selfadjoint operator with respect to the inner product $\langle \cdot, \cdot \rangle$.

In summary, starting with a cubic form on an inner product vector space, one can construct an algebra structure that translates the standard calculus into appropriate algebraic concepts. Then the Peirce decomposition relative to an idempotent is a significant tool in identifying the underlying nonassociative algebra structure, see [32, 47]. Therefore, it is important to know whether the set of idempotents in $V(u)$ is empty or not. The next proposition answers this question in the positive; its proof can be found in the author's papers [54, 56].

Proposition 2.5. *Let $V(u)$ be the algebra of a cubic form $u \not\equiv 0$ on an inner product vector space V . Denote by E the (nonempty) set of stationary points of the variational problem*

$$\langle x, x^2 \rangle \rightarrow \max \quad \text{subject to the constraint} \quad \langle x, x \rangle = 1. \quad (19)$$

Then for any $x \in E$, either $x^2 = 0$ or $c := \frac{x}{\langle x^2, x \rangle}$ is an idempotent in V . In particular, the set of idempotents of $V(u)$ is nonempty. Furthermore, if $E_0 \subset E$ denotes the subset of local maxima in (19), then the corresponding idempotent $c = x/\langle x^2, x \rangle$ satisfies the extremal property

$$\text{spectrum}(L_c|_{c^\perp}) \subset (-\infty, \frac{1}{2}], \quad (20)$$

where $c^\perp = \{x \in V : \langle x, c \rangle = 0\}$. In particular, the eigenvalue 1 of L_c is simple (i.e., c is a primitive idempotent).

2.4. Jordan algebras. Jordan algebras constitute an important class in the context of cubic forms and having numerous applications in mathematics and mathematical physics. Recall that a commutative algebra V is said to be *Jordan* if

$$[L_x, L_{x^2}] = 0. \quad (21)$$

It is well known that (21) implies that any Jordan algebra is power-associative, i.e., the powers x^n do not depend on associations, see [23].

In the most interesting case for applications when V is formally real, the bilinear form $b(x, y) := \text{tr } L_{xy}$ (the generic trace form) is positive definite and associating, i.e., satisfies (10). In their famous work [24], Jordan, von Neumann, and Wigner proved that the only simple formally real Jordan algebras are the Jordan algebras of $n \times n$ selfadjoint matrices over \mathbb{A}_d , $d = 1, 2, 3$, the exceptional 27-dimensional Albert algebra of 3×3 selfadjoint matrices over octonions, and the so-called spin-factors, or Clifford type Jordan algebras.

Recall that the (Peirce) *spectrum* of an idempotent c in a commutative algebra is called the spectrum of the corresponding multiplication operator L_c . The algebra spectrum is the union of all idempotent spectra.

An important structure property of any Jordan algebra V is that the spectrum of *any* idempotent is a subset of $\{0, \frac{1}{2}, 1\}$. In particular, V admits the so-called *Peirce decomposition*:

$$V = V_0(c) \oplus V_{\frac{1}{2}}(c) \oplus V_1(c), \quad (22)$$

where V_λ is the λ -subspace of L_c . Then the $V_c(\lambda)$ are subalgebras of V for $\lambda \in \{0, 1\}$.

The eigenvalue $\frac{1}{2}$ and the corresponding subspace is distinguished in many ways. For example, a Jordan algebra is simple if and only its $\frac{1}{2}$ -eigenspace is nontrivial for any idempotent. Moreover, $V_{\frac{1}{2}}$ satisfies the Jordan fusion laws

$$V_c(\lambda) V_c(\frac{1}{2}) \subset V_c(\frac{1}{2}), \quad \lambda \in \{0, 1\}; \quad V_c(\frac{1}{2}) V_c(\frac{1}{2}) \subset V_c(0) \oplus V_c(1). \quad (23)$$

As we shall see, $\frac{1}{2}$ plays an essential role in constructing nonclassical solutions.

2.5. Generic cubic forms. As we remarked in Subsection 1.3 above, some numerical simulations support the believe that a randomly chosen cubic form is not special. To give this observation a rigorous meaning, we need to formalize what we mean by a generic cubic form. One natural way to do this is to combine the correspondence given in the proceeding section with the concept of a generic algebra introduced recently in [25].

Recall that an algebra over \mathbb{C} is said to be *generic* if it has exactly 2^n idempotents (counting $x = 0$). The idea dates back to the classical work of Segre [48]: given a fixed basis of an arbitrary nonassociative algebra V of dimension $n \geq 2$ over complex numbers \mathbb{C} , one can identify the multiplication with a degree 2 homogeneous map on V . In this setting, the idempotent defining relation $x^2 = x$ becomes a system of quadratic polynomial equations on \mathbb{C}^n . It is well known that a generic (in the Zariski sense) polynomial system has the Bézout number of solutions equal to the product of the principal degrees of the system equations. In our case, the Bézout number is 2^n .

More precisely, let $\{e_i\}_{1 \leq i \leq n}$ be an arbitrary fixed basis of a finite dimensional vector space V . Then any algebra structure on V is uniquely determined

by its multiplication table $M := (a_{ijk})_{1 \leq i, j, k \leq n}$, where $e_i e_j = \sum_{k=1}^n a_{ijk} e_k$. In this notation, the algebra product $f(x, y) := xy$ in V is given component-wise by

$$f_k(x, y) := \sum_{k=1}^n a_{ijk} x_i y_j, \quad k = 1, 2, \dots, n. \quad (24)$$

Also, the idempotent defining relation becomes

$$f(x, x) = x. \quad (25)$$

Then the algebra V is generic if (25) has exactly 2^n distinct solutions over \mathbb{C} . In particular, this fact implies that each solution is a nondegenerate point of (25). The last condition is important because it can be reformulated in terms of the spectrum of the algebra V .

Now it is natural to call a cubic form u on V generic if so is the corresponding algebra $V(u)$.

Note that the above definitions work equally well in both analytic and algebraic settings. However, it is more preferable to work with the latter because the genericity concept is easily translated to the well-developed nonassociative algebra theory, including the Peirce decomposition.

In [25], the following criterion was established.

Proposition 2.6. *A commutative algebra V over \mathbb{C} is generic if and only if the spectrum of any idempotent in V does not contain $\frac{1}{2}$.*

Because of the essential role $\frac{1}{2}$ plays in the present context and also for the reader convenience, we give a sketch of the proof.

Proof. Note that $\tilde{f}(x) := f(x, x)$ is an endomorphism of V homogeneous of degree 2. Since V is commutative, we have

$$\tilde{f}(x + y) - \tilde{f}(x) - \tilde{f}(y) = f(x, y) + f(y, x) = 2f(x, y).$$

In particular, $\frac{\partial \tilde{f}_k}{\partial x_j}(x) = 2f_k(x, e_j)$, which yields

$$D\tilde{f}(x) = 2L_x,$$

where $D\tilde{f}(x)$ is the Jacobi matrix of g at x . Rewriting system (25) as $g_k(x) := \tilde{f}_k(x) - x_k = 0$, $k = 1, 2, \dots, n$, we see that

$$Dg(x) = 2L_x - I.$$

Thus, $Dg(x)$ is nondegenerate if and only if $2L_x - I$ is so, i.e., $\frac{1}{2}$ is not in the spectrum of L_x . In order to finish the proof, it remains to note that system (25) has the maximal finite number of solution if and only if each solution is a nondegenerate point of g , see for instance [12]. \square

We have already seen in Section 2.4 that any Jordan algebra has $\frac{1}{2}$ in its spectrum and plays a distinguished role in the classification of formally real algebras by Jordan, von Neumann, and Wigner, see [24]. In general, the exceptionality of the Peirce eigenvalue $\lambda = \frac{1}{2}$ in the spectrum of many well-established nonassociative algebra structures (including Jordan algebras, Bernstein algebras, and general genetic algebras) is a rather common phenomenon. We mention its very recent appearance in the context of axial algebras (generalizing the Monster algebra of the largest sporadic finite simple group), see [18]. Hsiang algebras emerging in the context of cubic minimal cones present another example of commutative nonassociative algebras where the presence of $\frac{1}{2}$ plays a distinguished role, see Chapter 6 in [36]. In a wider context of algebras with identities, the universality of $\frac{1}{2}$ was discussed recently in [57].

Our next step is to explain why the Peirce eigenvalue $\frac{1}{2}$ is also distinguished in the context of nonclassical solutions.

2.6. The spectral properties of the Hessian of w . It was remarked in Lemma 3.2 in [39] that the hyperbolicity condition (Lemma 2.1 above) for a cubic form u would be fulfilled if there would exist $\delta > 0$ such that for any fixed direction $d \in V$ the spectrum $\lambda_1(d) \leq \dots \leq \lambda_n(d)$ of the quadratic form $\partial_d u(x)$ satisfies

$$\rho(d) := \max\left\{\frac{\lambda_1(d)}{\lambda_3(d)}, \frac{\lambda_n(d)}{\lambda_{n-2}(d)}\right\} < 2 - \delta. \quad (26)$$

Unfortunately (see the remark after Lemma 3.2 in [39]), the above condition is too strong and fails for certain d for *any* cubic form. Indeed, using (16), the quadratic form $\partial_d u(x)$ may be rewritten as

$$\partial_d u(x) = \frac{1}{2}\langle x^2, d \rangle = \frac{1}{2}\langle x, xd \rangle = \frac{1}{2}\langle x, L_d x \rangle,$$

therefore its spectrum coincides with the spectrum of the scaled multiplication operator $\frac{1}{2}L_d$. Now, choosing $d \in E_0$ in the notation of Proposition 2.5, we deduce from (20) and the fact that d is an idempotent that $\lambda_1(d) = 1$ and $\lambda_2(d) \leq \frac{1}{2}$, hence $\rho(d) \geq 2$.

Thus, the failure if the gap inequality (26) is directly related to the Peirce spectrum of certain idempotents³ in the corresponding algebra $V(u)$. Below we try to clarify its nature in more detail and relate it to generic cubic forms.

Let $u \neq 0$ be an arbitrary cubic form on an inner product vector space V . Using (15), we rewrite (8) for $\alpha = 1$ as

$$w(x) = \frac{6u(x)}{|x|} = \frac{\langle x^2, x \rangle}{|x|}.$$

³These idempotents have interesting extremal properties and can be related to Clifford type Jordan algebras, see our recent paper [56]

Then we have for the directional derivative

$$\langle \nabla w(x), y \rangle = \partial_y w|_x = \frac{3\langle x^2, y \rangle |x|^2 - \langle x^2, x \rangle \langle x, y \rangle}{|x|^3}, \quad (27)$$

implying by duality the expression for the gradient

$$\nabla w(x) = \frac{3x^2|x|^2 - \langle x^2, x \rangle x}{|x|^3}.$$

Arguing similarly, we find for the Hessian

$$H(x) := D^2w(x) = \frac{6}{|x|}L_x - \frac{\langle x^2, x \rangle}{|x|^3} - 3\frac{x\widehat{\otimes}x^2}{|x|^3} + 3\frac{x \otimes x}{|x|^5}\langle x^2, x \rangle, \quad (28)$$

where

$$(a \otimes b)(x) = a\langle b, x \rangle, \quad a\widehat{\otimes}b = a \otimes b + b \otimes a.$$

Note that the Hessian is homogeneous degree 0 but it is an *odd* map. i.e., $H(-x) = -H(x)$. Furthermore,

$$\Delta w(x) = \text{tr } D^2w(x) = \frac{6}{|x|} \text{tr } L_x - (n+3)\frac{\langle x^2, x \rangle}{|x|^3}. \quad (29)$$

Recall that u is special if there exists $M \geq 1$ such that $H(x) - H(y)$ is M -hyperbolic for any pair $x, y \in V$. The following elementary observation immediately follows if one chooses $y = -x$ in (27) and refers to the oddness of $H(x)$.

Lemma 2.7. *If w is special, then the set $\{H(x)\}_{|x|=1}$ is hyperbolic.*

The particular case when $x = c$ is an idempotent of $V(u)$ is very special because x^2 and x coincide, implying considerable simplifications in (28), such that the spectrum of $H(c)$ can be calculated explicitly.

Lemma 2.8. *If c is a nonzero idempotent of $V(u)$, then*

$$\text{spectrum}(H(c)) = \left\{ \frac{2}{|c|}, \frac{6\lambda_1 - 1}{|c|}, \dots, \frac{6\lambda_{n-1} - 1}{|c|} \right\} \quad (30)$$

where $1, \lambda_1, \dots, \lambda_{n-1}$ are the eigenvalues of L_c counting the multiplicities. In particular, the characteristic polynomial of $H(c)$ is given by

$$\chi_{H(c)}(t) = \frac{6^n(6|c|t - 2)}{|c|^n(|c|t - 5)} \cdot \chi_c\left(\frac{1 + |c|t}{6}\right) \quad (31)$$

where $\chi_c(z)$ is the characteristic polynomial of L_c .

Proof. Indeed, since $c^2 = c$, from (28) we obtain

$$H(c) = \frac{1}{|c|} \left(6L_c - 1 - \frac{3}{|c|^2} c \otimes c \right). \quad (32)$$

In particular, $H(c)c = \frac{2}{|c|}c$, i.e., c is an eigenvector of $H(c)$ with eigenvalue $\frac{2}{|c|}$. Since $H(c)$ is selfadjoint, the orthogonal complement c^\perp is its invariant subspace. We have

$$H(c) = \frac{1}{|c|} (6L_c - 1) \quad \text{on } c^\perp.$$

This yields (30), and therefore (31). \square

Remark 2.9. As an immediate consequence of (31), we obtain

$$\chi_c(z) = (z - 1) \frac{|c|^n}{6^n} \cdot \frac{\chi_{H(c)} \left(\frac{6z-1}{|c|} \right)}{z - \frac{1}{2}}. \quad (33)$$

We point out a remarkable appearance of the eigenvalue $\frac{1}{2}$ in the denominator. This implies by virtue of Proposition 2.6 that if u is a generic cubic form, then the eigenvalue $\lambda = \frac{2}{|c|}$ must be a simple eigenvalue of the Hessian $H(c)$.

§3. Algebras of Cartan's isoparametric cubics

In this section we consider the principal model example, which clarifies some algebraic ingredients of the construction of nonclassical solutions. We establish the corresponding Peirce decomposition and derive the finiteness of the eiconal algebras in Proposition 3.5.

3.1. Cartan–Münzner equations. Presently, the only known example of a nonclassical solution in the lowest dimension $n = 5$ (see Theorem 1.2) is based on the cubic form u_5 , which naturally appears in the context of isoparametric hypersurfaces in the Euclidean spheres.

According to É. Cartan, a hypersurface of the unit sphere $S^{n-1} \subset \mathbb{R}^n$ is said to be *isoparametric* if it has constant principal curvatures. Isoparametric hypersurfaces have been shown to be useful in various areas of mathematics, see [9] for the modern account of the isoparametric theory. Cartan himself classified all isoparametric hypersurfaces with $g = 1, 2, 3$ constant principal curvatures and also established that all these are homogeneous and algebraic. The celebrated result of Münzner [34] states that $g \in \{1, 2, 3, 4, 6\}$, all five possibilities are realized, and each isoparametric hypersurface with g distinct principal curvatures is obtained as a level set of a homogeneous degree g polynomial on the unit sphere.

The case of $g = 3$ is quite distinguished in many way. In [8] Cartan proved that for any isoparametric hypersurface M^{3d} with $g = 3$ distinct principal curvatures, each principal curvature must have the same multiplicity, and the possible multiplicities are $d = 1, 2, 4, 8$ corresponding the dimensions of classical division algebras \mathbb{A}_d . More precisely, M^{3d} is a tube of constant radius over a standard Veronese embedding of a projective plane into the standard sphere over the division algebra \mathbb{A}_d . Equivalently, M^{3d} is a locus of a cubic form $u(x)$ in $S^{3d+1} \subset \mathbb{R}^{3d+2}$ with u satisfying the Cartan–Münzner system

$$|\nabla u(x)|^2 = 9|x|^4, \quad (34)$$

$$\Delta u(x) = 0. \quad (35)$$

Cartan classified all cubic solutions of (34)–(35) and showed that the corresponding defining cubic polynomials are given explicitly by

$$u = x_1^3 + \frac{3}{2}x_1 \left(|z_1|^2 + |z_2|^2 - 2|z_3|^2 - 2x_2^2 \right) + \frac{3\sqrt{3}}{2}x_2 \left(|z_2|^2 - |z_1|^2 \right) + 3\sqrt{3} \operatorname{Re}(z_1 z_2) z_3, \quad (36)$$

where $x = (x_1, x_2, z_1, z_2, z_3)$ and $z_i \in \mathbb{R}^d \cong \mathbb{A}_d$ and $d \in \{1, 2, 4, 8\}$. We refer to (36) as to a *Cartan isoparametric cubic*. Each Cartan isoparametric cubic also satisfies a determinantal representation like (7) above, where the determinant involved should be properly understood in an appropriate sense. More precisely, u is the generic determinant in the Jordan algebra of (3×3) -Hermitian matrices with entries in the division algebra \mathbb{A}_d .

3.2. Algebras attached to (34). Below, we apply the definitions given in Subsection 2.3 to Cartan isoparametric cubics. A somewhat different approach employing the Freudenthal–Springer construction was suggested in [52].

Our starting point is an arbitrary cubic homogeneous polynomial solution $u(x)$ of the Cartan–Münzner equation (34) alone. By abusing terminology, we call u an *eiconal cubic*. The harmonic eiconal cubics are exactly the Cartan isoparametric cubics. Using (16), we introduce the commutative algebra structure $V(u)$ on $V = \mathbb{R}^n$ equipped with the standard Euclidean inner product. Then $u(x) = \frac{1}{6} \langle x^2, x \rangle$ and (34) becomes

$$\langle x^2, x^2 \rangle = 36 \langle x, x \rangle^2.$$

The exact value of the constant factor 36 is not essential and may be chosen arbitrarily (positive) by a suitable scaling of the inner product.

Now we want to consider an arbitrary algebra satisfying the above identity. This motivates the following definition.

Definition 3.1. A commutative, maybe nonassociative, algebra with a positive definite associating form \langle , \rangle satisfying

$$\langle x^2, x^2 \rangle = \langle x, x \rangle^2 \quad (37)$$

is called an *eiconal* algebra.

If V is an arbitrary eiconal algebra, then the cubic form $u(x) = \frac{1}{6}\langle x, x^2 \rangle$ satisfies the (scaled) eiconal equation (34). This translates the study of (34) into a purely algebraic context.

Note also that the sense of (37) becomes more clear if one introduces the norm $N(x) = \langle x, x \rangle$. Then (37) takes form of the composition algebra identity

$$N(x^2) = N(x)^2. \quad (38)$$

Note, however, that (38) does not imply that $N(xy) = N(x)N(y)$

Our goal is the Peirce decomposition of V . To this end we need the standard linearization technique, which is an important tool in nonassociative algebra, see [32]. More precisely, we linearize (37) at x in the direction y to get

$$4\langle xy, x^2 \rangle = 4|x|^2\langle x, y \rangle.$$

Since the inner product is associating, we have

$$\langle xy, x^2 \rangle = \langle y, x^2x \rangle = \langle y, x^3 \rangle.$$

(Note that by virtue of the commutativity of V the third power in V is well defined: $x^3 = x^2x = xx^2$.) Therefore, $\langle x^3 - \langle x, x \rangle x, y \rangle = 0$ for all $y \in V$, implying by the nondegeneracy of the inner product that $x^3 - \langle x, x \rangle x = 0$ for all x . Conversely, if the last identity holds, one easily gets (37). This proves the following statement.

Proposition 3.2. *An arbitrary commutative algebra with a positive definite associating form \langle , \rangle is eiconal if and only if*

$$x^3 = \langle x, x \rangle x, \quad x \in V. \quad (39)$$

Next, note that the harmonicity condition (35) is equivalent (if we take (18) into account) to the trace free condition

$$\text{tr } L_x = 0, \quad x \in V. \quad (40)$$

A further linearization of (39) in direction $y \in V$ yields

$$x^2y + 2x(xy) = \langle x, x \rangle y + 2\langle x, y \rangle x,$$

and eliminating y , we get

$$L_{x^2} + 2L_x^2 = \langle x, x \rangle + 2x \otimes x. \quad (41)$$

Remark 3.3. The last identity implies that eiconal algebras are ‘nearly Jordan’. Indeed, recall that any Jordan algebra satisfies (21). On the other hand, from (39) and (41) we have

$$[L_{x^2}, L_x^2] = 2[x \otimes x, L_x^2] = 2(x \otimes x^3 - x^3 \otimes x) = 0,$$

i.e., L_{x^2} commutes with L_x^2 . In fact, from [52] it follows that any eiconal algebra has a natural structure of the trace free subspace in a rank 3 Jordan algebra.

3.3. The Peirce decomposition and fusion laws. From the definition of an eiconal algebra V , it follows that $x^2 \neq 0$ for any $x \neq 0$, i.e., V is a nonzero algebra. Therefore, Proposition 2.5 ensures that there are nonzero idempotents in V . Denote by $\text{Idm}(V)$ the set of all nonzero idempotents. Note also that by (37) $|c| = 1$ for any $c \in \text{Idm}(V)$.

The multiplication operator L_c is selfadjoint with respect to the inner product $\langle \cdot, \cdot \rangle$, hence V decomposes into orthogonal sum of L_c -invariant subspaces. We determine the spectrum of L_c . To this end, note that $c^2 = c$, hence c is an eigenvector of L_c with eigenvalue 1, thus, $\mathbb{R}c$ is an invariant subspace of L_c . Therefore, the orthogonal complement c^\perp is also an invariant subspace of L_c . Applying (41), we obtain

$$2L_c^2 + L_c - 1 = 2c \otimes c = 0 \quad \text{on } c^\perp, \quad (42)$$

hence $\text{spectrum}(L_c|_{c^\perp}) \subset \{-1, \frac{1}{2}\}$. Note that from the above inclusion it follows that the eigenvalue 1 has multiplicity one, i.e., any idempotent in V is *primitive*.

Remark 3.4. We point out that the presence of the eigenvalue $\frac{1}{2}$ is crucial for constructing the nonclassical solutions and closely related to the condition of being generic, discussed in Subsections 1.3 and 2.5 above, cf. Lemma 3.2 in [39].

Let $V_\lambda(c)$ be the λ -eigenspace of L_c . Then the Peirce decomposition of V is

$$V = \mathbb{R}c \oplus V_{-1}(c) \oplus V_{\frac{1}{2}}(c).$$

In order to extract the multiplication table (the so-called fusion laws) between the eigenspaces $V_\lambda(c)$, we linearize (41) further. This yields

$$L_{yx} + (L_x L_y + L_y L_x) = \langle x, y \rangle + x \widehat{\otimes} y. \quad (43)$$

Applying (43) to an arbitrary element $z \in V$ yields the full linearization

$$x(yz) + y(zx) + z(xy) = \langle x, y \rangle z + \langle y, z \rangle x + \langle z, x \rangle y. \quad (44)$$

Specializing $z = c \in \text{Idm}(V)$ in (44) and setting $x, y \in c^\perp$ yields

$$(cx)y + x(cy) + c(xy) = \langle x, y \rangle c.$$

Taking the scalar product with z in the last identity and assuming that $x \in V_{\lambda_1}(c)$, $y \in V_{\lambda_2}(c)$, and $z \in V_{\lambda_3}(c)$, where $\lambda_i \in \{\frac{1}{2}, -1\}$, we obtain

$$\langle x_1 x_2, x_3 \rangle (\lambda_1 + \lambda_2 + \lambda_3) = 0. \quad (45)$$

As a corollary, we have

$$V_{\lambda_1}(c)V_{\lambda_2}(c) \perp V_{\lambda_3}(c) \quad \text{whenever } \lambda_1 + \lambda_2 + \lambda_3 \neq 0.$$

For example, setting $\lambda_1 = \lambda_2 = -1$ immediately implies that $V_{-1}(c)V_{-1}(c)$ is perpendicular to both $V_{-1}(c)$ and $V_{\frac{1}{2}}(c)$, hence

$$V_{-1}(c)V_{-1}(c) \subset \mathbb{R}c. \quad (46)$$

Similarly, for $\lambda_1 = -1$ and $\lambda_2 = \frac{1}{2}$ one has $V_{-1}(c)V_{\frac{1}{2}}(c) \subset \mathbb{R}c \oplus V_{\frac{1}{2}}(c)$. On the other hand, since eigenspaces $V_{\lambda}(c)$ are perpendicular for distinct λ we have

$$\langle V_{-1}(c)V_{\frac{1}{2}}(c), c \rangle = \langle V_{-1}(c), V_{\frac{1}{2}}(c)c \rangle = \langle V_{-1}(c), V_{\frac{1}{2}}(c) \rangle = 0,$$

implying that $V_{-1}(c)V_{\frac{1}{2}}(c) \subset V_{\frac{1}{2}}(c)$. Arguing similarly for $\lambda_1 = \lambda_2 = \frac{1}{2}$, one arrives at the fusion (multiplication) laws shown in Table 1.

	V_{-1}	$V_{\frac{1}{2}}$
V_{-1}	$\mathbb{R}c$	$V_{\frac{1}{2}}$
$V_{\frac{1}{2}}$	$V_{\frac{1}{2}}$	$\mathbb{R}c \oplus V_{-1}$

Table 1. Fusion laws of an eiconal algebra

Recall that a linear map $A: X \times Y \rightarrow Y$ such that $A(x, \cdot): Y \rightarrow Y$ is selfadjoint for all $x \in X$ and $A^2(x, \cdot) = \langle x, x \rangle \text{id}_Y$ is called a symmetric Clifford system, cf. [9, 45]. It is well known that in this case $\dim Y$ is even and $\dim X \leq 1 + \rho(\frac{1}{2} \dim Y)$, and the Hurwitz–Radon function ρ is defined by

$$\rho(m) = 8a + 2^b \quad \text{if } m = 2^{4a+b} \cdot \text{odd}, \quad 0 \leq b \leq 3. \quad (47)$$

Proposition 3.5. *If $\dim V \geq 2$, then*

$$\dim V_{-1}(c) - 1 \leq \rho(\frac{1}{2} \dim V_{\frac{1}{2}}). \quad (48)$$

If additionally V satisfies (40), then the possible dimensions of the Peirce subspace $V_{-1}(c)$ coincide with those of classical division algebras. In particular, there are only finitely many isomorphy classes of harmonic eiconal algebras in dimensions $\dim V \in \{5, 8, 14, 26\}$.

Proof. First we show that $V_{-1}(c)$ is nontrivial. Indeed, assume by contradiction that $V_{-1}(c) = \{0\}$. Then $V = \mathbb{R}c \oplus V_{\frac{1}{2}}(c)$, hence $\dim V_{\frac{1}{2}}(c) = n - 1 \geq 1$. From Table 1, $V_{\frac{1}{2}}(c)V_{\frac{1}{2}}(c) \subset \mathbb{R}c$. Hence for any $x \in V_{\frac{1}{2}}(c)$ we have

$$x^2 = \langle x^2, c \rangle c = \langle x, xc \rangle c = \frac{1}{2} \langle x, x \rangle c.$$

Therefore $x^3 = \frac{1}{4} \langle x, x \rangle x \neq \langle x, x \rangle x$, and a contradiction with (39) follows. Thus $\dim V_{-1}(c) \geq 1$. Next, note that for any $x \in V_{-1}(c)$ Table 1 shows that L_x is an endomorphism of $V_{\frac{1}{2}}(c)$. Since

$$x^2 = \langle x^2, c \rangle c = \langle x, xc \rangle c = -\langle x, x \rangle c,$$

by (46), from (41) we find

$$L_x^2 = \frac{3}{4} \langle x, x \rangle \text{ on } V_{\frac{1}{2}}(c). \quad (49)$$

Thus, $\frac{2}{\sqrt{3}}L_x$ is a symmetric Clifford system. This imposes the dimensional obstruction (48). If additionally V satisfies (40), then

$$0 = \text{tr } L_c = 1 + \frac{1}{2} \dim V_{\frac{1}{2}}(c) - \dim V_{-1}(c),$$

thus $\dim V_{\frac{1}{2}}(c) = 2m$, where $m = \dim V_{-1}(c) - 1$, implying by (48) the inequality $\rho(m) \geq m$. It is well known and also easily follows from (47) that the last inequality holds true only if $m = 1, 2, 4, 8$, which implies that $\dim V = 3m + 2$. \square

§4. Concluding remarks and open questions

Minimal cones constitute an important subclass of singular minimal hypersurfaces in the Euclidean space \mathbb{R}^n for $n \geq 4$, playing a crucial role in the study of both the local and global structures of general (regular) minimal hypersurfaces. Many examples known so far of minimal hypercones are algebraic varieties coming essentially from two classical algebraic structures: the Jordan and the Clifford algebras.

An important class is minimal cones defined by a degree three homogeneous polynomials, i.e., cubic minimal cones. These were considered by W. Y. Hsiang in the late 1960s in [21]. On the other hand, all known examples of cubic forms satisfying the hyperbolicity criteria (i) or (ii) (including the examples discussed above and the determinant form u_9 in Remark 2.3) are *exceptional* Hsiang cubic minimal cones (or radial eigencubics in the sense of Chapter 6 in [36]). It is known that there are only finitely many (congruence classes of) such cubics in some distinguished dimensions $5 \leq n \leq 72$, see Table 1 on p. 158 in [36]. From an analytical point of view, any exceptional Hsiang cubic form can be

characterized as a cubic polynomial solution of the following Hessian trace equations:

$$\begin{aligned}\operatorname{tr} D^2 u(x) &= 0, \\ \operatorname{tr}(D^2 u(x))^2 &= C_1 |x|^2, \\ \operatorname{tr}(D^2 u(x))^3 &= C_2 u(x),\end{aligned}\tag{50}$$

Natural questions arise: Does any exceptional Hsiang cubic produce a nonclassical solution? Do there exist special cubics that are not Hsiang eigencubics? In general, it would be interesting to clarify the relationships between minimal cones and construction of nonregular solutions.

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