Classification of the 1st order PDE's

Standard notation:

For u = u(x, y) one denotes the first derivatives by

$$p = u'_x$$

 $q = u'_y$

• The most general 1st order PDE

$$F(x, y, u, p, q) = 0$$

• A **linear** equation

$$a(x, y)u'_x + b(x, y)u'_y = c(x, y)$$

• A homogeneous (linear) equation

$$a(x, y)u'_x + b(x, y)u'_y = 0$$

Generalizations of the linear case:

• A **semilinear** equation

$$a(x,y)u'_x + b(x,y)u'_y = c(x,y,u)$$

• A quasilinear equation

$$a(x, y, u)u'_{x} + b(x, y, u)u'_{y} = c(x, y, u)$$

Fully non-linear equation:

$$F(x, y, u, p, q) = 0$$

with F chosen arbitrarily; then additionally required that

$${F_p'}^2+{F_q'}^2\neq 0$$

1. Characteristics for a homogeneous linear equation



Characteristic equations:

$$\frac{dx}{dt} = a(x, y), \qquad \frac{dy}{dt} = b(x, y)$$

this yields

$$\frac{du(x(t), y(t))}{dt} = 0$$

hence u(x, y) = const along each characteristic curve. In particular u(A) = u(B) on the picture above and one can determine solution (uniquely) if one knows the values of the solution at some points. For instance, if one knows the values along a curve γ which is transversal to characteristic curves:



2. The method of characteristics for general quasilinear equation

$$a(x, y, u)u'_x + b(x, y, u)u'_y = c(x, y, u)$$

If we introduce a vector field $V = (V_1, V_2, V_3)$ with coordinates

$$V_1 = a(x, y, u),$$
 $V_2 = b(x, y, u),$ $V_3 = c(x, y, u)$

then V is orthogonal to the normal vector

$$N_0 = (-u'_x(x_0, y_0), -u'_y(x_0, y_0), 1)$$

at the point $(x_0, y_0, u(x_0, y_0))$ on the graph of a solution z = u(x, y):



- z = u(x, y) are *integral surfaces* of the vector field V
- a characteristic curve (in red)

• a Cauchy data Γ (in green): a curve in \mathbb{R}^3 transversal to the vector field V

Characteristic equations:

$$\frac{dx}{dt} = a(x, y, u), \qquad \frac{dy}{dt} = b(x, y, u), \qquad \frac{dz}{dt} = c(x, y, u)$$

The Cauchy problem: given a curve Γ in \mathbb{R}^3 , find a solution *u* of the 1st order equation whose graph contains Γ :

$$u|_{\gamma} = h(x, y)$$
.

The inviscid Burgers' equation

$$u \, u'_x + u'_y = 0$$

or equivalently

$$\left(\frac{u^2}{2}\right)'_x + u'_y = 0$$

is an example of a conservation law in the general form

$$(G(u))'_x + u'_y = 0$$

Mechanical interpretation: 1D stream of particles is in motion, each particle having constant velocity; a velocity field is given by u(x, y), where y denotes time. If we follow an individual particle, we get a function x = x(t) for which u(x(t), t) remains constant:

$$0 = \frac{d}{dt} \left(u(x(t), t) \right) = u'_x(x(t), t) \cdot \frac{dx(t)}{dt} + u'_y(x(t), t) \cdot \frac{dt}{dt} = u \, u'_x + u'_y(x(t), t) \cdot \frac{dt}{dt} = u \, u'_y($$

An initial velocity (the Cauchy problem):

$$u(x,0) = h(x).$$

• Characteristics lines are

$$x = h(s)t + s,$$
 $y = t,$ $z = h(s)$ (**)

• Then the general solution is u = h(x - uy)

One must distinguish two cases subject to the global behavior of the solution:



Indeed, a simple analysis of (**) shows that two characteristics will intersect if and only if the system

$$x = h(s_1)t + s_1$$
$$x = h(s_2)t + s_2$$

has a *positive* solution t (i.e. for positive time), which is equivalent to saying that

$$t = \frac{s_2 - s_1}{h(s_1) - h(s_2)} > 0$$

If this condition is satisfied for some values s_1 and s_2 then solution suffers a *gradient catastrophe* type of singularity. Otherwise the solution exists globally.

An example of the gradient catastrophe:



Weak solutions for quasilinear equations

Demonstration by an example of a conservation law in the form

$$(G(u))'_{\chi} + u'_{\gamma} = 0 \tag{(*)}$$

where G is some smooth function of u.

The weak form of (*), also called a conservation law, is the following integrated identity (cf. the integral form of the heat equation given in lecture 1)

$$G(u(b,y)) - G(u(a,y)) + \frac{d}{dy} \int_a^b u(x,y) \, dx = 0.$$

Roughly speaking, a weak solution may contain discontinuities, may not be differentiable, and will require less smoothness to be considered a solution than a classical solution. Working with the weak solution of a PDE usually requires that the PDE be reformulated in an integral form. If a classical solution to the problem exists, it will also satisfy the definition of a weak solution.

We consider the simplest case when discontinuity of *u*, also called a shock front, is projected to a smooth curve front $x = \xi(y)$ in the *xy*-plane. In other words, for a fixed *y* function *u* has a jump discontinuity at $x = \xi(y)$:



Then we have the following necessary condition for the shock front $x = \xi(y)$:

Jump condition (or the Rankine-Hugoniot condition):

$$\xi'(y) = \frac{G(u_r) - G(u_l)}{u_r - u_l}$$

Here u_l and u_r denote the limiting values of u from the left and right sides of the shock.