

Partiella differentialekvationer: Solutions to 2009-03-10

Problem 1. Solve the boundary value problem for the Laplace equation in a square

$$\begin{aligned} u''_{xx} + u''_{yy} &= 0, & 0 \leq x \leq 1, & \quad 0 \leq y \leq 1, \\ u(0, y) = u(x, 0) &= 0, & u(1, y) = 2y, & \quad u(x, 1) = 3 \sin \pi x + 2x. \end{aligned}$$

Solution. First we reduce the problem to a homogeneous one. Notice that $2xy$ is a harmonic function and then $u_0 = u - 2xy$ is also harmonic. The boundary conditions for u_0 are found as

$$\begin{aligned} u_0(0, y) &= u(0, y) = 0, \\ u_0(x, 0) &= u(x, 0) = 0, \\ u_0(1, y) &= u(1, y) - 2y = 0, \\ u_0(x, 1) &= u(x, 1) - 2x = 3 \sin \pi x. \end{aligned}$$

Now we can apply the method of separation of variables to u_0 : we try find the solution in the form

$$u_0(x, y) = \sum_{n=1}^{\infty} A_n v_n(x) w_n(y),$$

where v_n and w_n the “elementary solutions”, that is

$$\Delta(v(x)w(y)) = 0.$$

The latter implies $v''(y)w(y) + v(x)w''(y) = 0$, and consequently $\frac{v'''(x)}{v(x)} = -\frac{w''(y)}{w(y)} = C$, where C is some constant. Applying our boundary conditions $v(0) = v(1) = 0$ we obtain that v must be a trigonometric function, moreover a sinus function:

$$v_n(x) = \sin \pi n x, \quad n = 1, 2, 3, \dots$$

This function corresponds to $C_n = -(\pi n)^2$. We have then for the second component

$$w'' - (\pi n)^2 w = 0,$$

that is the solution is found by means of the hyperbolic functions:

$$w_n(y) = A_n \cosh \pi n y + B_n \sinh \pi n y.$$

Then the boundary condition $w(0) = 0$ implies $A_n = 0$. In summary, we have

$$u_0(x, y) = \sum_{n=1}^{\infty} A_n \sin \pi n x \sinh \pi n y.$$

It remains only to find the coefficients. Notice that

$$u_0(x, 1) = 3 \sin \pi x = \sum_{n=1}^{\infty} (A_n \sinh \pi n) \sin \pi n x,$$

hence $A_n \sinh \pi n = 3$ for $n = 1$ and $A_n = 0$ otherwise. This gives

$$u_0(x, y) = \frac{3}{\sinh \pi} \sin \pi x \sinh \pi y,$$

and finally

$$u = u_0 + 2xy = 2xy + 3 \sin \pi x \frac{\sinh \pi y}{\sinh \pi}.$$

Problem 2. Show that $\frac{1}{|x|}$ is locally integrable in \mathbb{R}^3 ($|x| = \sqrt{x_1^2 + x_2^2 + x_3^2}$) and find

$$\lim_{\epsilon \rightarrow +0} \int_{|x| > \epsilon} \frac{\Delta \phi(x)}{|x|} dx_1 dx_2 dx_3$$

for any twice-differentiable function ϕ with compact support.

Solution. Because $f \equiv \frac{1}{|x|}$ is continuous everywhere in the punctured space $\mathbb{R}^3 \setminus \{0\}$ then it is locally integrable there. It remains only to show that f is integrable in any ball

$$B_R = \{x: |x| < R\}.$$

This is equivalent to the existence of absolute convergence of the integral

$$I_\epsilon = \int_{R > |x| > \epsilon} \frac{1}{|x|} dx_1 dx_2 dx_3.$$

We calculate this by the spherical Fubini theorem:

$$I_\epsilon = \int_\epsilon^R dr \int_{S_r} \frac{dS}{|x|} = \int_\epsilon^R dr \int_{S_r} \frac{dS}{r} = \int_\epsilon^R \frac{dr}{r} \text{Area}(S_r) = 4\pi \int_\epsilon^R r dr = 2\pi(R^2 - \epsilon^2).$$

(S_r is the sphere of radius r and with center at the origin). We see that I_ϵ converges as $\epsilon \rightarrow 0$, hence we have proved that $f = \frac{1}{|x|}$ is locally integrable.

In order to find the limit we apply the 2nd Green identity

$$\int_\Omega (v \Delta u - u \Delta v) dx = \int_{\partial\Omega} v \frac{\partial u}{\partial \nu} - u \frac{\partial v}{\partial \nu} dS$$

which is valid for any functions u, v continuous in $\bar{\Omega}$ and twice differentiable in the interior Ω . We take $u = \phi$ and $v = \frac{1}{|x|}$. We notice also that $f = \frac{1}{|x|}$ is harmonic. Indeed,

$$\nabla f = \nabla \frac{1}{r} = -\frac{1}{r^2} \nabla r = -\frac{1}{r^2} \cdot \frac{x}{r},$$

$$\Delta \frac{1}{r} = -\text{div} \left(\frac{1}{r^2} \cdot \frac{x}{r} \right) = -\frac{1}{r^3} \text{div } x + \left\langle x, \frac{3x}{r^5} \right\rangle = -\frac{3}{r^3} + \frac{3}{r^3} = 0.$$

Since ϕ has compact support, it vanishes outside a large ball, say in B_R . Then the Green identity for $\Omega = B_R \setminus B_\epsilon$ reads as

$$\int_{\Omega} \Delta \phi(x) \frac{dx}{|x|} = \int_{\partial \Omega} \frac{\partial \phi}{\partial \nu} \cdot \frac{dS}{r} - \int_{\partial \Omega} \phi(x) \partial_{\nu} \left(\frac{1}{r} \right) dS = -\frac{1}{\epsilon} \int_{|x|=\epsilon} \frac{\partial \phi}{\partial \nu} dS - \frac{1}{\epsilon^2} \int_{|x|=\epsilon} \phi dS$$

Here we used that $\phi(R) = \partial_{\nu} \phi(R) = 0$ and $\partial_{\nu} \left(\frac{1}{r} \right) = -\frac{1}{r^2}$. We have for the integrals: the first integral converges to zero because $\frac{\partial \phi}{\partial \nu}$ is bounded in the unit ball (say by constant M) and the surface integral $\int_{|x|=\epsilon} \frac{\partial \phi}{\partial \nu} dS$ then can be estimated as

$$\left| \frac{1}{\epsilon} \int_{|x|=\epsilon} \frac{\partial \phi}{\partial \nu} dS \right| \leq M \cdot \frac{4\pi\epsilon^2}{\epsilon} \rightarrow 0.$$

For the second integral: for any small δ we chose ϵ such that $|\phi(x) - \phi(0)| < \delta$ if $|x| < \epsilon$ (this is possible since ϕ is continuous at the origin). Then

$$\left| \frac{1}{\epsilon^2} \int_{|x|=\epsilon} \phi(x) dS - 4\pi\phi(0) \right| = \left| \frac{1}{\epsilon^2} \int_{|x|=\epsilon} (\phi(x) - \phi(0)) dS \right| \leq 4\pi\delta.$$

This shows that $\frac{1}{\epsilon^2} \int_{|x|=\epsilon} \phi(x) dS \rightarrow 4\pi\phi(0)$ as $\epsilon \rightarrow 0$.

Thus we have

$$\lim_{\epsilon \rightarrow +0} \int_{|x|>\epsilon} \frac{\Delta \phi(x)}{|x|} dx_1 dx_2 dx_3 = -4\pi\phi(0).$$

Problem 3. Find the radial symmetric *continuous* solution of the Poisson equation

$$\Delta u = 1, \quad x \in \mathbb{R}^3,$$

satisfying $u(1,2,1) = 1$.

Solution. Denote $u(x) = f(r)$, where $r = \sqrt{x_1^2 + x_2^2 + x_3^2}$. Then

$$\nabla f(r) = f'(r) \nabla r = f' \cdot \frac{x}{r},$$

$$\Delta f = \operatorname{div} \left(\frac{f'(r)}{r} x \right) = \frac{3f'(r)}{r} + \left(\frac{f'(r)}{r} \right)' r = f'' + \frac{2}{r} f'.$$

We arrive at the equation

$$f'' + \frac{2}{r} f' = 1,$$

and the Cauchy condition must be determined now by $u(1,2,1) = 1$. We have

$$r(1,2,1) = \sqrt{1 + 4 + 1} = \sqrt{6},$$

hence

$$u(1,2,1) = f(\sqrt{6}) = 1.$$

Solve the above differential equation by substitution $y = f'$:

$$y' + \frac{2}{r} y = 1, \quad y = \frac{r}{3} + \frac{C_1}{r^2} \equiv f', \quad f = \frac{r^2}{6} - \frac{C_1}{r} + C_2.$$

Since the solution must be continuous, it follows that $C_1 = 0$. Moreover, by the Cauchy data,

$$f(\sqrt{6}) = \frac{6}{6} + C_2 = 1, \quad C_2 = 0.$$

We find finally

$$u = \frac{r^2}{6} = \frac{1}{6}(x_1^2 + x_2^2 + x_3^2).$$

Problem 4. Find all solutions of the heat equation

$$u_{xx} - u_t = 0$$

which have the form $u = t^{-\frac{1}{2}}f(t^\alpha x^2)$, where f is a function and α is a real number to be found. (Hint: reduce to a second order ODE with respect to $\xi = t^\alpha x^2$ and find the possible values of α .)

Solution. Denote $\xi = t^\alpha x^2$ and find the derivatives:

$$u_t = -\frac{1}{2}t^{-\frac{3}{2}}f(\xi) + \alpha t^{\alpha-\frac{3}{2}}x^2 f'(\xi) = t^{-\frac{3}{2}}\left(-\frac{1}{2}f(\xi) + \alpha\xi f'(\xi)\right),$$

$$u_x = 2xt^{\alpha-\frac{1}{2}}f'(\xi),$$

$$u_{xx} = 4x^2 t^{2\alpha-\frac{1}{2}}f''(\xi) + 2t^{\alpha-\frac{1}{2}}f'(\xi).$$

Substitution into $u_{xx} - u_t = 0$ then yields

$$4\xi t^{\alpha+1}f''(\xi) + 2t^{\alpha+1}f'(\xi) = -\frac{1}{2}f(\xi) + \alpha\xi f'(\xi),$$

hence

$$t^{\alpha+1} = \frac{-\frac{1}{2}f(\xi) + \alpha\xi f'(\xi)}{4\xi f''(\xi) + 2f'(\xi)}$$

The right hand side of the latter equation depends on $\xi = x^2 t^\alpha$ which is independent variable with respect to t , hence $\alpha + 1 = 0$. We have $\alpha = -1$, therefore in summary,

$$8\xi f''(\xi) + 4f'(\xi) = -(f(\xi) + 2\xi f'(\xi))$$

Dividing both sides by $\xi^{\frac{1}{2}}$ we combine the terms by the Leibnitz product formula

$$8\xi^{1/2}f''(\xi) + 4\xi^{-1/2}f'(\xi) = 8\left(\xi^{\frac{1}{2}}f'(\xi)\right)',$$

$$\xi^{-1/2}f(\xi) + 2\xi^{1/2}f'(\xi) = 2\left(\xi^{\frac{1}{2}}f(\xi)\right)'$$

We have

$$8\left(\xi^{\frac{1}{2}}f'(\xi)\right)' = -2\left(\xi^{\frac{1}{2}}f(\xi)\right)',$$

and after integration

$$4f'(\xi) + f(\xi) = C_1 \xi^{-\frac{1}{2}}.$$

The solution of the homogeneous equation is

$$f_0 = C_2 e^{-\frac{\xi}{4}},$$

and the solution of the nonhomogeneous is found by variation of constants. We find finally the required solution in terms of the Gauss error integral

$$f = C_2 e^{-\frac{\xi}{4}} + C_1 e^{-\frac{\xi}{4}} \int^{\xi} e^{\frac{\xi}{4}} \xi^{-\frac{1}{2}} d\xi = C_2 e^{-\frac{\xi}{4}} + C_3 e^{-\frac{\xi}{4}} \int^{\sqrt{\xi}} e^{-\frac{s^2}{4}} ds.$$

It remains only to substitute $\xi = \frac{x^2}{t}$.

Problem 5. Solve the pure initial value problem for the heat equation

$$u_{xx} + u_{yy} = u_t, \quad u(x, y, 0) = e^{-x^2 - y^2}.$$

Find the maximal temperature at time $t = 2$.

Solution. We have the general representation

$$u(x, t) = \frac{1}{(4\pi t)^{\frac{n}{2}}} \int_{\mathbb{R}^n} e^{-\frac{|x-\xi|^2}{4t}} g(\xi) d\xi.$$

In our case, $n = 2$, $x = (x_1, x_2)$, $\xi = (\xi_1, \xi_2)$ and $g(x_1, x_2) = e^{-x_1^2 - x_2^2}$. Thus

$$u(x, t) = \frac{1}{4\pi t} \int_{\mathbb{R}^2} e^{-\frac{|x-\xi|^2}{4t}} e^{-\xi_1^2 - \xi_2^2} d\xi_1 d\xi_2$$

We have

$$|x - \xi|^2 = |x|^2 - 2\langle x, \xi \rangle + |\xi|^2,$$

and for the integrand:

$$e^{-\frac{|x-\xi|^2}{4t}} e^{-\xi_1^2 - \xi_2^2} = \exp\left(-\frac{|x|^2 - 2\langle x, \xi \rangle + |\xi|^2}{4t} - |\xi|^2\right).$$

Simplify, by completing the square:

$$\begin{aligned} \frac{|x|^2 - 2\langle x, \xi \rangle + |\xi|^2}{4t} + |\xi|^2 &= \frac{|x|^2}{4t} + |\xi|^2 \left(1 + \frac{1}{4t}\right) - \frac{2\langle x, \xi \rangle}{4t} \\ &= \frac{|x|^2}{4t} + \left| \xi \sqrt{1 + \frac{1}{4t}} - \frac{x}{4t\sqrt{1 + \frac{1}{4t}}} \right|^2 - \left| \frac{x}{4t\sqrt{1 + \frac{1}{4t}}} \right|^2 = \end{aligned}$$

$$= \frac{|x|^2}{4t+1} + |\xi a - b|^2,$$

where $a = \sqrt{1 + \frac{1}{4t}}$ and $b = \frac{x}{4t\sqrt{1 + \frac{1}{4t}}}$. This gives

$$u(x, t) = \frac{1}{4\pi t} e^{-\frac{|x|^2}{4t+1}} \cdot \int_{\mathbb{R}^2} e^{-|\xi a - b|^2} d\xi_1 d\xi_2 =$$

(change the variables $\eta = \xi a - b$, the determinant of the Jacobi matrix is a^2)

$$\begin{aligned} &= \frac{1}{4\pi t a^2} e^{-\frac{|x|^2}{4t+1}} \cdot \int_{\mathbb{R}^2} e^{-|\eta|^2} d\eta_1 d\eta_2 = \frac{1}{4\pi t a^2} e^{-\frac{|x|^2}{4t+1}} \cdot (\sqrt{\pi})^2 \\ &= \frac{1}{4t \left(1 + \frac{1}{4t}\right)} e^{-\frac{|x|^2}{4t+1}} = \frac{1}{4t+1} e^{-\frac{|x|^2}{4t+1}}. \end{aligned}$$

This is the required solution.

Now the maximal temperature for $t = 2$ is found by

$$u(x, 2) = \max_{\forall x} \left(\frac{1}{9} e^{-\frac{|x|^2}{9}} \right) = \frac{1}{9}.$$